

## Molecular and dissociative O<sub>2</sub> adsorptions on the Cu<sub>2</sub>O(111) surface

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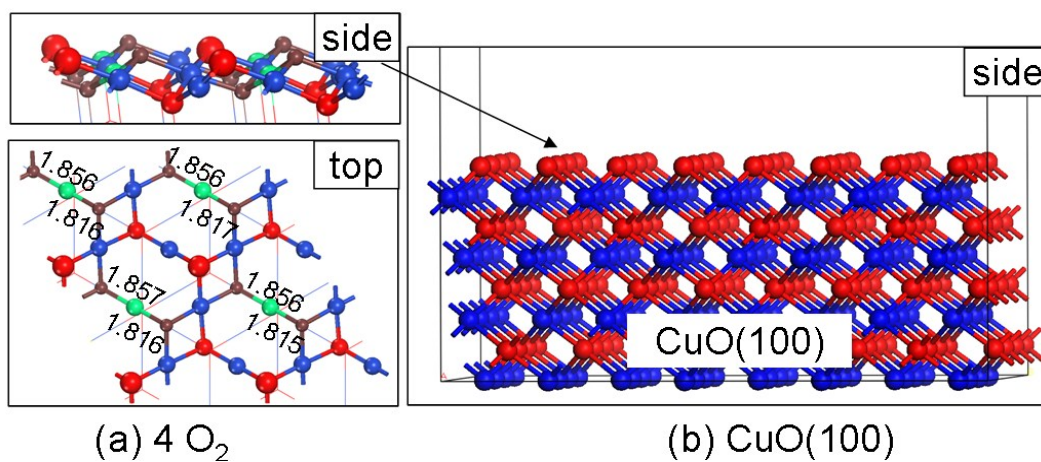
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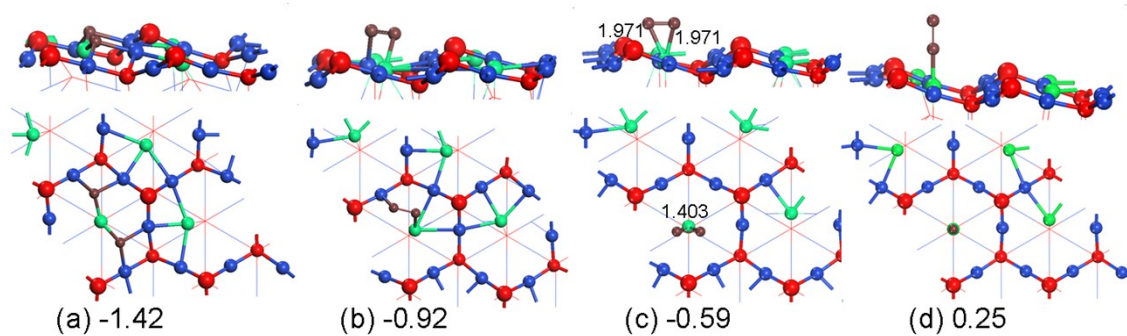
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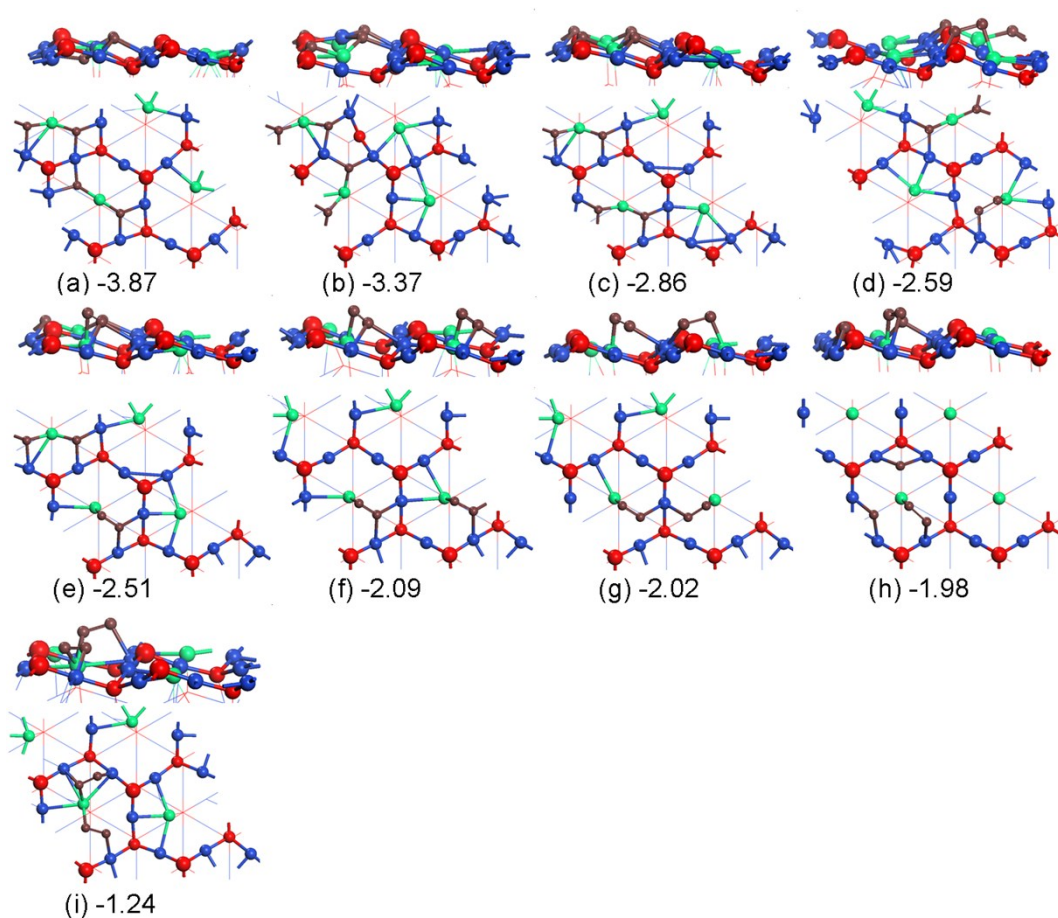
**Fig. S1** (a) Most stable configuration of four  $O_2$  adsorption on the  $Cu_2O(111)$  surface and (b)  $CuO(100)$  surface. (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed  $O_2$  in small brown ball.).



**Fig. S2** Least stable configurations of one O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface. (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball. Adsorption energies in eV).

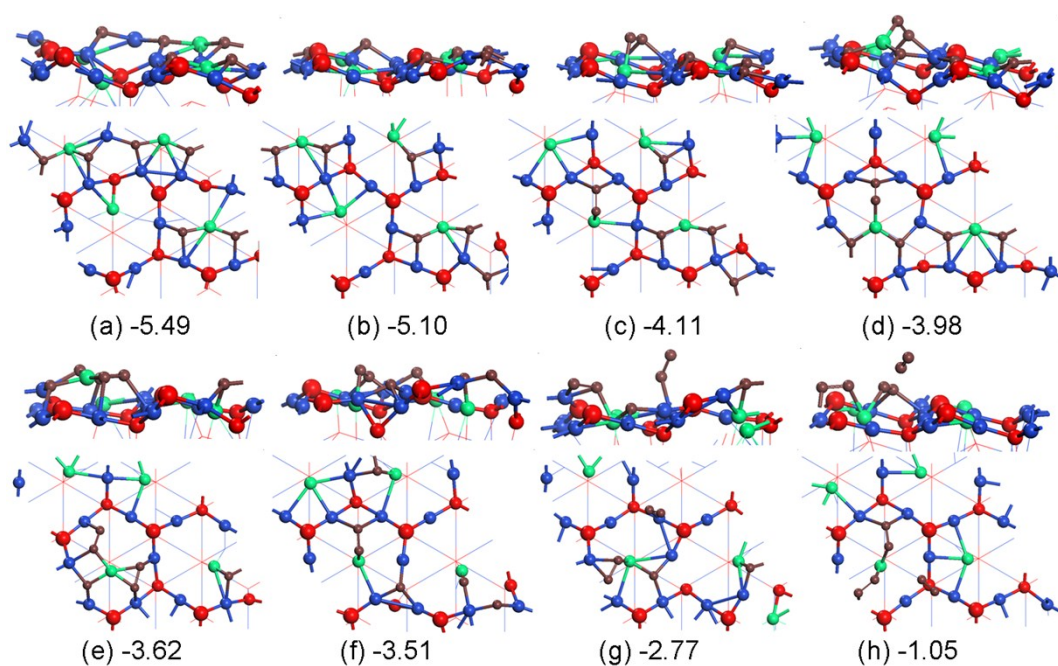


**Fig. S3** Least stable configurations of the two O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball. Adsorption energies in eV).

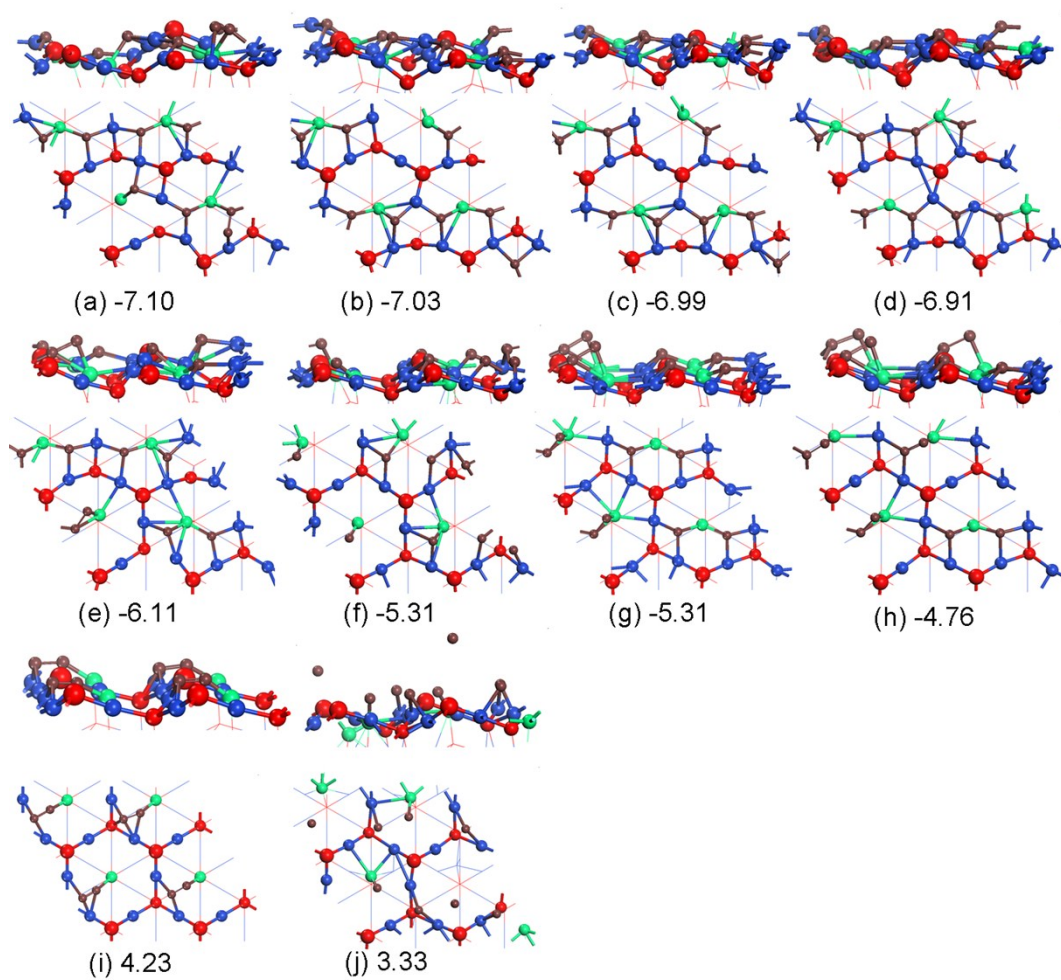


**Fig. S4** Metastable configurations of the three O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface. (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball.

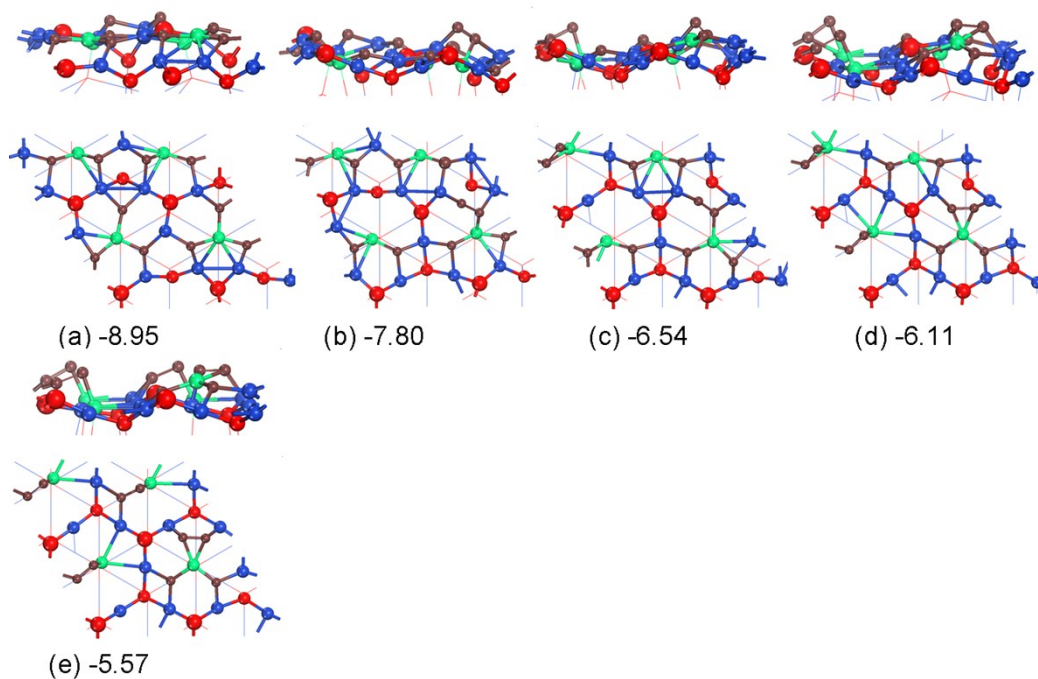
Adsorption energies in eV).



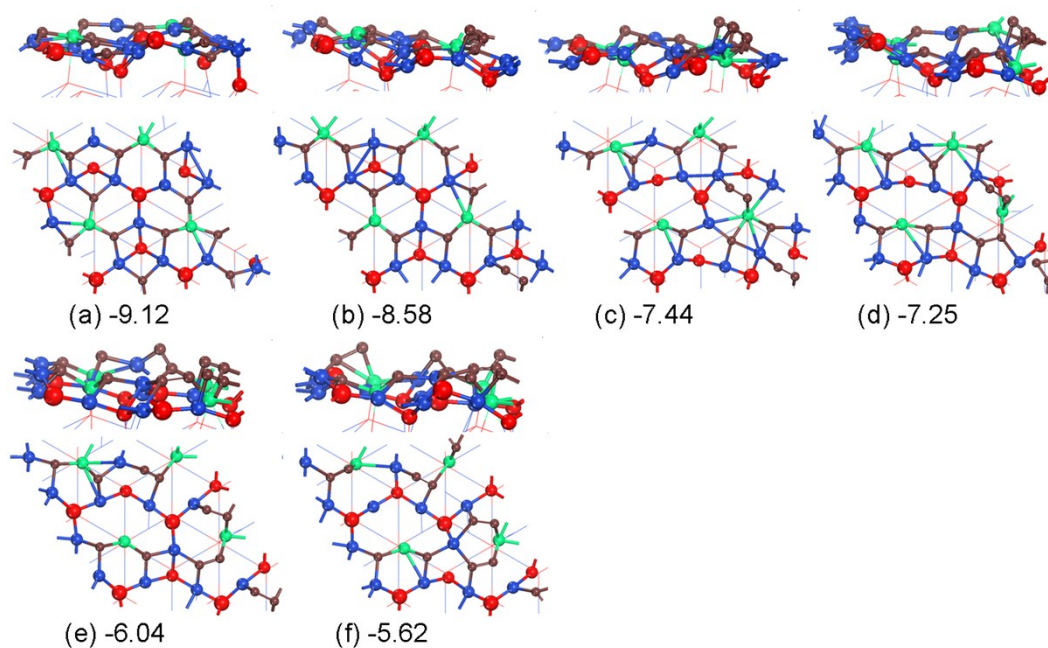
**Fig. S5** Metastable configurations of the four O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball. Adsorption energies in eV).



**Fig. S6** Metastable configurations of the five O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball. Adsorption energies in eV).

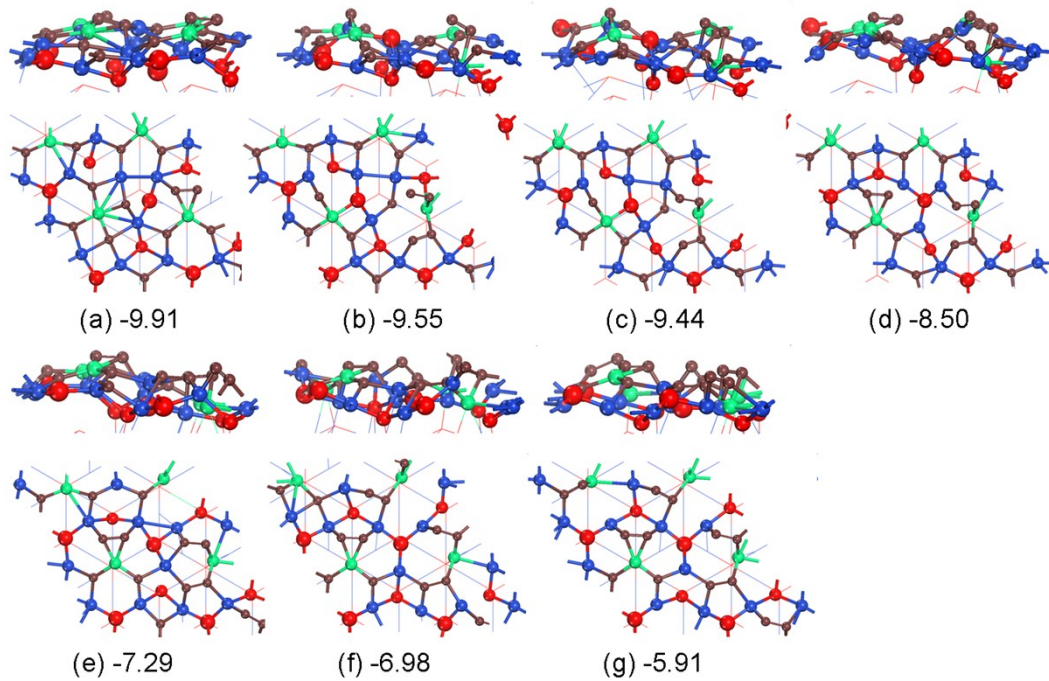


**Fig. S7** Metastable configurations of the six O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball. Adsorption energies in eV).





**Fig. S8** Metastable configurations of the seven O<sub>2</sub> adsorption on the Cu<sub>2</sub>O(111) surface (coordinated saturated Cu atoms in blue ball, coordinated unsaturated Cu atoms in green ball, first layer coordinated unsaturated O atoms in big red ball, third coordinated saturated O atoms in red ball, O atoms of adsorbed O<sub>2</sub> in small brown ball. Adsorption energies in eV).



The full description of Boltzmann statistics and atomistic thermodynamics is shown as follows.

According to atomic thermodynamics, the Gibbs free energy ( $\Delta G$ ) of  $nO_2$  adsorption on the Cu<sub>2</sub>O(111) surface is chosen as the criterion and defined as:

$$\Delta G_{Cu_2O}^{ads}(T, P, nO_2) = G[nO_2/Cu_2O(111)] - G[Cu_2O(111)] - nG_{O_2}(T, p_{O_2}) \quad (1)$$

where  $G[nO_2/Cu_2O(111)]$ ,  $G[Cu_2O(111)]$  and  $G_{O_2}(T, p_{O_2})$  are the Gibbs free energies of adsorbed system, support and adsorbates, respectively,  $n$  is number of adsorbed O<sub>2</sub> molecules,  $T$  is the temperature,  $p_{O_2}$  is the partial pressure of O<sub>2</sub> in the gas atmosphere. It is clear that a more

negative  $\Delta G$  indicates the more stable adsorption structure. We can write  $G_{O_2}(T, p_{O_2})$  as:

$$G_{O_2}(T, p_{O_2}) = E_{O_2}^{\text{total}} + \mu_{O_2}(T, p^0) + k_B T \ln(p/p^0) \quad (2)$$

where  $E_{O_2}^{\text{total}}$  is the total energy of  $O_2$  molecules including zero point vibration energy,  $\mu_{O_2}(T, p^0)$  term includes vibrational and rotational contributions for  $O_2$  gas, and can be taken from tables of thermodynamic data<sup>1</sup>.  $k_B T \ln(p/p^0)$  is the contribution of temperature and  $O_2$  partial pressure to the  $H_2$  chemical potential and  $k_B$  is the Boltzmann constant. The vibration contribution of the  $Cu_2O(111)$  surface with adsorbed  $O_2$  was included in the DFT-calculated total energy to substitute the Gibbs free energies of the solid surface since frequency motions of the surface slab have a large entropy contribution to the surface. Then, equation (1) can be rewritten as:

$$\Delta G_{Cu_2O}^{\text{ads}}(T, P, nO_2) = G[nO_2/Cu_2O(111)] - E[Cu_2O(111)] - nG_{O_2}(T, p_{O_2}) \quad (3)$$

Inserting equation (2) into (3), we can derive:

$$\Delta G_{Cu_2O}^{\text{ads}}(T, P, nO_2) = G[nO_2/Cu_2O(111)] - E[Cu_2O(111)] - nE_{O_2}^{\text{total}} - n\mu_{O_2}(T, p^0) - nk_B T \ln(p/p^0) \quad (4)$$

Each adsorption configuration has a probability of occurrence,  $P_m$ , which is a function of the temperature:  $P_m = 1/Z \exp[-\Delta E/(k_B T)]$ , where  $Z$  is the canonical partition function,  $\Delta E$  is the adsorption energy of a single  $O_2$  adsorption on the  $Cu_2O(111)$  surface, and  $k_B$  is the Boltzmann constant, respectively.

#### References:

- (1) Stull, D. R.; Prophet, H. *JANAF thermochemical tables*, DTIC Document, 1971.