

**Electronic Structure of Germanium Phosphide Monolayer and Li-diffusion in Its  
Bilayer**

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### **Description 1: Li-diffusion along $a$ direction on the GeP monolayer**

Here, we investigated the Li diffusion along the  $a$  ( $//X$ ) direction. It will be quite reasonable to assume that the diffusion follows stepwise paths along the hollow sites in Figure 2, i.e.,  $A \rightarrow B \rightarrow E \rightarrow C \rightarrow F \rightarrow D \rightarrow A$ , repeatedly. Figure 4 shows that the activation barriers are 0.09, 0.34, 0.08, 0.12, 0.04, and 0.01 eV for steps  $A \rightarrow B$ ,  $B \rightarrow E$ ,  $E \rightarrow C$ ,  $C \rightarrow F$ ,  $F \rightarrow D$ , and  $D \rightarrow A$ , respectively. We note that site E practically corresponds to the saddle point.

Therefore, the rate-determining step is  $B \rightarrow C$ , and the activation barrier ( $E_a$ ) for the diffusion along the  $a$  axis is 0.43 eV. In turn, the barrier is inaccessible at room temperature, because the Boltzmann factor corresponding to the fraction of Li atoms with the barrier is only  $5 \times 10^{-8}$ . In short, the Li atoms cannot diffuse along the armchair path on the monolayer.

### **Description 2: Li-diffusion path $P_b^1$ in the interlayer region of GeP bilayer**

Considering that site A is the most stable for the monolayer, we began our discussion on four nonequivalent steps to adjacent sites starting from site  $A_1F_2$ , i.e.,  $A_1F_2 \rightarrow B_1T_2$ ,  $A_1F_2 \rightarrow D_1T_2$ ,  $A_1F_2 \rightarrow T_1C_2$ , and  $A_1F_2 \rightarrow T_1D_2$ . Our NEB calculation shows that the diffusion occurs along the first path exclusively, considering that their activation barriers are 0.24, 0.52, 0.53, and 0.57 eV, respectively. [Here, we recall that  $A_1F_2 \rightarrow D_1T_2$  and  $A_1F_2 \rightarrow T_1D_2$  steps are inequivalent, although  $D_1T_2$  and  $T_1D_2$  sites are equivalent.] Next, there are again four inequivalent steps from  $B_1T_2$  to adjacent sites, i.e.,  $B_1T_2 \rightarrow A_1F_2$ ,  $B_1T_2 \rightarrow E_1E_2$ ,  $B_1T_2 \rightarrow T_1C_2$ , and  $B_1T_2 \rightarrow T_1D_2$ , for which barriers are 0.07, 0.35, 0.37, and 0.71 eV, respectively. Therefore, the preferred path will be  $A_1F_2 \rightarrow B_1T_2 \rightarrow A_1F_2$  repeatedly with barriers of 0.24, 0.07, and 0.24 eV, as shown in Figure 6(a).

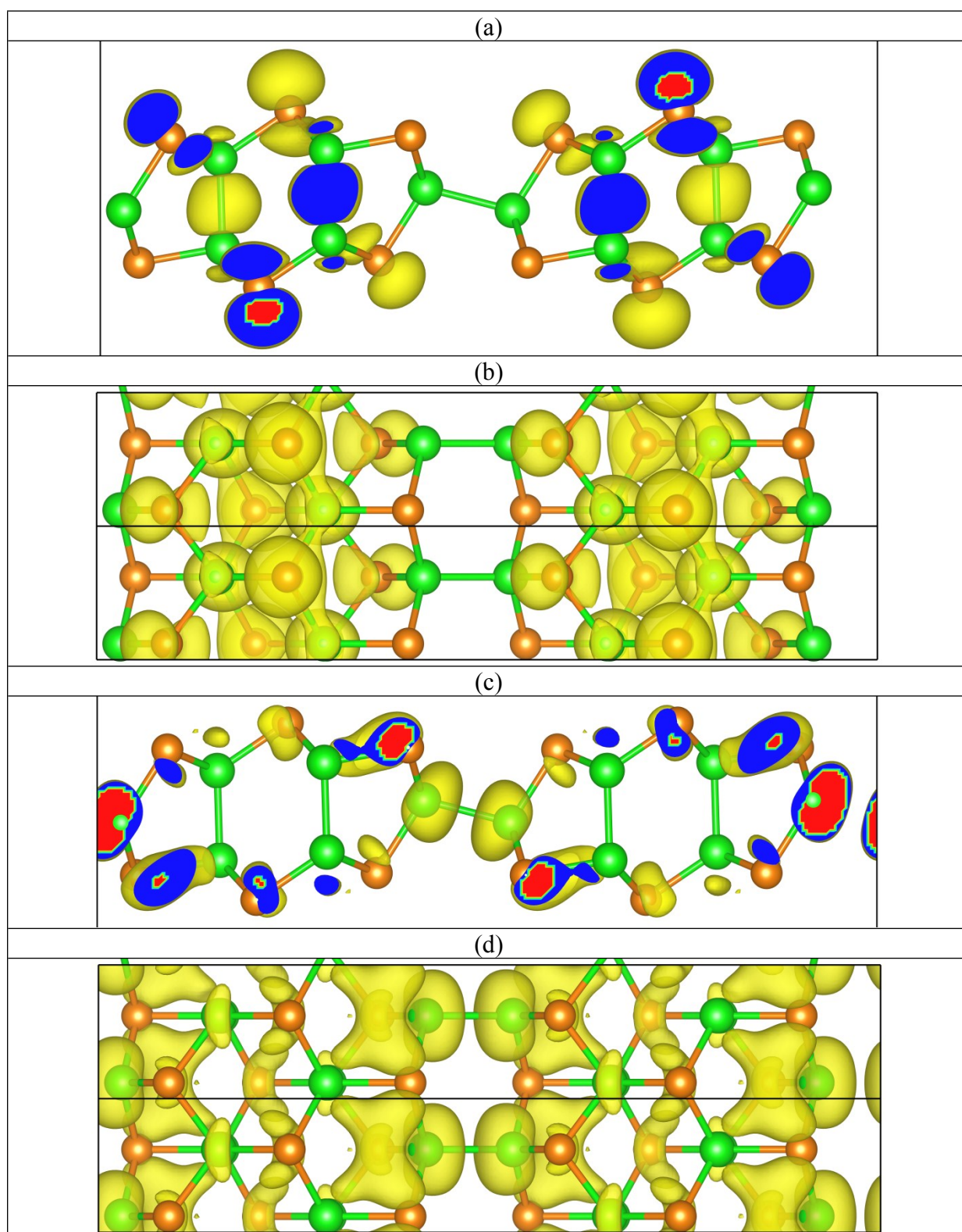
### **Description 3: Alternative Li-diffusion paths connecting $P_b^1$ to $P_b^2$ in the interlayer region of GeP bilayer**

Next, we focused on various paths connecting  $P_b^1$  to  $P_b^2$  starting from  $B_1T_2$ . Its first step will be  $B_1T_2 \rightarrow E_1E_2$ , considering that its barrier is comparable to that for graphene bilayer. There are also four inequivalent steps from  $E_1E_2$ :  $E_1E_2 \rightarrow T_1B_2$ ,  $E_1E_2 \rightarrow B_1T_2$ ,  $E_1E_2 \rightarrow C_1T_2$  and,  $E_1E_2 \rightarrow T_1C_2$ , for which barriers are 0.19, 0.20, 0.28, and 0.30, respectively, indicating that all of them are accessible. This is because their activation barriers are lower than or comparable to that ( $= 0.34$  eV) for graphite. Now, we can define alternative paths connecting  $P_b^1$  and  $P_b^2$  paths. The one with the lowest barrier is denoted by  $P_{a1}^1$ :  $B_1T_2 \rightarrow E_1E_2 \rightarrow T_1B_2$ , for which stepwise barriers of 0.35 and 0.19 eV are certainly higher than those for a set of paths  $P_b$  by at least 0.11 eV. Consideration of the Boltzmann factor suggests that the diffusion from  $P_b^1$  to  $P_b^2$  along the  $a$  axis through  $P_{a1}^1$  will be  $\sim 100$  times slower than those along the  $b$  axis. It is clear that there are other less accessible paths connecting  $P_b^1$  to  $P_b^2$  in terms of the activation barriers comparable to that for graphite bilayer, such as  $B_1T_2 \rightarrow E_1E_2 \rightarrow C_1T_2 \rightarrow F_1A_2$ .

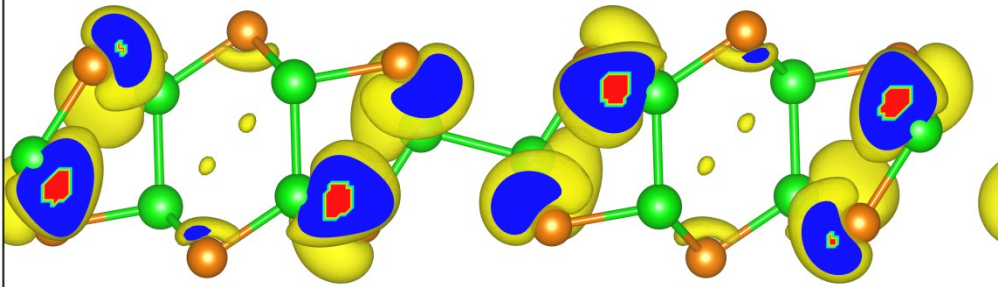
**Description 4: an alternative Li-diffusion path connecting  $P_b^2$  to  $P_b^3$  in the interlayer region of GeP bilayer**

For comparison, we have also considered another path starting from  $T_1B_2$  instead. Along the same direction, there is only one inequivalent step:  $T_1B_2 \rightarrow D_1T_2$ , for which the barrier of 0.64 eV is even larger.

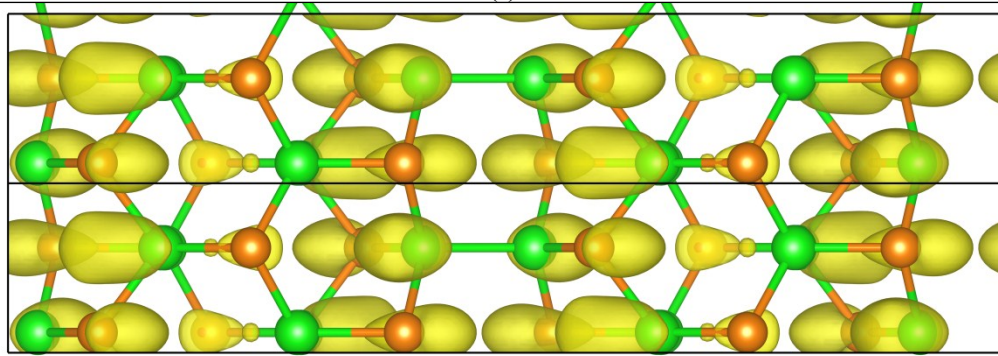
**Figure S1.** Two different views of the charge density plot for the VBM at  $S'$  point (a,b), the CBM at  $Y''$  (c,d), the second CBM at the  $\Gamma$  points (e,f), and the VB-1 at  $\Gamma$  (g,h). The brown and green colors represent phosphorus (P) and germanium (Ge) atoms, respectively.



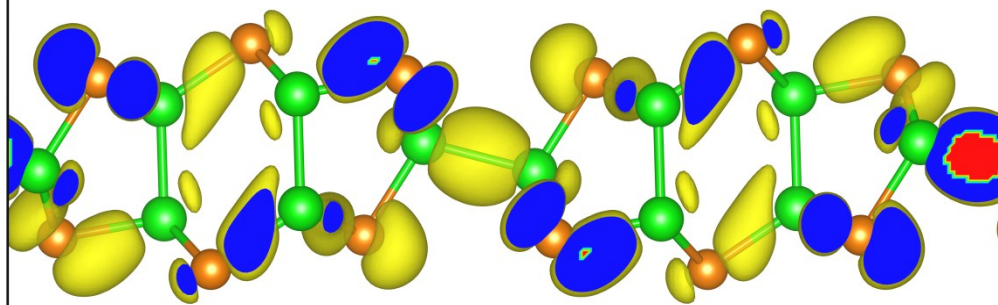
(e)



(f)



(g)



(h)

