Morphological transitions of axially-driven microfilaments

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Supplemental Materials

We derived an elasto-hydrodynamic model of a filament driven by active force with an arbitrary force profile. The governing equations in dimensionless form is

$$\theta_t = -\theta_{ssss} - \Lambda \theta_{ss} - \left(1 + \frac{1}{\gamma}\right) \Lambda_s \theta_s + 3\left(1 + \frac{1}{\gamma}\right) \theta_s^2 \theta_{ss} + \frac{1}{\gamma} f_{\parallel} \theta_s + \frac{\partial f_{\perp}}{\partial s}, \quad (1a)$$

$$\Lambda_{ss} - \gamma \theta_s^2 \Lambda = 3\theta_{ss}^2 + (3+\gamma)\theta_s \theta_{sss} - \gamma \theta_s^4 + \frac{\partial f_{\parallel}}{\partial s} - \gamma f_{\perp} \theta_s.$$
(1b)

The boundary conditions for the filament with free-ends are given by the moment- and force-free conditions at both ends s = 0 and s = 1, leading to

$$\theta_s(0) = \theta_{ss}(0) = \theta_s(1) = \theta_{ss}(1) = 0,$$
(2a)

$$\Lambda(0) = \Lambda(1) = 0. \tag{2b}$$

We consider a force at $s = s_p$ acting along the local tangential direction, with force profile given by a Dirac delta function,

$$f_{\parallel} = f_{\rm p}\delta(s - s_{\rm p}), \quad f_{\perp} = 0.$$
(3)

We discretize the arc length s uniformly into N segments and apply the second-order central difference formulae for all the derivatives of θ and Λ , in the bulk,

$$(X_s)_j = \frac{1}{2\Delta s} (-X_{j-1} + X_{j+1}),$$
 (4a)

$$(X_{ss})_j = \frac{1}{\Delta s^2} (X_{j-1} - X_j + X_{j+1}),$$
 (4b)

$$(X_{ssss})_j = \frac{1}{\Delta s^4} (X_{j-2} - 4X_{j-1} + 6X_j - 4X_{j+1} + X_{j+2}),$$
(4c)

where $\Delta s = 1/N$, j = 2, 3...N, and X represents either θ or Λ . At the boundaries, the discretization is

$$(X_s)_1 = \frac{1}{2\Delta s} (-3X_1 + 4X_2 - X_3),$$
 (5a)

$$(X_{ss})_1 = \frac{1}{\Delta s^2} (2X_1 - 5X_2 + 4X_3 - X_4),$$
 (5b)

$$(X_s)_{N+1} = \frac{1}{2\Delta s} (3X_{N+1} - 4X_N + X_{N-1}),$$
(5c)

$$(X_{ss})_{N+1} = \frac{1}{\Delta s^2} (2X_{N+1} - 5X_N + 4X_{N-1} - X_{N-2}).$$
(5d)

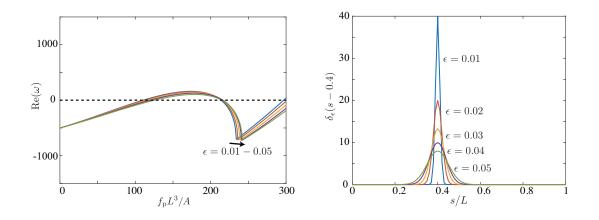


FIG. 1. Left: Real part of dominant eigenvalues when $s_{\rm p} = 0.4$, with regularization coefficient $\epsilon = 0.01 - 0.05$. Right: Regularized Dirac delta function with $\epsilon = 0.01 - 0.05$.

At each time step, we first calculate the tension Λ from (1b), then march (1a) in time with a modified Crank-Nicolson method. We treat the linear terms implicitly and the nonlinear terms explicitly. At time step k, the discrete equations are given by

$$\Lambda_{ss}^{k} - \gamma \left(\theta_{s}^{k}\right)^{2} \Lambda^{k} = 3(\theta_{ss}^{k})^{2} + (3+\gamma)\theta_{s}^{k}\theta_{sss}^{k} - \gamma(\theta_{s}^{k})^{4} + f_{p}\frac{\partial\delta(s-s_{p})}{\partial s}$$

$$\frac{\theta^{k+1} - \theta^{k}}{\Delta t} = -\frac{1}{2} \left(\theta_{ssss}^{k} + \theta_{ssss}^{k+1}\right) - \frac{1}{2} \left[\Lambda^{k} \left(\theta_{ss}^{k} + \theta_{ss}^{k+1}\right)\right]$$

$$-\frac{1}{2}(1+\gamma^{-1}) \left[\Lambda_{s}^{k} \left(\theta_{s}^{k} + \theta_{s}^{k+1}\right)\right] + \frac{3}{2}(1+\gamma^{-1}) \left[\left(\theta_{s}^{k}\right)^{2} \left(\theta_{ss}^{k} + \theta_{ss}^{k+1}\right)\right]$$

$$+\frac{1}{2\gamma} f_{p}\delta(s-s_{p}) \left(\theta_{s}^{k} + \theta_{s}^{k+1}\right).$$
(6a)
(6b)

We use N = 100 in our simulations, the time step is around 10^{-4} . To avoid the singularity, the Dirac delta function is regularized using the regularization function

$$\delta_{\epsilon}(x) \approx \frac{1}{\sqrt{2\pi\epsilon}} e^{-\frac{x^2}{2\epsilon^2}}.$$
(7)

In Fig. 1, we calculate the eigenvalues for $s_p = 0.4$ (shown in Fig. 5 of the main document) with regularization coefficient $\epsilon = 0.01-0.05$. Clearly, the instability features are not affected qualitatively by the exact value of ϵ . In the main document, we set $\epsilon = 0.01$ throughout.

We next choose a finer discretization N = 200 with $\epsilon = 0.01$. The results show the same level of accuracy, see Fig. 2.

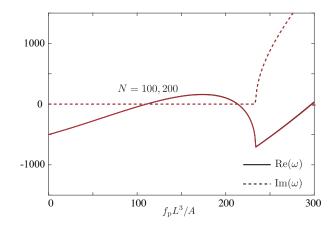


FIG. 2. Dominant eigenvalues when $s_p = 0.4$, with regularization coefficient $\epsilon = 0.01$. Two sets of line overlap, with number of discretization N = 100 and 200 respectively.