

## Electronic Supplementary Information

### Intercalation of Cobaltocene into WS<sub>2</sub> Nanosheets for Enhanced Catalytic Hydrogen Evolution Reaction

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#### **IV. References**

## I. Experimental Section

**Synthesis.** All reactants were purchased from Sigma-Aldrich. 88 mg (0.254 mmol) of Ammonium tetrathiotungstate ((NH<sub>4</sub>)<sub>2</sub>WS<sub>4</sub>, molecular weight (MW) = 260.28 g mol<sup>-1</sup>, 99.97%) and 4.8-48 mg (0.0254-0.254 mmol) of cobaltocene (bis(cyclopentadienyl) cobalt (II), Co(C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>, MW = 189.12 g mol<sup>-1</sup>) are added into 10 mL of oleylamine (C<sub>18</sub>H<sub>35</sub>NH<sub>2</sub>, OA, MW = 267.49 g mol<sup>-1</sup>) in a 50 mL three-neck flask. Non-intercalated WS<sub>2</sub> nanosheets were also synthesized under the same growth condition without CoCp<sub>2</sub>. The thermal decomposition of (NH<sub>4</sub>)<sub>2</sub>WS<sub>4</sub> would produce WS<sub>2</sub>, which is described as (i) (NH<sub>4</sub>)<sub>2</sub>WS<sub>4</sub> → WS<sub>3</sub>+2NH<sub>3</sub>+H<sub>2</sub>S; (ii) WS<sub>3</sub> → WS<sub>2</sub> + S.<sup>S1</sup> The CoCp<sub>2</sub> can catalyse the reduction step (ii) by injection of the electron into the WS<sub>3</sub>. The reaction mixture is degassed at 100 °C for 30 min. Then, the solution is heated to 280 °C under Ar environment at a rate of 10 °C min<sup>-1</sup>. When the temperature reaches to 280 °C, the color of reaction solution turns from yellowish green to dark black. The reaction temperature is maintained at 280 °C for an hour. Next, anhydrous ethanol is added into the reaction mixture while cooling down the reaction mixture to room temperature. The reaction solution is precipitated using centrifugation at 6000 rpm for 10 min. Purification process via precipitation is conducted by 5 times using ethanol (2) toluene (2), and acetone (1). The black powders were obtained after drying under vacuum.

**Characterization.** The products were characterized by scanning electron microscopy (SEM, Hitachi S-4700), and field-emission transmission electron microscopy (FE TEM, FEI TECNAI G2 200 kV, Jeol JEM 2100F, HVEM). Energy-dispersive X-ray fluorescence spectroscopy (EDX) with elemental maps was measured using a TEM (FEI Talos F200X) operated at 200 keV that equipped with high-brightness Schottky field emission electron source (X-FEG) and Super-X EDS detector system (Bruker Super-X). Fast Fourier-transform (FFT) images were

generated by the inversion of the TEM images using Digital Micrograph GMS1.4 software (Gatan Inc.).

High-resolution TEM (HRTEM)/STEM system equipped with a probe Cs corrector (JEM-ARM 200F with Schottky type FEG operated at 200 kV equipped with CEOS Cs-corrector) was used to obtain the high-angle annular dark-field (HAADF) and annular bright-field (ABF) STEM images. The electron energy loss spectroscopy (EELS) experiments were carried out using an Enfina system from Gatan Inc., respectively. Cs-corrected STEM analysis was also carried out using a Titan 80-300<sup>TM</sup> (FEI, The Netherlands) microscope operated at 300 kV. The STEM convergence semi-angle ( $\alpha$ ) used was  $\sim 18$  mrad. The STEM minimum and maximum acceptance semi-angles ( $\beta$ ) were  $\sim 20$  and  $122$  mrad, respectively.

High-resolution X-ray diffraction (XRD) patterns were obtained using the 9B and 3D beamlines of the Pohang Light Source (PLS) with monochromatic radiation ( $\lambda = 1.54595$  Å). XRD pattern measurements were also carried out in a Rigaku D/MAX-2500 V/PC using Cu K $_{\alpha}$  radiation ( $\lambda = 1.54056$  Å). X-ray photoelectron spectroscopy (XPS) measurements were performed using the 8A1 beam line of the PLS, as well as a laboratory-based spectrometer (Thermo Scientific Theta Probe) using a photon energy of  $1486.6$  eV (Al K $_{\alpha}$ ).

X-ray absorption near edge spectra (XANES) and extended X-ray absorption fine structure (EXAFS) spectra at the W L $_3$ -edge and Co K-edge were collected in transmission mode using the 10C beam line of the PLS with a ring current of  $350$  mA at  $3.0$  GeV. Energy calibration was carried out by simultaneously measuring the reference spectrum of W and Co metal foils. Least-squares fits of EXAFS data were performed using the Athena and Artemis software packages, version 0.9.25.

Inductively coupled plasma atomic emission spectrometry (ICP-AES) was performed on a ThermoFisher iCAP-7600. Solid-state  $^{13}\text{C}$  (100.64 MHz) NMR spectra were acquired on a Bruker AVANCE II<sup>+</sup> 400 MHz NMR system (at the KBSI Seoul Western Center) equipped with a Bruker 3.2 mm bore HXY probe operating in HX mode. The magic angle spinning  $^{13}\text{C}$  NMR experiments (one pulse method) were performed using a pulse length of 2  $\mu\text{s}$  for a  $\pi/2$  pulse length of 5  $\mu\text{s}$ , and a pulse repetition delay time of 3 s. The spectra were referenced to an external adamantane standard in which the peak at higher chemical shift was set at 38.43 ppm. The spectra were processed using the Bruker Topspin software (version 3.2) using conventional techniques, and a 50 Hz line broadening window function was applied in all cases.

Electron paramagnetic resonance (EPR) measurements were performed on a Bruker EMX-Plus spectrometer at room temperature. The samples (4 mg) were loaded in a quartz tube. The microwave frequency was 9.644564 GHz, and the microwave power was fixed to 20 mW to avoid saturation. Attenuated total reflectance Fourier transform infrared (ATR FTIR) spectra were obtained on a LabRam ARAMIS IR<sup>2</sup> (HORIBA JOBIN YVON). Spectra were collected by pressing the sample onto a diamond crystal with a pressure setting of 90 on the DuraScope. Each spectrum consisted of 16 spectra co-added accumulated between 4000  $\text{cm}^{-1}$  and 560  $\text{cm}^{-1}$  with a spectral resolution of 4  $\text{cm}^{-1}$ . Raman spectra were measured with a micro-Raman spectrometer (Horiba ARAMIS IR2), using a diode laser with an excitation wavelength of 532 nm.

***Electrochemical Measurements.*** Experiments were carried in a three-electrode cell connected to an electrochemical analyzer (CompactStat, Ivium Technologies). HER electrocatalysis (in 0.5 M  $\text{H}_2\text{SO}_4$  electrolyte) was measured using a linear sweeping from 0 to -0.8 V (vs. RHE) with a scan rate of 2  $\text{mV s}^{-1}$ . A saturated calomel electrode (SCE, KCl saturated, Basi Model

RE-2BP) was used as reference electrode, and a graphite rod (6 mm dia. × 102 mm long, 99.9995%, Pine Instrument) was used as counter electrode. The electrolyte was purged with H<sub>2</sub> (ultrahigh grade purity) during the measurement. The applied potentials (E) reported in our work were referenced to the reversible hydrogen electrode (RHE) through standard calibration. In 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte (pH 0),  $E \text{ (vs. RHE)} = E \text{ (vs. SCE)} + E_{\text{SCE}} (= 0.241 \text{ V}) + 0.0592 \text{ pH} = E \text{ (vs. SCE)} + 0.241 \text{ V}$ . The overpotential ( $\eta$ ) was defined as  $E \text{ (vs. RHE)}$ . 4 mg sample was mixed with 1 mg carbon black (Vulcan XC-72) dispersed in Nafion (20  $\mu$ L) and isopropyl alcohol (0.98 mL). The catalyst materials (0.390 mg cm<sup>-2</sup>) were deposited on a glassy carbon (GC) rotating disk electrode (RDE, area = 0.1641 cm<sup>2</sup>, Pine Instrument), and a rotation speed of 1600 rpm was used for the linear sweep voltammetry (LSV) measurements. The Pt/C (20 wt.% Pt in Vulcan carbon black, Aldrich-Sigma) tested as reference sample using the same procedure.

Electrochemical impedance spectroscopy (EIS) measurements were carried out for the electrode in an electrolyte by applying an AC voltage of 10 mV in the frequency range of 100 kHz to 0.1 Hz at a bias voltage of -0.15 V (vs. RHE). To measure double-layer capacitance, a potential range in which no apparent Faradaic processes occur was determined from static via cyclic voltammograms (CV). This range is 0.1–0.2 V. All measured current in this non-Faradaic potential region is assumed to be due to double-layer capacitance. The charging current,  $i_c$ , is then measured from CVs at multiple scan rates. The working electrode was held at each potential vertex for 10 s before beginning the next sweep. The current density ( $J$ ) is equal to the product of the scan rate ( $v$ ) and the electrochemical double-layer capacitance ( $C_{dl}$ ), as given by equation  $J = v C_{dl}$ . Thus, a plot of  $J$  as a function of  $v$  yields a straight line with a slope equal to  $C_{dl}$ . The scan rates were 20–100 mV s<sup>-1</sup>.

**TOF Calculation.** The active site density and per-site turnover frequency (TOF) have been estimated as follows. It should be emphasized that since the nature of the active sites of the catalysts is not clearly understood yet and the real surface area for the nanostructured heterogeneous catalyst is hard to accurately determine, the following result is really just an estimation.

To estimate the electrochemically active surface site density, we used the  $C_{dl}$  value determined by CV measured at 0.1-0.2 V, in a non-Faradaic region, using various scan rates; 5.8, 22.5, and 10.2 mF cm<sup>-2</sup>, respectively, for 7%, 14%, and 20% CoCp<sub>2</sub>-intercalated WS<sub>2</sub>. The roughness factor, which is basically the surface area ratio between the catalyst vs. the metal electrodes (0.035 mF cm<sup>-2</sup>),<sup>S2</sup> is calculated as 165 (= 5.8 mF cm<sup>-2</sup>/0.035 mF cm<sup>-2</sup>), 643, and 291 for 7%, 14%, and 20% CoCp<sub>2</sub>-intercalated WS<sub>2</sub>, respectively.

The number of catalytic sites on the flat surface of catalyst can be estimated based on the crystal structure of distorted octahedral-phase 1T' WS<sub>2</sub>. Using the lattice parameters of 1T' phase WS<sub>2</sub> (we calculated as  $a = 3.27 \text{ \AA}$ ,  $b = 3.20 \text{ \AA}$ ,  $\gamma = 119^\circ$ ) and assuming one active site per WS<sub>2</sub> (which translates into one reactive sites per unit cell), the density of surface active sites is  $1/(0.5 \times 3.27 \times 3.20 \times \sin 119^\circ) \times 10^{16} \text{ cm}^{-2} = 2.2 \times 10^{15} \text{ atom cm}^{-2}$ .<sup>S3</sup> Our DFT calculation shows that the basal S sites above the center site (Co) of CoCp<sub>2</sub> are the most active sites, so the number of surface active sites is calculated as 0.14 (or 0.07, 0.20)  $\times 2.2 \times 10^{15} \text{ cm}^{-2}$ , for 14% (or 7%, 20%) CoCp<sub>2</sub>-intercalated WS<sub>2</sub>. The density of surface active sites ( $m$ ) of WS<sub>2</sub>-CoCp<sub>2</sub> (14%) on geometric area:  $0.14 \times 2.2 \times 10^{15} \text{ atom cm}^{-2} \times \text{roughness factor} (= 643) = 2.0 \times 10^{17} \text{ cm}^{-2}$ .

The total number of hydrogen (H<sub>2</sub>) gas turns overs was calculated from the current density (J in mA cm<sup>-2</sup>) according to  $n_{H_2} = J (\text{mA cm}^{-2})/1000 \text{ mA} \times 1 \text{ C s}^{-1} \times 1 \text{ mol e}^-/96486 \text{ C} \times (1 \text{ mol}$

$\text{H}_2/2 \text{ mol e}^{-1} \times (6.022 \times 10^{23} \text{ H}_2 \text{ molecules}/1 \text{ mol H}_2) = 3.12 \times 10^{15} \text{ H}_2 \text{ s}^{-1} \text{ cm}^{-2}$  per  $\text{mA cm}^{-2}$ .

For  $\text{WS}_2\text{-CoCp}_2$  (14%), the TOF for  $J = 30 \text{ mA cm}^{-2}$  at  $\eta = 0.2 \text{ V}$  is  $n_{\text{H}_2}/m = 30 \times 3.12 \times 10^{15} \text{ H}_2 \text{ s}^{-1} \text{ cm}^{-2} / 2.0 \times 10^{17} \text{ cm}^{-2} = 0.47 \text{ H}_2 \text{ s}^{-1}$ . We summarized the TOF values at 0.2 V as follows.

Concentration of $\text{CoCp}_2$	$J$ ( $\text{mA cm}^{-2}$ ) at $\eta = 0.2 \text{ V}$	$n_{\text{H}_2}$ ( $\text{s}^{-1} \text{ cm}^{-2}$ )	Roughness factor	$m$ ( $\text{cm}^{-2}$ )	TOF ( $\text{s}^{-1}$ )
7%	3.2	$1.0 \times 10^{16}$	165	$2.5 \times 10^{16}$	0.40
14%	30	$9.4 \times 10^{16}$	643	$2.0 \times 10^{17}$	0.47
20%	18	$5.8 \times 10^{16}$	291	$1.3 \times 10^{17}$	0.45

**Computations.** Geometry optimizations were performed using the Vienna ab-initio simulation package (VASP).<sup>S4</sup> The electron-ion interactions were described using the projector-augmented wave (PAW) method.<sup>S5</sup> Attractive van der Waals interactions were included using Grimme's correction for the Perdew–Burke–Ernzerhof (PBE)-D3 method.<sup>S6</sup> For structural optimization, the atoms were relaxed in the direction of the Hellmann-Feynman force using the conjugate gradient method with an energy cut-off of 520 eV, until a stringent convergence criterion of 1 meV/Å was satisfied. Lattice constants were optimized using the PBE-D3 exchange-correlation functional.<sup>S7</sup> The  $k$ -point sampling was performed using  $\Gamma$ -centered  $6 \times 6 \times 4$  points, which gave the total energy of the  $(4 \times 4)$   $\text{WS}_2$  complex in the 1T' phase within 1 meV.

The adsorption energies and activation barriers for an  $\text{H}^+$  ion at different adsorption sites were calculated using the climbing image-nudged elastic band (CINEB) method.<sup>S8</sup> For simplicity, we employed a slab geometry in which a supercell included two  $\text{WS}_2$  layers and one  $\text{CoCp}_2$  molecule. The coordinate system is defined so that the  $\text{WS}_2$  layer lies on the  $ab$  plane parallel to the  $XY$  plane. Two lowest sublayers (i.e., all W ions as well as one-half of S



atoms on the lower WS<sub>2</sub> layer) were fixed, while all other atoms were permitted to relax freely. For computational efficiency, we adopted a supercell of (4×4) 2WS<sub>2</sub> with one CoCp<sub>2</sub> molecule. For the Volmer reaction, the *c* constant of the supercell perpendicular to the slab geometry was maintained at a sufficiently large value (31.08 Å). In the initial configuration (IC), a H<sup>+</sup> ion was placed far from the WS<sub>2</sub> surface so that its closest distances to two WS<sub>2</sub> image surfaces along the *c* direction were at least 6.10 and 7.05 Å, respectively. In Heyrovsky reaction, the same *c* constant was used to identify the activated complex (AC). In the next step, the supercell was further elongated along the *c* direction by 19.00 Å and the second H<sup>+</sup> ion was placed and separated from the two image surfaces by 9.00 and 10.00 Å in the IC of the reaction, so that the activation barrier can be estimated more accurately by improving the accuracy of the calculated energies of the IC and AC.

## II. Supporting Tables

**Table S1.** Characteristics of WS<sub>2</sub> samples and their HER performance.

I <sup>a</sup>	[CoCp <sub>2</sub> ]/[(NH <sub>4</sub> ) <sub>2</sub> WS <sub>4</sub> ]	[CoCp <sub>2</sub> ] <sup>b</sup>	[Co <sup>3+</sup> ] <sup>c</sup>	[Co <sup>2+</sup> ] <sup>d</sup>	η <sub>J=10</sub> <sup>e</sup>	b <sup>f</sup>
None	0	0	0	0	0.36	120
CoCp <sub>2</sub>	0.5	7	4.4	2.6	0.24	72
CoCp <sub>2</sub>	1	14	8.2	5.8	0.17	40
CoCp <sub>2</sub>	2	20	9.2	10.8	0.19	47

<sup>a</sup> Intercalated molecules; <sup>b</sup> Concentration (%) of intercalated CoCp<sub>2</sub> ([CoCp<sub>2</sub>]/[WS<sub>2</sub>]) determined using EDX, XPS (Co 2p and W 4f peaks), and ICP-AES; <sup>c</sup> Concentration (%) of intercalated cation form ([Co<sup>3+</sup>]/[WS<sub>2</sub>]); <sup>d</sup> Concentration (%) of Co<sup>2+</sup> form = [CoCp<sub>2</sub>] - [Co<sup>3+</sup>]; <sup>e</sup> HER overpotential (V) at  $J = 10 \text{ mA cm}^{-2}$ ; <sup>f</sup> Tafel slope (mV dec<sup>-1</sup>) for HER.

**Table S2.** Fitting parameters of EXAFS data for the WS<sub>2</sub> samples (see **Figure 2** and **Figure S5**). The FT curves of EXAFS were fitted to two scattering shells. A least-squares curve parameter method was performed using the ARTEMIS module of IFEFFIT and USTCXAFFS software packages.

	Sample	Scattering Path	$R$ (Å) <sup>a</sup>	CN <sup>b</sup>	$\Delta E$ (eV) <sup>c</sup>	$\sigma^2$ (Å <sup>2</sup> ) <sup>d</sup>
W L <sub>3</sub> -edge	2H WS <sub>2</sub>	W-S	2.40540	6.0	6.479	0.00193
		W-W	3.18070	6.0	4.811	0.00320
	WS <sub>2</sub>	W-S	2.39880	4.1 ± 0.7	5.398	0.00407
		W-W	2.76410	2.3 ± 1.2	4.036	0.00539
	WS <sub>2</sub> -CoCp <sub>2</sub> -7%	W-S	2.39810	3.3 ± 1.5	5.842	0.00671
		W-W	2.76152	3.1 ± 1.9	3.716	0.00581
	WS <sub>2</sub> -CoCp <sub>2</sub> -14%	W-S	2.39880	3.2 ± 1.0	5.328	0.00149
		W-W	2.76410	2.1 ± 2.1	-6.499	0.00490
WS <sub>2</sub> -CoCp <sub>2</sub> -20%	W-S	2.40943	3.0 ± 2.0	4.830	0.00886	
	W-W	2.75942	2.2 ± 2.0	3.518	0.00681	
	CoCp <sub>2</sub>	Co-C	2.04	10.4 ± 2.4	6.422	0.00508
Co K-	WS <sub>2</sub> -CoCp <sub>2</sub> -7%	Co-C	2.08	10.2 ± 2.1	5.962	0.00517

WS <sub>2</sub> -CoCp <sub>2</sub> -14%	Co-C	2.10	10.3 ± 2.2	4.817	0.00221
WS <sub>2</sub> -CoCp <sub>2</sub> -20%	Co-C	2.14	10.5 ± 2.3	4.114	0.01234

<sup>a</sup> Distance between scattering atoms. The FT curve of the 2H phase WS<sub>2</sub> are characterized by two main peaks at 2.40 and 3.18 Å, corresponding to the nearest W–S and W–W bonds, respectively. In contrast, in the FT curves of WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub>, the distance of W–W bond is decreased to 2.76 Å.

<sup>b</sup> Coordination number of W atoms. Since the coordination number of W–W decreased from 6 (2H phase) to 2 (1T' phase), the intensity of 1T' phase WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub> is nearly reduced by 1/3. This result indicates that the WS<sub>2</sub> adopts a distorted octahedral coordination.

<sup>c</sup> Edge energy shift, representing between the energy grids of experimental and theoretical data.

<sup>d</sup> Debye-Waller factor, which measures the static and thermal disorder, is larger for the intercalated samples than 2H-WS<sub>2</sub>. It suggests that the intercalation produces a broad range of W-S and W-W distances.

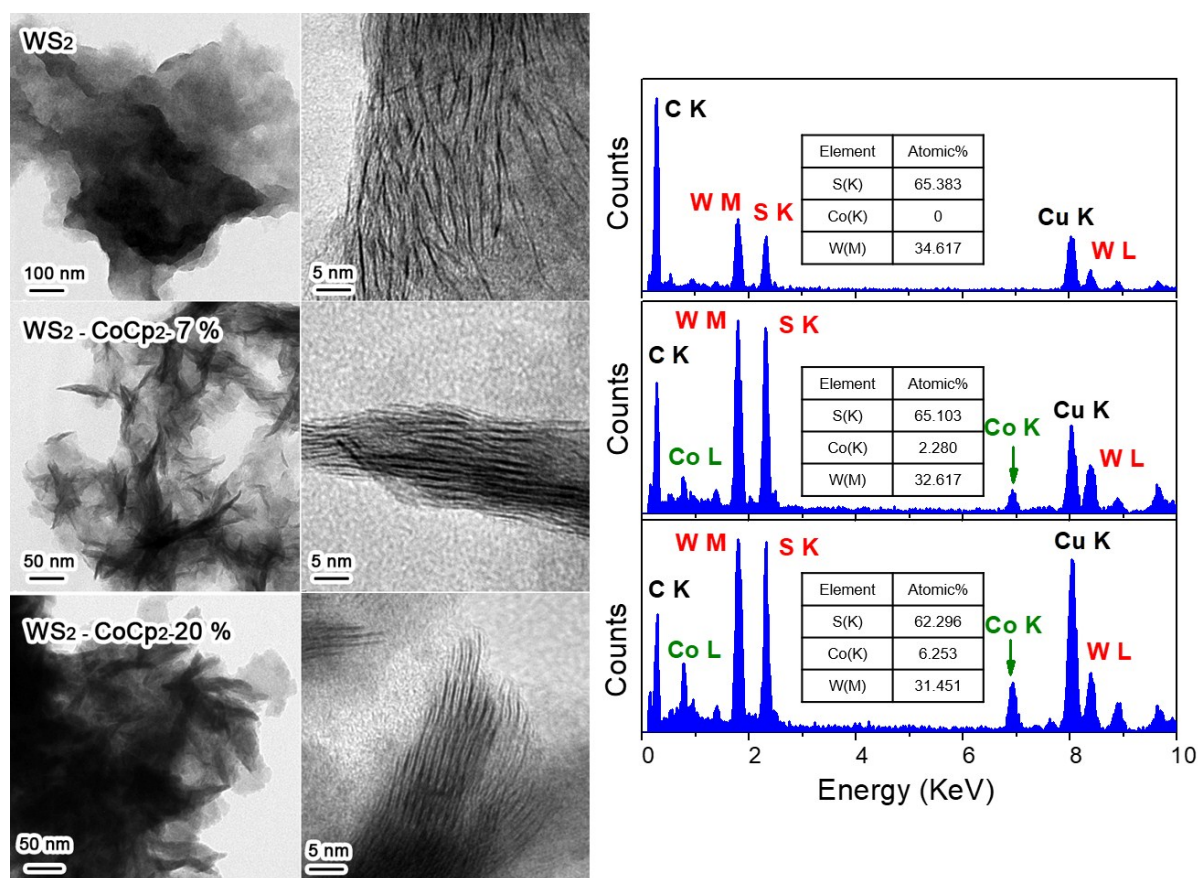
**Table S3.** Impedance parameters for the equivalent circuit that was shown in **Figure S8**, and the double-layer capacitance ( $C_{dl}$ ) as shown in **Figure S9**.

Samples	EIS		$C_{dl}$ (mF cm <sup>-2</sup> )
	$R_s$ ( $\Omega$ )	$R_{ct}$ ( $\Omega$ )	
WS <sub>2</sub>	5.0	360	2.2
WS <sub>2</sub> -CoCp <sub>2</sub> -7%	5.2	69.1	5.8
WS <sub>2</sub> -CoCp <sub>2</sub> -14%	4.0	34.2	22.5
WS <sub>2</sub> -CoCp <sub>2</sub> -20%	4.5	42.3	10.2

**Table S4.** Comparison of HER performance (in pH 0) of WS<sub>2</sub> in the literatures (N/A: Not applicable).

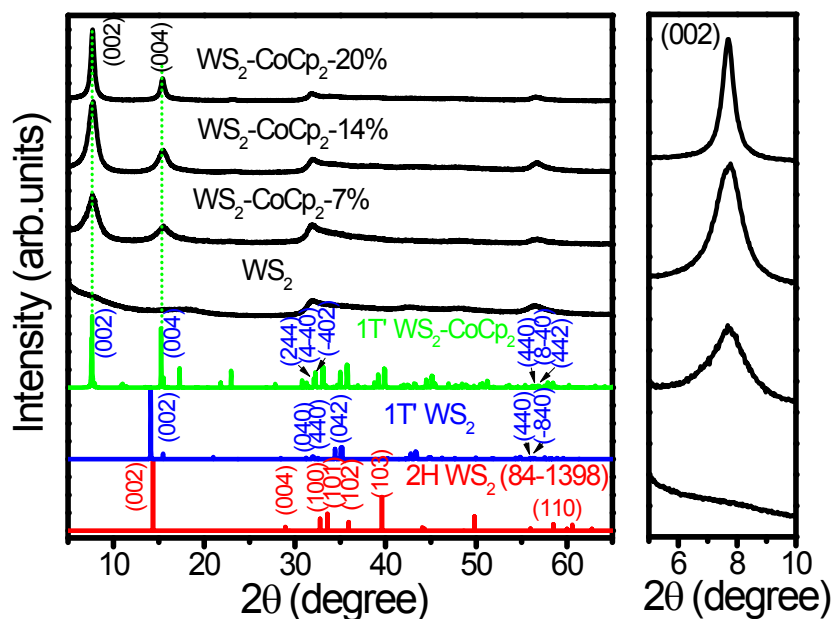
Reference	Materials	Phase	$E_{J=10}$ (mV) at 10 mA cm <sup>-2</sup>	Tafel slope (mV dec <sup>-1</sup> )	TOF
S9	Exfoliated WS <sub>2</sub>	1T	N/A (~280)	60	175 s <sup>-1</sup> at 288 mV
S10	WS <sub>2</sub> nanoribbon	2H	225	68	N/A
S11	WS <sub>2</sub> Nanoflake	2H	N/A (~420)	48	N/A
S12	WS <sub>2</sub> @heteroatom doped graphene	1T	125	52.7	N/A
S13	WS <sub>2</sub> /Graphene	2H	229	73	N/A
S14	N doped WS <sub>2</sub>	2H	N/A (~140)	69	0.725 s <sup>-1</sup> at 108 mV
S15	CoS <sub>2</sub> @WS <sub>2</sub> /CC	2H	97.2	66	N/A
S16	WS <sub>2</sub> @graphene	1T	260	45.4	N/A
S17	Colloidal WS <sub>2</sub>	1T'	200	50	N/A
S18	Co:WS <sub>2</sub> @Co:W <sub>18</sub> O <sub>49</sub>	1T	240	49	N/A
Present work	WS <sub>2</sub> -CoCp <sub>2</sub>	1T'	170	40	0.47 s <sup>-1</sup> at 200 mV

### III. Supporting Figures



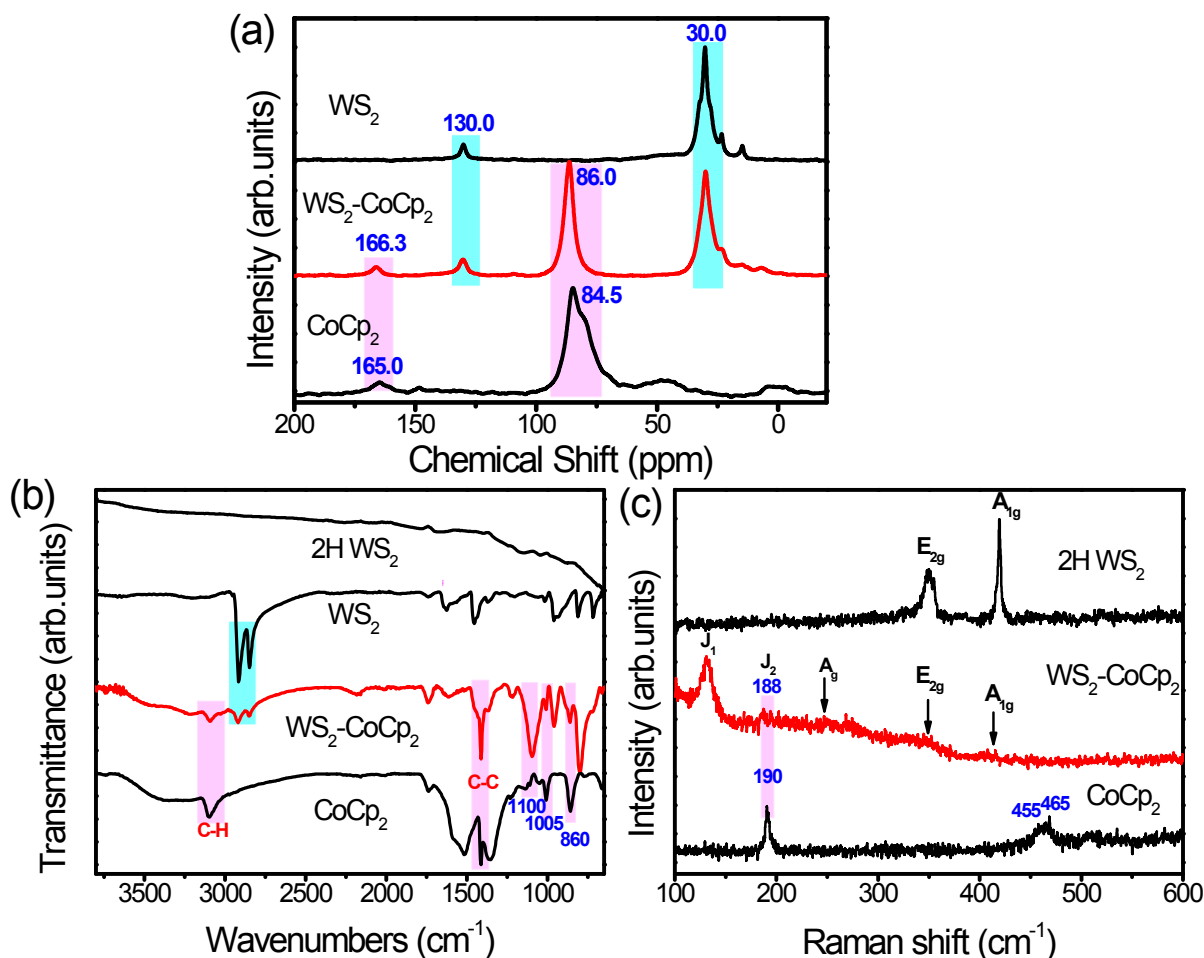
**Figure S1.** HRTEM images of WS<sub>2</sub>, WS<sub>2</sub>-CoCp<sub>2</sub>-7%, and WS<sub>2</sub>-CoCp<sub>2</sub>-20% and their EDX spectrum showing that the atomic ratio of [Co]/[W] is 0%, 7%, and 20%, respectively.

The WS<sub>2</sub> nanosheets consisted of the aggregated thin layers. The distance between adjacent WS<sub>2</sub> layers ( $d_{002}$ ) is in the wide range of 0.7-3 nm. The WS<sub>2</sub>-CoCp<sub>2</sub> complexes exhibit the similar morphology in which the average size of nanosheets is 80 nm, and the average thickness is 10 nm. The atomically resolved TEM images show that the interlayer distance ( $d_{002}$ ) is uniformly expanded to 11 Å. The EDX spectrum revealed the homogeneous distribution of the composition in the entire sample. The peaks of Co (K shell) and W (M shell) provides the ratio of CoCp<sub>2</sub> and WS<sub>2</sub>.



**Figure S2.** XRD pattern of  $\text{WS}_2$  and  $\text{WS}_2\text{-CoCp}_2$  complexes with 7%, 14%, and 20%. The peaks were referenced to those of 2H phase  $\text{WS}_2$  (JCPDS Card 84-1398;  $a = 3.153 \text{ \AA}$  and  $c = 12.323 \text{ \AA}$ ), 1T' phase  $\text{WS}_2$  ( $a = 3.26 \text{ \AA}$ ,  $b = 3.19 \text{ \AA}$ , and  $c = 12.02 \text{ \AA}$ ), and 1T' phase  $\text{WS}_2\text{-CoCp}_2$  ( $a = 3.26 \text{ \AA}$ ,  $b = 3.19 \text{ \AA}$ , and  $c = 23.2 \text{ \AA}$ ), as shown at the bottom.

VESTA program (<http://jp-minerals.org/vesta/en/>) was used to generate the simulated XRD pattern for 1T' phase whose lattice parameters were obtained by the present calculation. The (002) and (004) peak position of  $\text{WS}_2\text{-CoCp}_2$  complexes is matched well with that of the simulated XRD pattern of  $\text{WS}_2\text{-CoCp}_2$ . The broad peaks at  $2\theta = 32^\circ$  and  $57^\circ$  are ascribed to the overlapped many peaks. For non-intercalated  $\text{WS}_2$ , the (002) and (004) peaks exhibit almost zero intensity, probably due to the wide range of interlayer distance as identified by the TEM data.



**Figure S3.** (a) Solid state  $^{13}\text{C}$  NMR spectra of  $\text{CoCp}_2$ ,  $\text{WS}_2$ , and  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$ . (b) IR spectrum of  $\text{CoCp}_2$ ,  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$ ,  $\text{WS}_2$ , and  $2\text{H}$  phase  $\text{WS}_2$  ( $2\text{H-WS}_2$ ) powders. (c) Raman spectra of  $\text{CoCp}_2$ ,  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$ , and  $2\text{H-WS}_2$ . The  $\text{CoCp}_2$  and  $2\text{H-WS}_2$  powders were purchased from Sigma-Aldrich.

(a) The  $^{13}\text{C}$  NMR spectrum of  $\text{CoCp}_2$  powders (purchased from Sigma-Aldrich) show the C-C(=C) peak of cyclopentadienyl ring ( $\text{Cp}_2$ ) of  $\text{CoCp}_2$  at 84.5 and that of cationic form at 165 ppm. The  $\text{WS}_2\text{-CoCp}_2$  shows those peaks at 86.0 and 166.3 ppm. The blue shift from the peaks of  $\text{CoCp}_2$  could be due to the charge transfer from the Cp to  $\text{WS}_2$ . The peaks at 30.8 and 130.0 ppm can be assigned to the C-C and C=C of oleylamine, respectively, since the  $\text{WS}_2$  shows the same peak.

(b) We assigned the IR peaks of  $\text{CoCp}_2$  based on the references;<sup>S19,S20</sup> C-H stretching at 3090  $\text{cm}^{-1}$ , C-C stretching at 1415  $\text{cm}^{-1}$ , C-C breathing at 1100  $\text{cm}^{-1}$ , C-C-H (in-plane) bending at

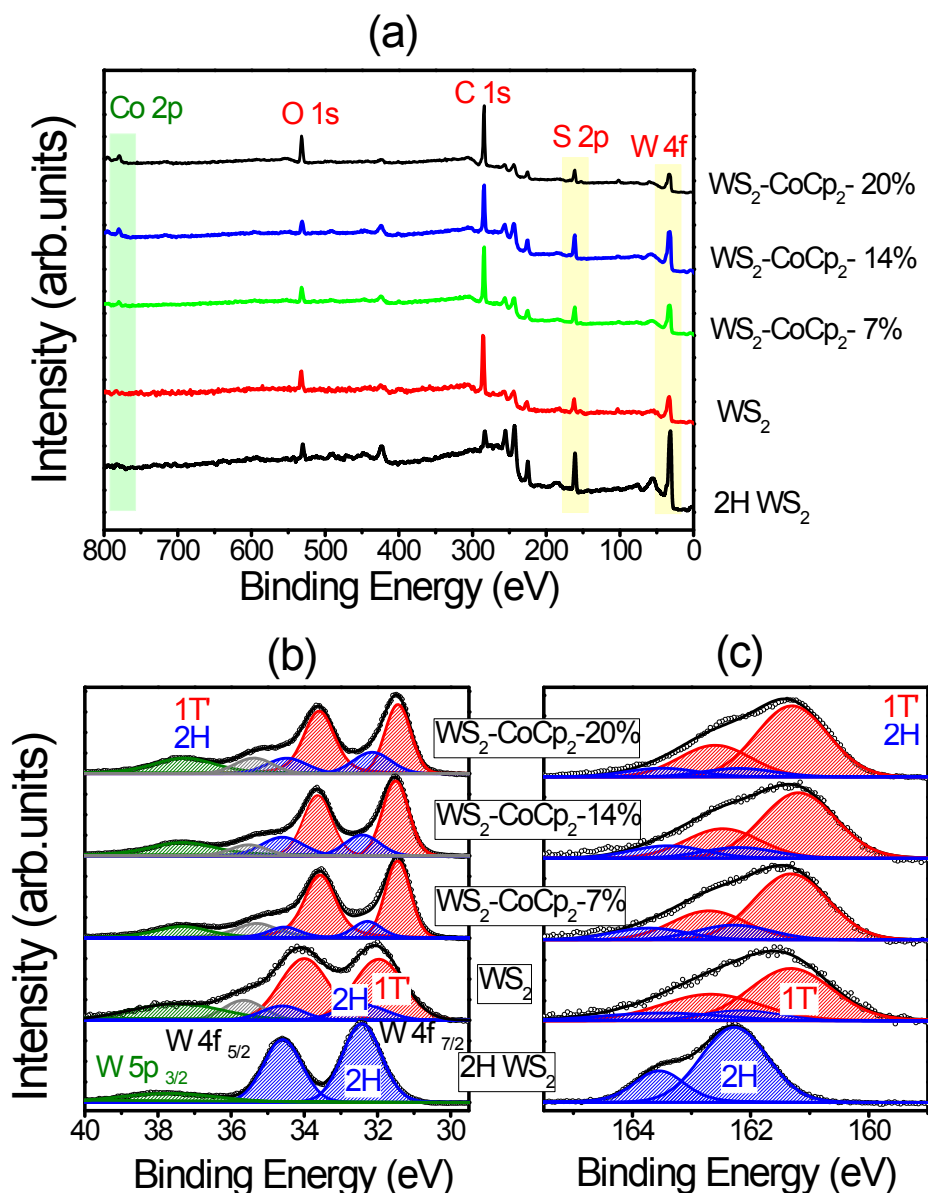
1005 $\text{cm}^{-1}$ , and C-H (out-of-plane) bending at 860  $\text{cm}^{-1}$ . The peak at 1500  $\text{cm}^{-1}$  is originated from the scissoring mode of water. The  $\text{WS}_2\text{-CoCp}_2$  shows the same peaks; C-H at 3088  $\text{cm}^{-1}$ , C-C at 1415  $\text{cm}^{-1}$ , C-C at 1095  $\text{cm}^{-1}$ , C-C-H at 1008  $\text{cm}^{-1}$ , and C-H bending at 860  $\text{cm}^{-1}$ . It is known that the last peak is sensitive to oxidation; the peak position of  $\text{CoCp}_2$  and  $\text{WS}_2\text{-CoCp}_2$  are close to those of cation form.<sup>S19</sup> The  $\text{WS}_2$  shows the strong C-H peaks at 2915 and 2850  $\text{cm}^{-1}$  are originated from the oleylamine. The  $\text{WS}_2\text{-CoCp}_2$  also shows the same peaks.

### Summary of IR peak assignment

Samples	C-H stretching	C-C stretching	C-C breathing	C-C-H (in-plane) bending	C-H (out-of-plane) bending
$\text{CoCp}_2$	3090	1415	1100	1005	860
$\text{WS}_2\text{-CoCp}_2$	3088	1415	1095	1008	860

(c) The 2H phase  $\text{WS}_2$  exhibit two characteristic Raman peaks at 350 and 418  $\text{cm}^{-1}$ , corresponding to the in-plane  $E_{2g}^1$  and out-of-plane  $A_{1g}$  vibration modes, respectively. The  $\text{WS}_2\text{-CoCp}_2$  shows the peaks of 1T' phase:  $J_1$  peak at 133  $\text{cm}^{-1}$ ,  $J_2$  peak at 188  $\text{cm}^{-1}$ , and  $A_g$  at 258  $\text{cm}^{-1}$ , with the  $E_{2g}^1$  mode at 350  $\text{cm}^{-1}$  and the  $A_{1g}$  mode at 410  $\text{cm}^{-1}$ .<sup>S17</sup> The  $\text{CoCp}_2$  shows the Raman peaks at 190, 455, and 465  $\text{cm}^{-1}$ , which can be assigned to the Cp ring deformation vibrational mode of  $\text{CoCp}_2^+$ , the Co-Cp bond stretching vibrational mode of  $\text{CoCp}_2^+$ , and ring tilt vibrational mode of  $\text{CoCp}_2$ , respectively.<sup>S21</sup> The peak of  $\text{WS}_2\text{-CoCp}_2$  at 188  $\text{cm}^{-1}$  could be also originated from the Cp ring deformation vibrational mode.





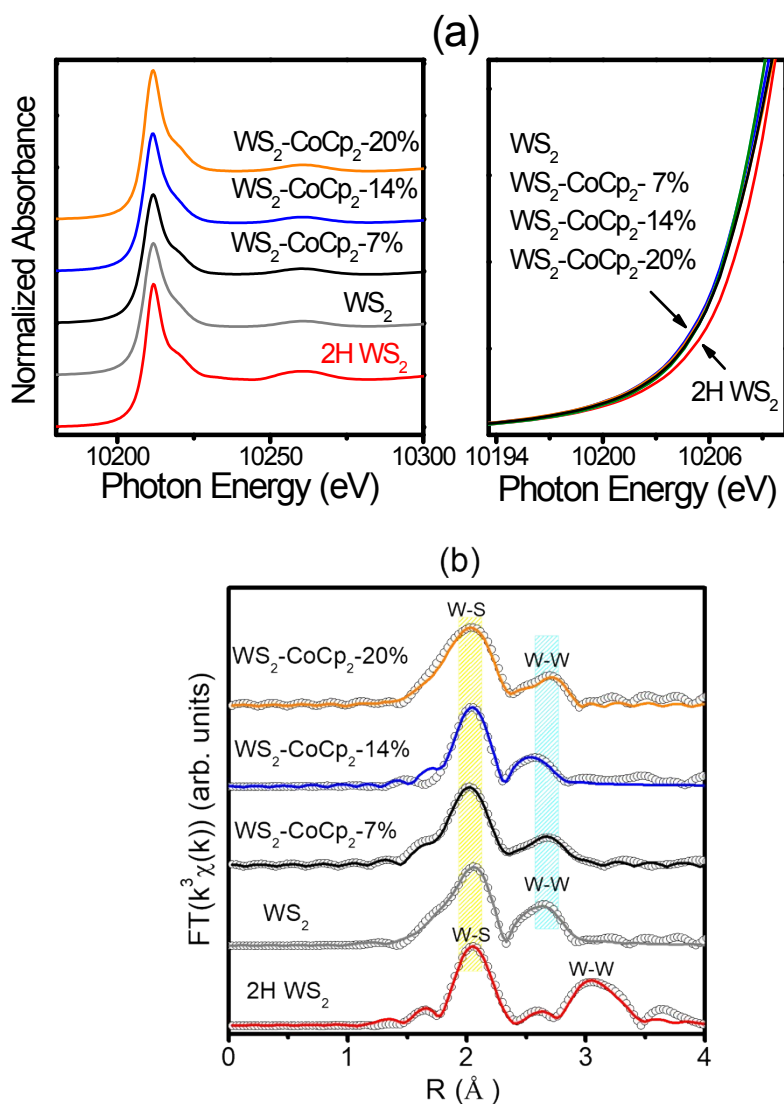
**Figure S4.** (a) Full-range XPS data of 2H-WS<sub>2</sub> (purchased from Sigma-Aldrich, 99.99%), WS<sub>2</sub>, and WS<sub>2</sub>-CoCp<sub>2</sub> (7%, 14%, and 20%). Fine-scanned (b) W 4f and W 5p<sub>3/2</sub> peaks and (c) S 2p peaks.

(a) The atomic ratio of [S]/[W] is 2 for all samples. For three WS<sub>2</sub>-CoCp<sub>2</sub> samples, [Co]/[W] is 7, 14, and 20%. The C peak of WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub> is larger than that of 2H WS<sub>2</sub>, due to the residual solvent and/or intercalated CoCp<sub>2</sub>.

(b) The W 4f<sub>7/2</sub> and 4f<sub>5/2</sub> peaks (separated by about 2.18 eV), and W 5p<sub>3/2</sub> peak (green). The 2H-WS<sub>2</sub> shows the 4f<sub>7/2</sub> peak at 32.4 eV, which is 1 eV blue-shifted from the neutral W at 31.4 eV. The 4f<sub>7/2</sub> peak of WS<sub>2</sub>-CoCp<sub>2</sub> appears at 31.5 eV; this negligible shift is due to the metallicity

of 1T' phase. The  $4f_{7/2}$  peak was resolved into two bands: 1T' phase (red) at 31.5 eV and 2H phase (blue) at 32.4 eV. The fraction of the 1T' phase was determined as avg. 80%, indicating that this is the major phase. The small peak (grey) at 35.6 eV is assigned to the defective bonding with oxygen.

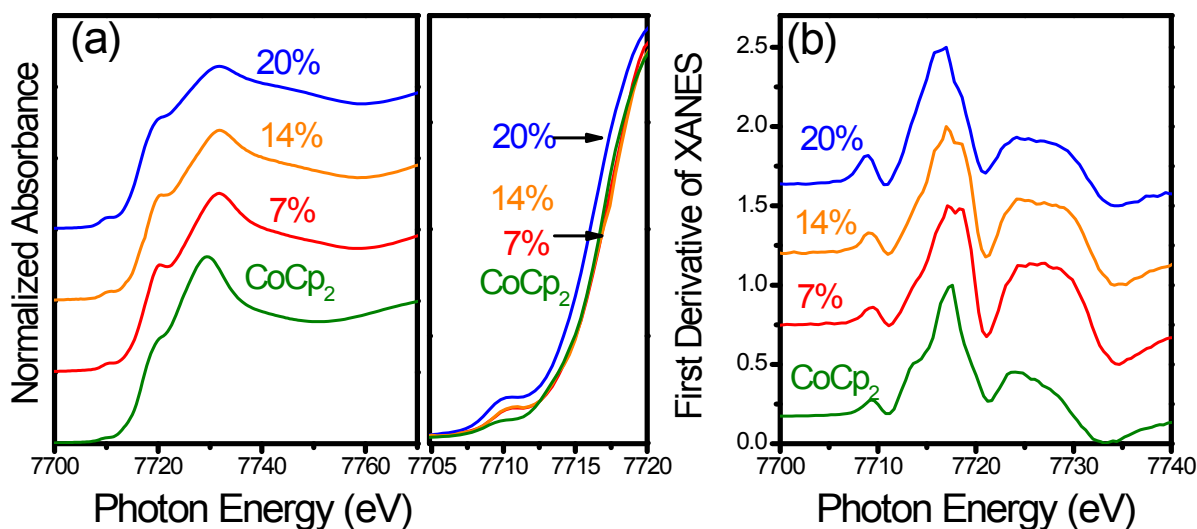
(c) The S  $2p_{3/2}$  and S  $2p_{1/2}$  peaks, which are separated by about 1.2 eV. The 2H-WS<sub>2</sub> shows peaks at 162.3 and 163.5 eV, which are 1.7 eV red-shifted with respect to the signal of neutral S (S<sup>0</sup>) at 164.0 and 165.2 eV. They correspond to the S<sup>2-</sup> anions bonded with the W cations in the 2H phase. For the WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub> samples, the broad peak was resolved into four bands; two each for the 2H phase (blue) and the 1T' phase (red). The larger red-shift band, S  $2p_{3/2}$  at 161.3 eV and S  $2p_{1/2}$  at 162.5 eV, are assigned to those of electron-rich 1T' phase.



**Figure S5.** XANES spectra at the (a) W L<sub>3</sub> edge for 2H-WS<sub>2</sub> (purchased from Sigma-Aldrich, 99.99%), WS<sub>2</sub>, and WS<sub>2</sub>-CoCp<sub>2</sub> (7%, 14%, and 20%) and (b) non-phase-corrected k<sup>3</sup> weighted FT EXAFS data (open circles) of WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub> (7%, 14%, and 20%) at the W L<sub>3</sub> edge and their fitting curves (line).

(a) The evolution of the local crystal structure of samples was probed with W L<sub>3</sub> edge X-ray absorption near edge spectra (XANES) analysis. Magnified scaled spectrum reveals that edge energies of all samples are lower than those of 2H-WS<sub>2</sub>. It suggests that the intercalation of CoCp<sub>2</sub> induces more metallic states of W.

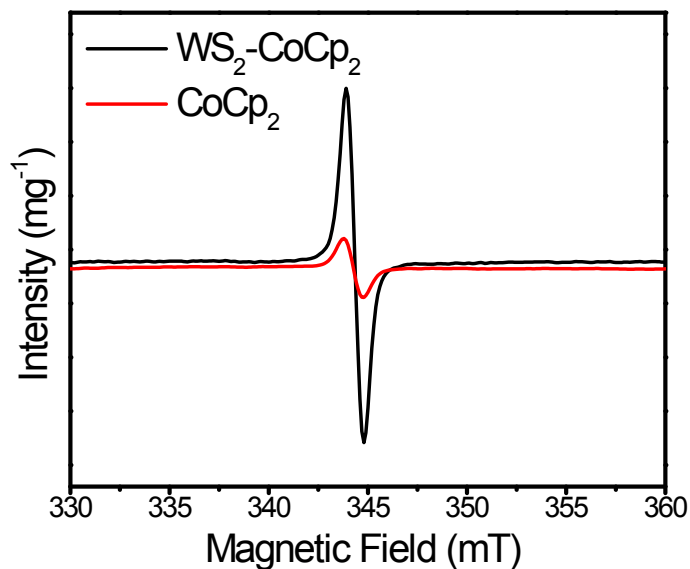
(c) In the FT profiles (in real space), the EXAFS peaks indicate the distances to nearest neighbor atoms. The profiles of WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub> are significantly different from that of 2H-WS<sub>2</sub>, suggesting a remarkable change in the local atomic arrangements due to the phase change to 1T'. The FT curves were fitted to two scattering shells, W-S and W-W bonds (see the parameters in **Table S2**). The shortest distance between W-S ( $d_{W-S}$ ) and W-W ( $d_{W-W}$ ) is 2.40 Å and 3.19 Å for the 2H-WS<sub>2</sub>. All of WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub> shows  $d_{W-S} = 2.40$  Å and  $d_{W-W} = 2.76$  Å, consistently with the values of 1T' phase.



**Figure S6.** (a) XANES spectra of CoCp<sub>2</sub> powders and WS<sub>2</sub>-CoCp<sub>2</sub> (7, 14, and 20%) at the Co K edge. (b) First derivative of absorbance curve in the onset region.

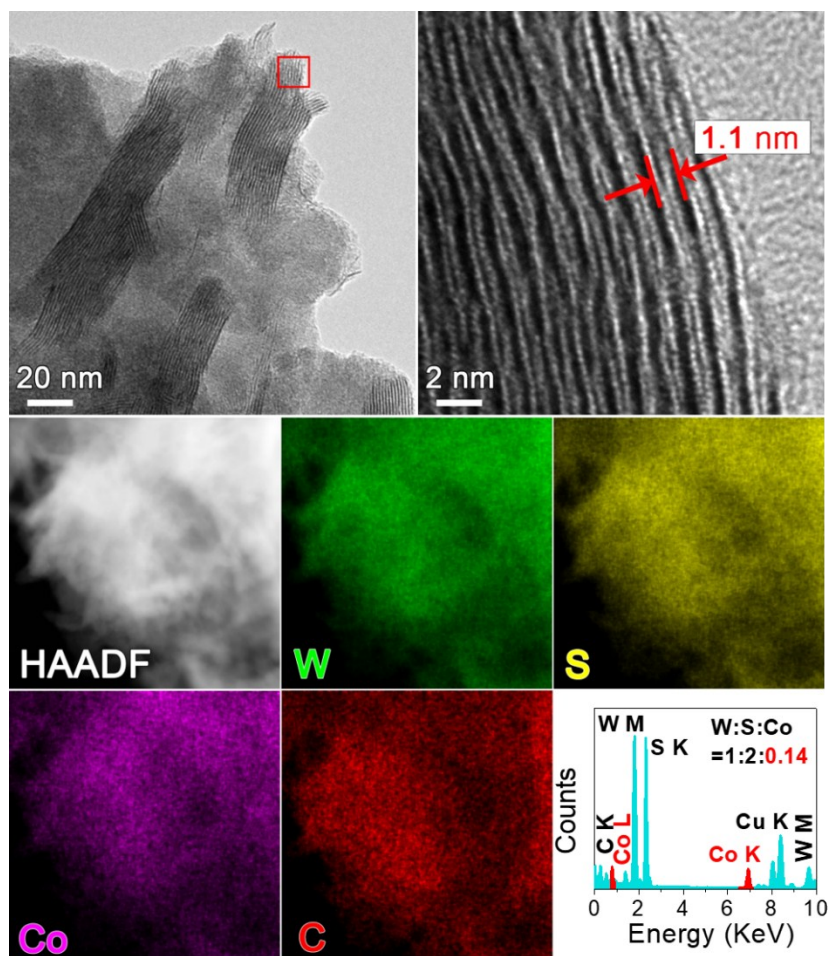
(a) The evolution of the local crystal structure of Co-containing samples was probed with the Co K edge XANES. The absorption edge in magnified scale shows a red shift for WS<sub>2</sub>-CoCp<sub>2</sub>-20%.

(b) The edge position was deduced from the first derivative of absorption curve in the onset region. The position of CoCp<sub>2</sub> powders is 7717.40 eV, which is blue shifted from than that of complexes, probably due to the higher concentration of cation (CoCp<sub>2</sub><sup>+</sup>). The CoCp<sub>2</sub> powder underwent the oxidation and exist as the cationic complexes with hydroxide. It should be noted that the data of CoCp<sub>2</sub> powder is just used as the reference to assign the peak. As the concentration of CoCp<sub>2</sub> increases, the position shift to the lower energy; 7717.28, 7717.25, and 7716.79 eV for 7%, 14%, and 20%, respectively, indicating that the oxidation number decreases due to the increased concentration of neutral form, which is consistent with the EXAFS and XPS data.



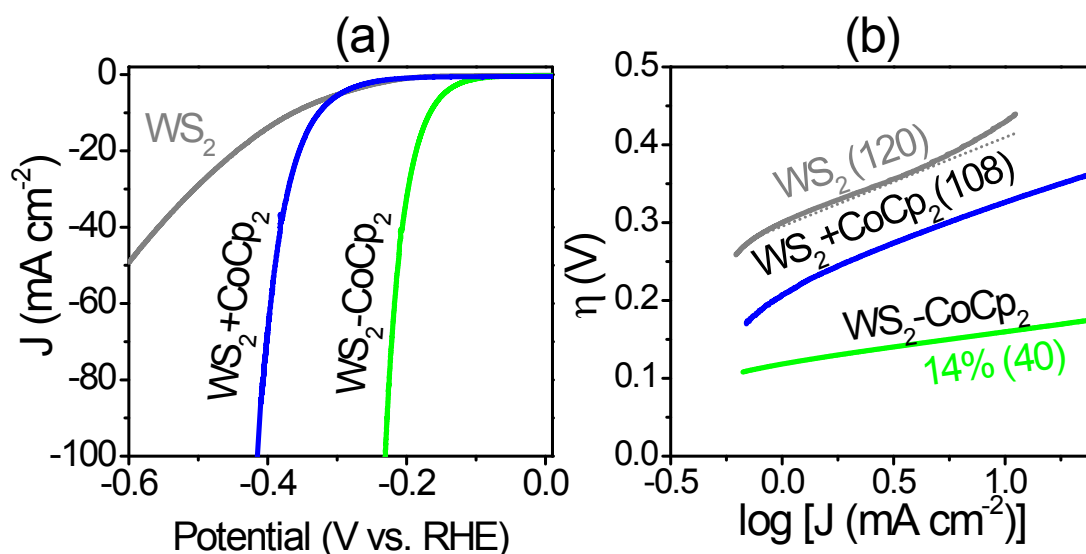
**Figure S7.** Electron paramagnetic resonance (EPR) spectra for  $\text{CoCp}_2$  and  $\text{WS}_2\text{-CoCp}_2$  (14%).

EPR measurements were performed for  $\text{CoCp}_2$  and  $\text{WS}_2\text{-CoCp}_2$  at room temperature. Both  $\text{CoCp}_2$  precursor and  $\text{WS}_2\text{-CoCp}_2$  exhibit a single peak, centered at 344 mT ( $g = 2.00$ ). The signal can be originated from the low spin of  $\text{Co}^{2+}$  state in  $\text{CoCp}_2$  neutral molecule.<sup>S22</sup> The larger intensity of  $\text{WS}_2\text{-CoCp}_2$  than that of  $\text{CoCp}_2$  is ascribed to the higher population of neutral form, consistently with the XPS, EXAFS, IR, and Raman data.



**Figure S8.** TEM images and EDX data of WS<sub>2</sub>-CoCp<sub>2</sub>-14% samples after 12h chronoamperometric measurement.

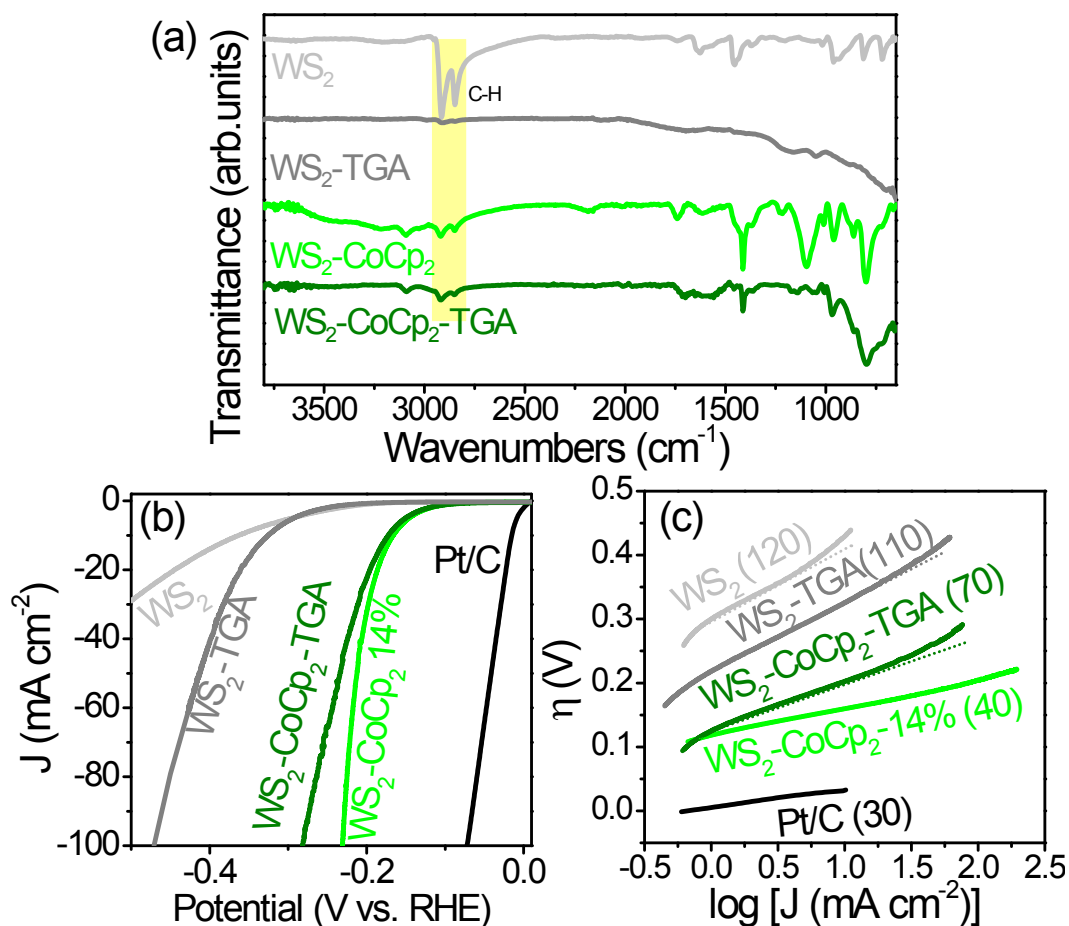
The average size (80 nm) and thickness (20 nm) of nanosheets are similar with those of the samples before HER. The lattice-resolved TEM image shows that the average distance between adjacent WS<sub>2</sub> layers ( $d_{002}$ ) is the same as that of the before samples; 1.1 nm. The EDX mapping and spectrum show that the Co atoms distribute homogeneously over the entire samples with 14%, which remain the same after the HER.



**Figure**

**e S9.** (a) LSV curves (scan rate: 2 mV s<sup>-1</sup>) for a physical mixture of WS<sub>2</sub> and 14% CoCp<sub>2</sub> molecules (referred to as WS<sub>2</sub>+CoCp<sub>2</sub>) in H<sub>2</sub>-saturated 0.5 M H<sub>2</sub>SO<sub>4</sub>. The data of WS<sub>2</sub> and WS<sub>2</sub>-CoCp<sub>2</sub>-14% is plotted for comparison. The overpotential for a current density of 10 mA cm<sup>-2</sup> is 0.32 V, which is close to 0.36 V of WS<sub>2</sub>. (b) Tafel plot derived from the LSV curve, where the Tafel slope is indicated in parentheses. The Tafel slope is 108 mV dec<sup>-1</sup>, which is close to 120 mV dec<sup>-1</sup> of WS<sub>2</sub>.



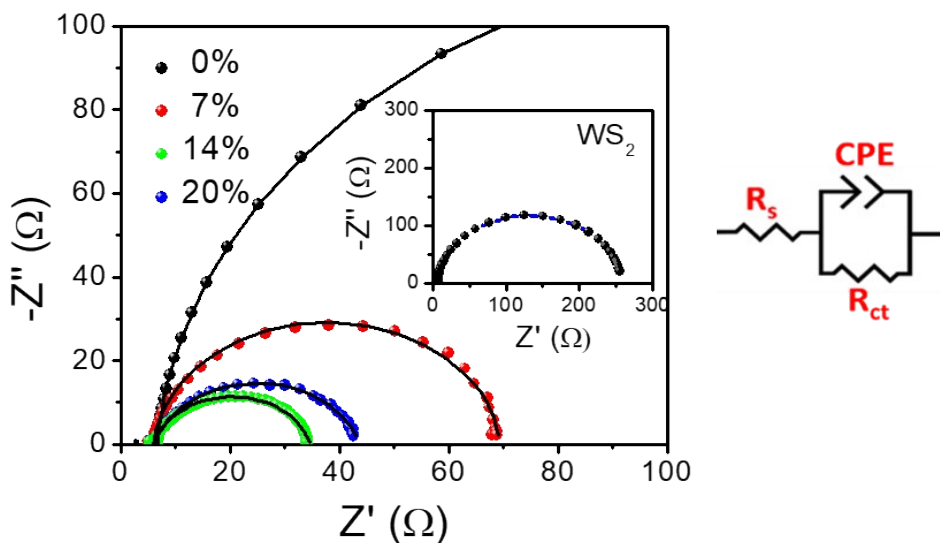


**Figure S10.** (a) FTIR spectrum of  $\text{WS}_2$  and  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$  before and after treatment with thioglycolic acid (TGA). The after samples were referred to as  $\text{WS}_2\text{-TGA}$  and  $\text{WS}_2\text{-CoCp}_2\text{-14\%-TGA}$ . (b) LSV curves (scan rate:  $2 \text{ mV s}^{-1}$ ) for  $\text{WS}_2\text{-TGA}$  and  $\text{WS}_2\text{-CoCp}_2\text{-14\%-TGA}$  in  $\text{H}_2$ -saturated  $0.5 \text{ M H}_2\text{SO}_4$ . The data of the before samples is plotted for comparison. (c) Tafel plot derived from the LSV curve, where the Tafel slope is indicated in parentheses.

(a) The as-synthesized powders ( $30 \text{ mg}$ ) were dispersed in anhydrous ethanol  $50 \text{ mL}$  and TGA (Sigma-Aldrich,  $99\%$ )  $3 \text{ mL}$  was added in the solution.<sup>S23</sup> The mixture was vigorously stirred at room temperature for  $12\text{h}$  under argon gas flow, and then washed with ethanol by several times. The intensity of oleylamine peak is significantly reduced after the TGA treatments.

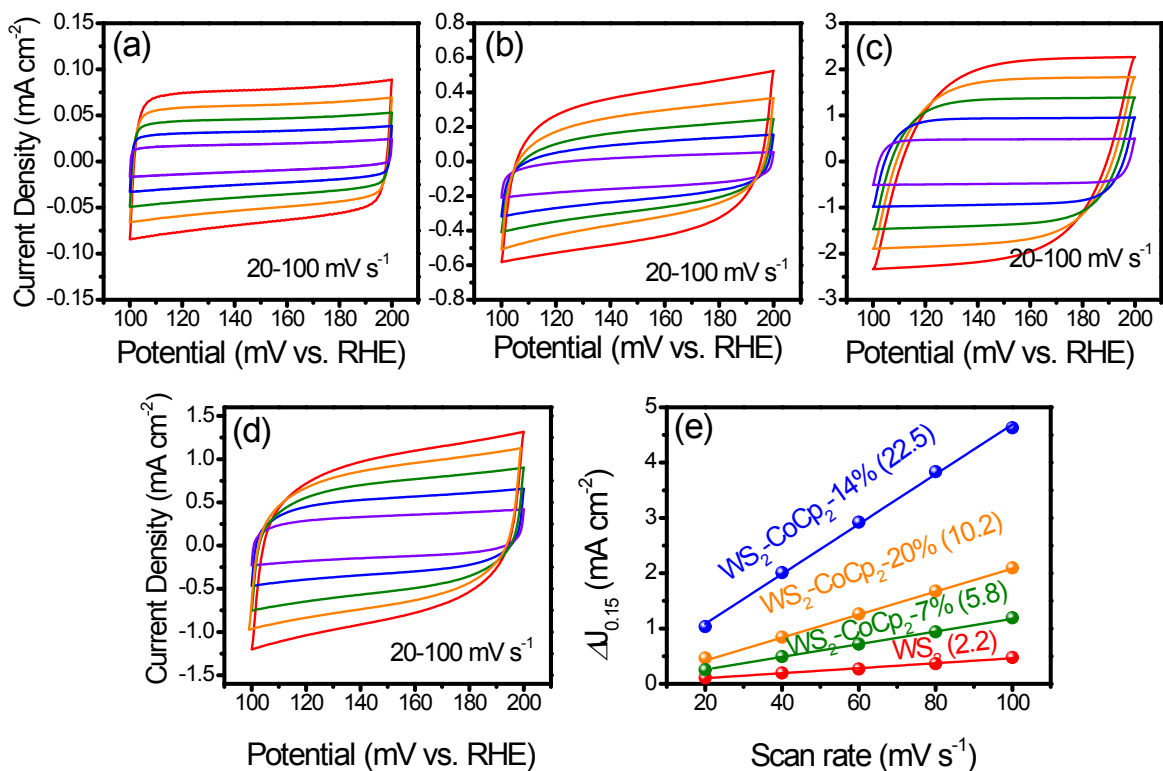
(b), (c) The overpotential for a current density of  $10 \text{ mA cm}^{-2}$  is  $0.36$  and  $0.32 \text{ V}$ , respectively, for  $\text{WS}_2$  and  $\text{WS}_2\text{-TGA}$ . The Tafel slope is  $70$  and  $110 \text{ mV dec}^{-1}$  for  $\text{WS}_2$  and  $\text{WS}_2\text{-TGA}$ . So the HER performance can be improved by the elimination of oleylamine ligand. In the case of  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$ , the HER performance was slightly reduced after the TGA treatment: the

overpotential for a current density of  $10 \text{ mA cm}^{-2}$  is  $0.17 \text{ V}$ , and the Tafel slope is  $40$  and  $70 \text{ mV dec}^{-1}$ , respectively, for before and after the TGA treatment. Nevertheless, the HER catalytic activity of  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$  is still higher than that of non-intercalated  $\text{WS}_2$  after the ligand elimination. This result confirmed the enhancement effect of the  $\text{CoCp}_2$  intercalation in the HER performance.



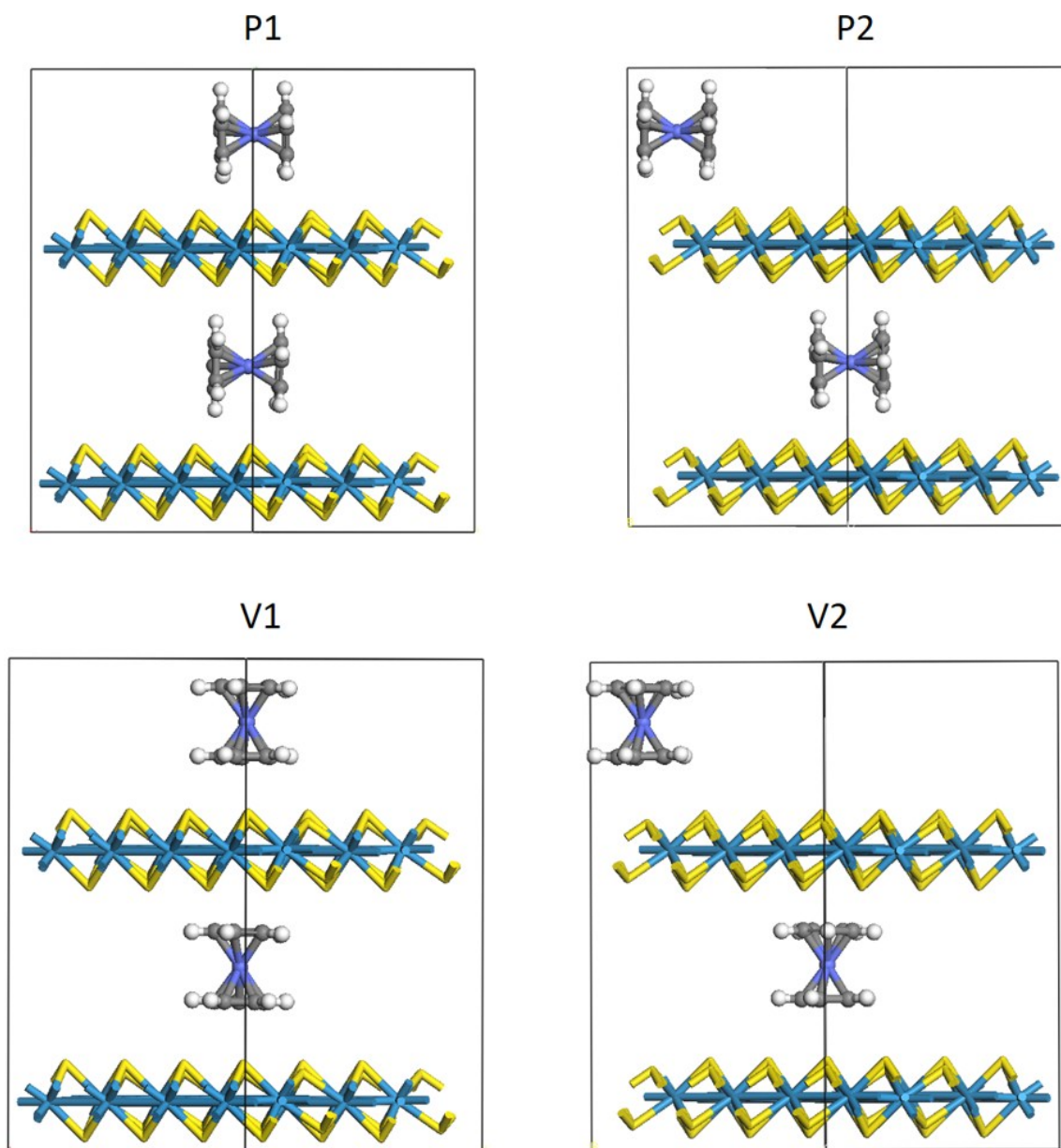
**Figure S11.** Nyquist plots for EIS measurements of  $\text{WS}_2$  (including the inset) and  $\text{WS}_2\text{-CoCp}_2$  (7%, 14%, and 20%) from 100 kHz to 0.1 Hz at a representative potential of -0.15 V (vs. RHE). The modified Randles circuit for fitting is shown on the right.

Electrochemical impedance spectroscopy (EIS) measurements of the samples were performed using a 100 kHz–0.1 Hz frequency range and an amplitude of 10 mV at  $\eta = 0.15$  V. In the high-frequency limit and under non-Faradaic conditions, the electrochemical system is approximated by the modified Randles circuit shown on the right panel, where  $R_s$  denotes the solution resistance, CPE is a constant-phase element related to the double-layer capacitance, and  $R_{ct}$  is the charge-transfer resistance from any residual Faradaic processes. A semicircle in the low-frequency region of the Nyquist plots represents the charge transfer process, with the diameter of the semicircle reflecting the charge-transfer resistance. The real ( $Z'$ ) and negative imaginary ( $-Z''$ ) components of the impedance are plotted on the  $x$  and  $y$  axes, respectively. The simulation of the EIS spectra using an equivalent circuit model allowed us to determine the charge transfer resistance,  $R_{ct}$ , which is a key parameter for characterizing the catalyst-electrolyte charge transfer process. The fitting parameters are listed in **Table S3**. The obtained  $R_{ct}$  values of  $\text{WS}_2$ ,  $\text{WS}_2\text{-CoCp}_2\text{-7\%}$ ,  $\text{WS}_2\text{-CoCp}_2\text{-14\%}$ , and  $\text{WS}_2\text{-CoCp}_2\text{-20\%}$  are 360, 69.1, 34.2, and 42.3  $\Omega$ , respectively, showing a significant decrease upon intercalation of  $\text{CoCp}_2$ . The  $R_{ct}$  is consistent with that of the HER performance. The reduced charge-transfer resistance plays a major role in enhancing the HER catalytic activity of the intercalated samples.

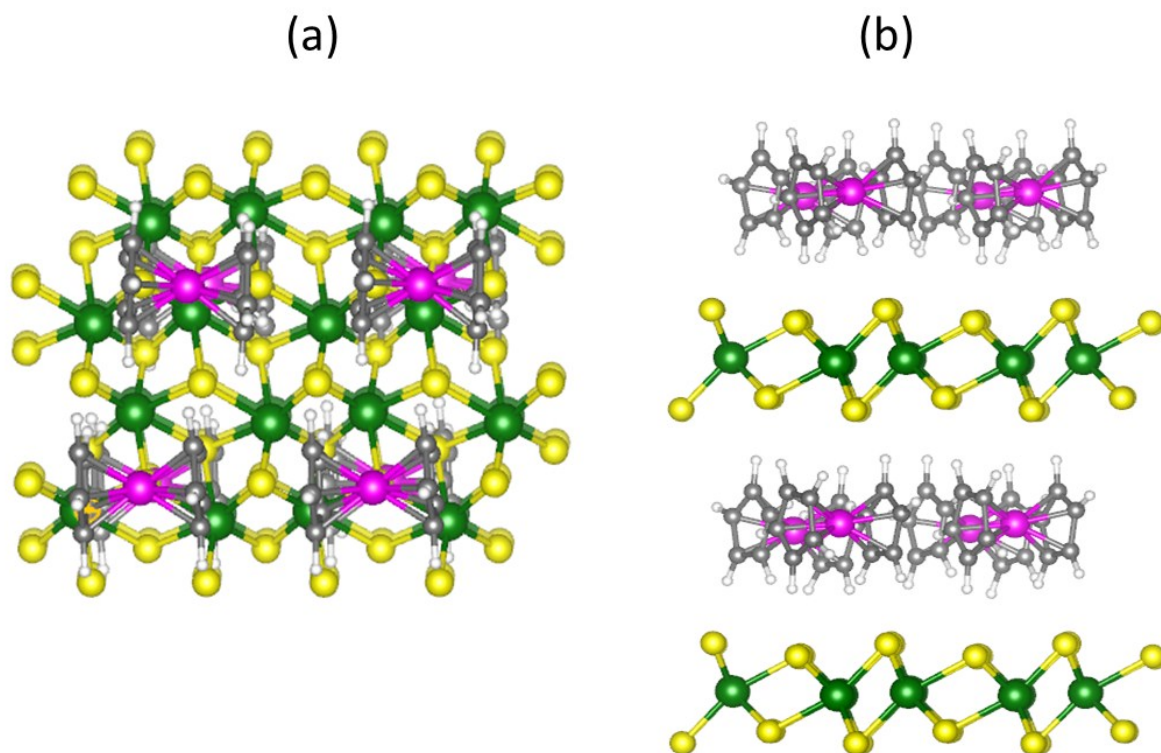


**Figure S12.** Cyclic voltammograms of (a) WS<sub>2</sub>, (b) WS<sub>2</sub>-CoCp<sub>2</sub>-7%, (c) WS<sub>2</sub>-CoCp<sub>2</sub>-14%, and (d) WS<sub>2</sub>-CoCp<sub>2</sub>-20% in a non-Faradaic region (0.1-0.2 V vs. RHE), at 20-100 mV s<sup>-1</sup> scan rates (with a step of 2 mV) and in 0.5 M H<sub>2</sub>SO<sub>4</sub> solution. (e) Difference ( $\Delta J$ ) between the anodic charging and cathodic discharging currents measured at 0.15 V (vs. RHE) and plotted as a function of the scan rate. The value in parenthesis represents the C<sub>dl</sub>, obtained by the half of the linear slope.

Cyclic voltammograms (CV) were measured at 0.1-0.2 V, in a non-Faradaic region, using various scan rates. The double-layer capacitance (C<sub>dl</sub>) was obtained as the slope of a linear fit of  $\Delta J$  vs. scan rate (20–100 mV s<sup>-1</sup>), where  $\Delta J$  is the difference between the anodic charging and cathodic discharging currents. The C<sub>dl</sub> values of WS<sub>2</sub>, WS<sub>2</sub>-CoCp<sub>2</sub>-7%, WS<sub>2</sub>-CoCp<sub>2</sub>-14%, and WS<sub>2</sub>-CoCp<sub>2</sub>-20% are 2.2, 5.8, 22.5, and 10.2 mF cm<sup>-2</sup>, respectively (see the summary in **Table S3**), showing a significant increase upon intercalation. The intercalated samples have very rough surfaces and can thus expose a large number of active sites. Therefore, the increased double-layer capacitance leads to the enhanced HER catalytic activity of the intercalated samples.

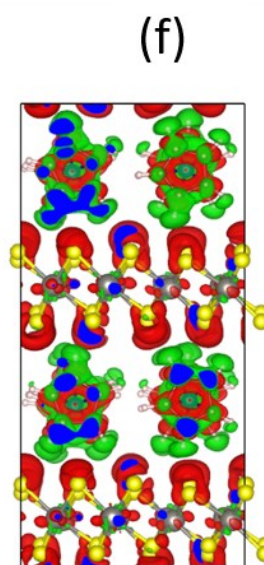
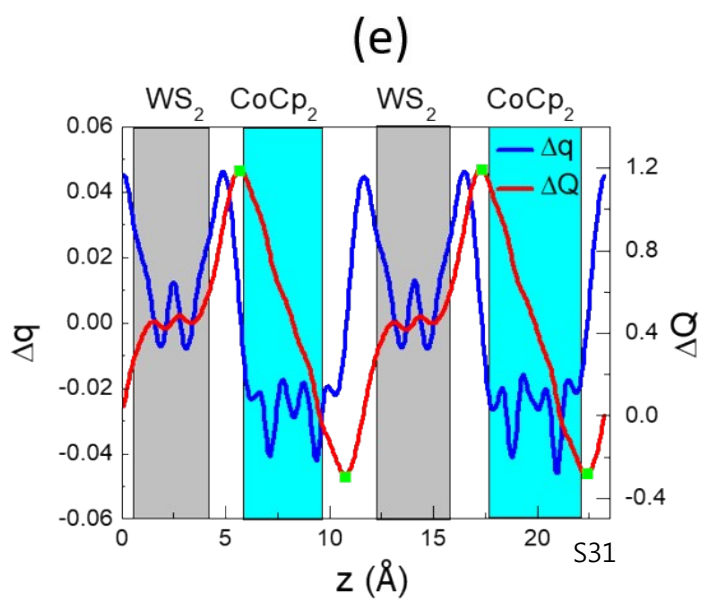
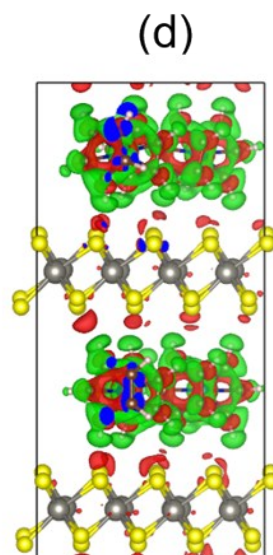
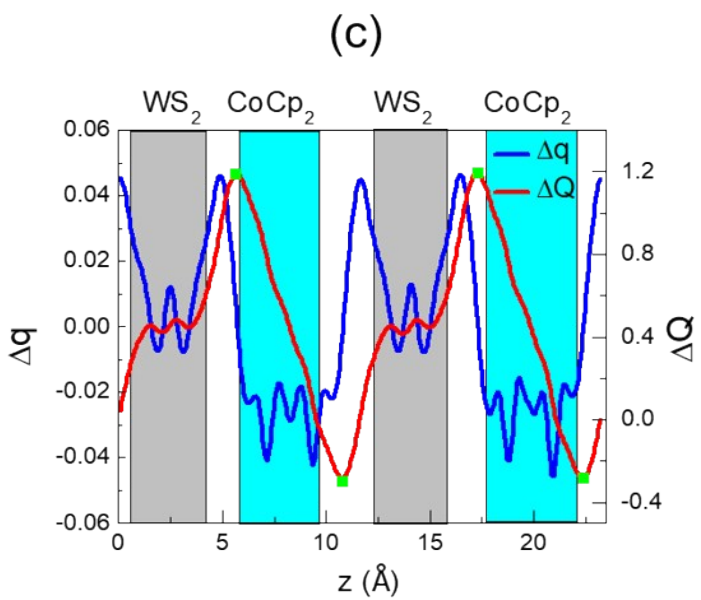
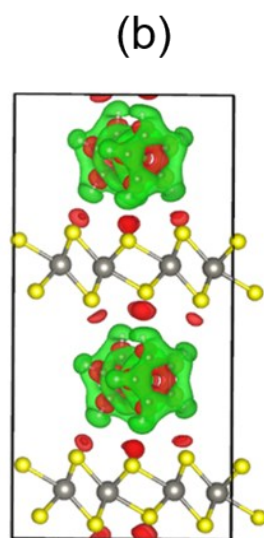
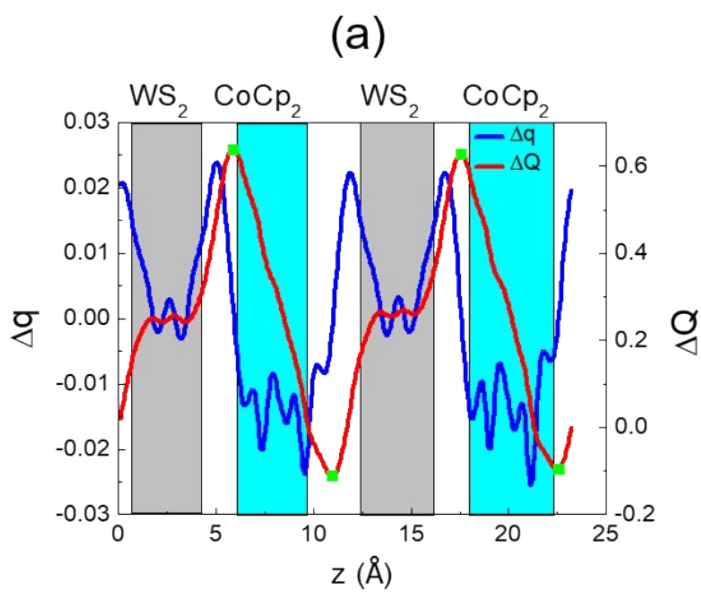


**Figure S13.** Structure of P1, P2, V1, and V2 configuration ( $4 \times 4$ )  $2WS_2-CoCp_2-2$ . Turquoise, yellow, blue, grey, and white balls represent the W, S, Co, C, and H atoms, respectively.



**Figure S14.** Structure (ball-and-stick model) of 1T' phase  $(4 \times 4)_2\text{WS}_2\text{-CoCp}_2\text{-8}$  (P4 configuration) in (a) top and (b) side views. Green, yellow, pink, gray, and white balls represent the W, S, Co, C, and H atoms, respectively.





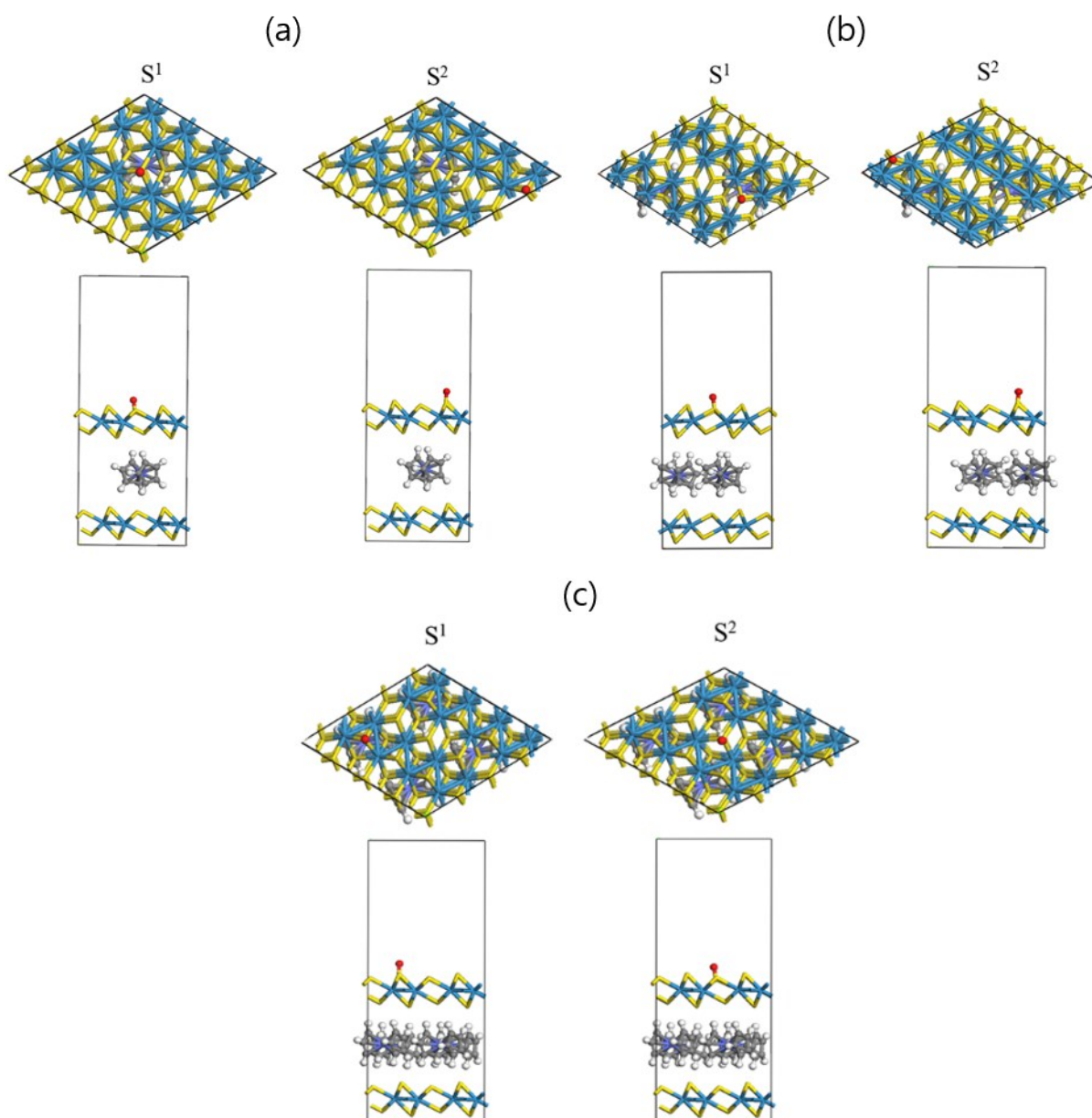
**Figure S15.** (a, c, e)  $\Delta q$  and  $\Delta Q$  of (4×4) 2WS<sub>2</sub>-CoCp<sub>2</sub>-2, (4×4) 2WS<sub>2</sub>-CoCp<sub>2</sub>-4, and (4×4) 2WS<sub>2</sub>-CoCp<sub>2</sub>-8 versus  $z$ :  $z_{\min}$  and  $z_{\max}$  values corresponding to  $Q_{\min}(z)$  and  $Q_{\max}(z)$  are shown by filled squares (■), respectively. (b, d, f) The differential charge density  $\Delta\rho(x, y, z)$  contour plot: charge accumulation and depletion regions are represented by red and green colors, respectively.

Following our previous studies, the amount of the charge transfer was calculated. The change in electron density (expressed in  $e \text{ \AA}^{-3}$ ) along the  $c$  ( $= z$ ) axis generated by the intercalation process was defined as  $\Delta\rho = \rho(z)\{\text{WS}_2\text{-CoCp}_2\} - \rho(z)\{\text{WS}_2\} - \rho(z)\{\text{CoCp}_2\}$ , averaged over the  $xy$  plane in a supercell.

The total electron density change ( $e$ ) was defined as  $\Delta q(z) = \Delta\rho(z)\Delta V$ , where  $\Delta V$  is the volume of a fine grid, *i.e.*,  $\Delta V = V_{\text{cell}}/N_c$ , in which  $V_{\text{cell}}$  is the total volume of the supercell and  $N_c$  is the number of fine grids. The thickness of each WS<sub>2</sub> layer was defined on the basis of the  $z$  coordinates of the S atoms in the upper and lower sublayers ( $S_L$  and  $S_U$ ), with  $z(S_L) < z(S_U)$ . The thickness of CoCp<sub>2</sub> molecules was obtained from the minimum and maximum  $z$  coordinates of its atoms. The actual thickness might be larger than that obtained using this definition, if the finite atomic size is taken into account. In addition,  $Q(z)$ , displayed on the right vertical axis represents the accumulated excess charge in the interval  $[0, z]$ :  $Q(z) = \sum_0^{z'} \Delta q(z')$ ,

*i.e.*, the integration of charge difference  $\Delta q(z')$  within  $z' < z < c$ . The amount of charge transfer was defined as  $\Delta Q = Q_{\max} - Q_{\min}$ , where  $Q_{\max}$  and  $Q_{\min}$  correspond to the maximum and minimum charge values in the regions (marked by ■) adjacent to the WS<sub>2</sub> layers and CoCp<sub>2</sub> molecules, respectively. The  $\Delta Q$  value per CoCp<sub>2</sub> is  $0.75e$ ,  $0.75e$ , and  $0.63e$ , respectively, for CoCp<sub>2</sub>-2, CoCp<sub>2</sub>-4, and CoCp<sub>2</sub>-8 complexes.





**Figure S16.** Structures of (a)  $(4 \times 4)$   $2\text{WS}_2\text{-CoCp}_2\text{-2}$ , (b)  $(4 \times 4)$   $2\text{WS}_2\text{-CoCp}_2\text{-4}$ , and (c)  $(4 \times 4)$   $2\text{WS}_2\text{-CoCp}_2\text{-8}$  in slab geometry, showing the two different adsorption sites ( $\text{S}^1$  and  $\text{S}^2$ ) for an  $\text{H}^+$  ion (red balls) for Volmer reaction. Turquoise, yellow, blue, grey, and white balls represent the W, S, Co, C, and H atoms, respectively. The top and side views are shown. The red ball denotes the H atoms adsorbed on the S atoms.

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