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Supplementary Information for

Large resistive switching and switchable photovoltaic response in ferroelectric doped BiFeO₃-based thin films by chemical solution deposition

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Topography image



Figure S1. Topography image of the BSFN films.

Ferroelectric resistive switching



Figure S2. Ferroelectric resistive switching for BSFN and BSF thin films with a thickness of 250 nm. The inset shows the curves with semilogarithmic scale.



Figure S3. Ferroelectric resistive switching for BSFN thin films with a thickness of 120 nm. a) Direct current *I-V* characteristics dependent on voltage at room temperature. The inset shows the curves with semilogarithmic scale. b) The resistive ON/OFF as a function of reading voltages, which exhibits a largest value of 194 at a reading voltage of 1.3 V with a small writing voltage of 4 V.

Mechanism	Heterostructure	Ratio	Ref.	
switchable diode	Pt/BFO/SRO/STO	~753	Adv. Mater., 2013, 25, 2339.	
switchable diode	Pt/BFO/Nb-STO	~5000	Appl. Phys. Lett., 2013, 102, 102901.	
switchable diode	Pt/BSFN/FTO	~12000	Current study	
tunnel junction	Pt/BTO/LSMO/NGO	~750	Nature, 2009, 460, 81.	
tunnel junction	Pt/BTO/Nb-STO	~12800	Nat. Mater., 2013, 12, 617.	

Fable S1	Comparis	son betwee	n different	t ferroeled	ctric me	emristor.
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Absorption spectrum



Figure S4. a) UV-visible absorption spectrum of the BSFN films on FTO glass substrate and the bare substrate. The inset shows the photograph. b) Plot of $(F(R_{\infty})E)^2$ vs. *E*. The linear extrapolation gives a band gap of 2.27 eV.



Ultraviolet photoelectron spectroscopy (UPS)

Figure S5. UPS cutoff spectra. a) The work function (φ) of BSFN and FTO. The difference between the two is less than 0.1 eV. The φ depends on the photon energy of UV and the bias voltage (in our case the photon energy of synchrotron radiation is 36.07 eV and the bias voltage is 5 eV). So the φ of BSFN is 36.07 - 26.86 - 5 = 4.21 eV. b) The valence band region of BSFN. $|E_{VB}-E_F|$ is 1.58 eV. From Figure S4, we have obtained that the band gap of BSFN is 2.27 eV. Thus $|E_F-E_{CB}| = 2.27 - 1.58 = 0.69$ ev, which indicates that the present films are n-type semiconductors due to the few existing oxygen vacancies caused by the small Ni substitution. Furthermore, the electron affinity (χ) of films is estimated to be 4.21 - 0.69 = 3.52 eV, consistent with that reported.¹

X-ray photoemission spectroscopy measurements (XPS)

XPS has been used to investigate the oxygen vacancies present in the different films of BSF and BSFN. The XPS measurement of BSFN thin films clearly shows that the main state of the Ni substitution in the lattice is Ni²⁺, which would increase more oxygen vacancies in the BSFN thin films. The deconvolution of the O1s line results in peaks around 529.9 eV and 532.3 eV, corresponding to lattice oxygen ($O_{lattice}$) and combinations of oxygen vacancies with surface adsorbed oxygen ($O_{ads&vac}$). The relative area of $O_{ads&vac}/O_{lattice}$ of the BSFN films is higher than the BSF films, which should be attributed to high concentration of oxygen vacancies with Ni substitution. Furthermore, we cannot obtain large resistive switching in BSF films (Fig. S2), which indicated that the more existence of oxygen vacancies in BSFN films would be important for the ferroelectric polarization modulation.



Figure S6. XPS spectra of the Ni signal of BSFN thin films.



Figure S7. XPS spectra of the O1s signal of BSF and BSFN thin films, respectively.

I-V conduction measurements



Figure S8 The *I-V* conduction curves of up-polarization and down-polarization states for BSFN thin films.

Conduction mechanism

In order to investigation the conduction mechanism of the DC I-V curve at 9 V, we divided the sweeping bias voltage into four parts, $0 \rightarrow 9 \text{ V}$, $9 \rightarrow 0 \text{ V}$, $0 \rightarrow -9 \text{ V}$, $-9 \rightarrow 0$ V, respectively. Figure S9 a_1)- a_3) show three possible conduction mechanisms at the sweeping bias voltage from 0 to 9 V. Log-log plot shows that the conduction is Ohmic at low field from 0 to 1.5 V due to the slope of 0.93 closing to 1.00 (Figure S9 a_1). The slope at higher field is larger than 1 indicating other two mechanisms. Hence, we plotted the semilogarithmic scale curves with log(J) and E, as seen in Figure S9 a_2 . The curve fit with the straight lines from 1.5 to 4.5 V, suggesting that the Schottky barrier model with thermionic emission process is predominant,² which is correspond to the log-log plot at the same voltage range with a large slope of 9.26. At high field from 4.5 to 9 V, Fowler-Norhdeim (FN) plot shows FN tunneling domination (Figure S9 a₃). As shown in Figure S9 b, the conduction mechanism is space-charge limited current (SCLC) for the sweeping bias voltage from 9 to 0 V, according that the relation between J and E satisfies the expression of $J^{\infty}E^{2,3}$ Meanwhile, the conduction mechanisms at the sweeping bias voltage from 0 to -9 V are SCLC, Ohmic, Schottky, and FN from 0 to -2.4 V, -2.4 to -3.6 V, -3.6 to -5.1 V, and -5.1 to -9 V, respectively. Also the conduction behavior follows Lampert's theory of SCLC from -9 to 0 V.







Figure S9. Conduction mechanism of the DC *I-V* curve at 9 V. a_1)- a_3) for $0 \rightarrow 9$ V, b) for $9 \rightarrow 0$ V, c_1)- c_3) for $0 \rightarrow -9$ V, d) for $-9 \rightarrow 0$ V, respectively. e) The summary of conduction mechanism.

References:

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