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Supplementary material (ESI)

for

"A Highly Sensitive Method for Time-Resolved Detection of O(¹D) Applied to Precise Determination of Absolute O(¹D) Reaction Rate Constants and O(³P) yields"

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The CH($A^2\Delta$) quasi-steady-state concentration, $[CH(A^2\Delta)]_{ss}$, and hence the CH($A^2\Delta \rightarrow X^2\Pi$) emission intensity at 431 nm, I_{em} , are proportional to the ratio of the $[CH(A^2\Delta)]$ formation and removal rates. Thus the general equation is:

$$I_{em} \propto [CH(A^2 \Delta)]_{ss} = B[C_2 H] \left([O(^1 D_2)] + \frac{[O(^3 P_J)]}{3.0} \right)$$
 (i)

Here, $B = k_{cheml}/\{A_{nm} + \Sigma_i(kI_{bXi}[X_i] + kI_{bR}[R])\}$, where the denominator is the sum of quenching rate constants for CH(A² Δ), kI_{bXi} and kI_{bR} , by species [X_i] and [R], respectively, and the Einstein coefficient, A_{nm} , for spontaneous emission for CH(A² $\Delta \rightarrow$ X² Π). [X_i] represents the concentrations of the photolytic precursors for O(¹D) and C₂H, and that of the buffer gas, [He], whereas [R] represents the concentration of any added reactant. [X_i] and [R] are treated separately because the kinetic effects of [X_i] remain essentially constant , even though [He] is altered slightly to compensate for added [R].

The temporal profile of the CH(A \rightarrow X) emission depends on the time dependences of $[O(^{1}D_{2})], [C_{2}H]$, and $[O(^{3}P_{J})]$, governed by:

$$O(^{1}D_{2}) + R \rightarrow \text{other products} \qquad k I_{aR} \qquad (5)$$

$$\rightarrow O(^{3}P_{J}) + \text{products} \qquad kl_{bR} \qquad (6)$$

$$O(^{1}D_{2}) + X_{i} \rightarrow \text{other products} \qquad \Sigma_{i}kl_{aXi} \qquad (7)$$

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$$\rightarrow O({}^{3}P_{J}) + \text{products} \qquad \Sigma_{i}kI_{bXi} \qquad (8)$$

$$C_{2}H + R \rightarrow \text{other products} \qquad k2_{aR} \qquad (9)$$

$$\rightarrow O({}^{3}P_{J}) + \text{products} \qquad k2_{bR} \qquad (10)$$

$$C_{2}H + X_{i} \rightarrow \text{products} \qquad \Sigma_{i}k2_{aXi} \qquad (11)$$

$$\rightarrow O({}^{3}P_{J}) + \text{products} \qquad \Sigma_{i}k2_{bXi} \qquad (12)$$

$$O({}^{3}P_{J}) + R \rightarrow \text{products} \qquad k3_{R} \qquad (13)$$

$$O({}^{3}P_{J}) + X_{i} \rightarrow \text{products} \qquad \Sigma_{i}k3_{Xi} \qquad (14)$$

Therefore, $[O(^1D_2)]_t = [O(^1D_2)]_0 \exp(-kl't)$ and $[C_2H]_t = [C_2H]_0 \exp(-k2't)$ where $k\delta' = k\delta_{aR}[R] + k\delta_{bR}[R] + \sum_i((k2_{aXi} + k2_{bXi})[X_i])$, with $\delta = 1$ or 2, and

$$\frac{d[O({}^{3}P_{J})]}{dt} = [O({}^{1}D_{2})]k_{1_{b}} + [C_{2}H]k_{2_{b}} - [O({}^{3}P_{J})]k_{3}$$
(ii)

Where the total removal rate of $O({}^{3}P_{J})$, $k3' = k3[R] + \Sigma_{i}(k3_{Xi}[X_{i}])$, giving

$$\left[O\left({}^{3}P_{J}\right)\right] = \left\{\frac{k1_{b}'}{k1'-k3'}\right\} \left[O\left({}^{1}D_{2}\right)\right]_{0} \left(\exp(-k3't) - \exp(-k1't)\right) + \left\{\frac{k2_{b}'}{k2'-k3'}\right\} \left[C_{2}H\right]_{0} \left(\exp(-k3't) - \exp(-k2't)\right)$$
(iii)

This leads to the full expression for I_{em} of

$$I_{em} \propto B[C_2H]_0[O(^1D_2)]_0 \left\{ \exp(-(k1'+k2')t) + \left\{ \frac{1}{3.0} \right\} \left\{ \frac{k1_b'}{k1'-k3'} \right\} \left\{ \exp(-(k3'+k2')t) - \exp(-(k1'+k2')t) \right\} \right\}$$

+ $B[C_2H]_0^2 \left\{ \left\{ \frac{1}{3.0} \right\} \left\{ \frac{k2_b'}{k2'-k3'} \right\} \left\{ \exp(-(k3'+k2')t) - \exp(-2k2')t \right\} \right\}$ (iv)

Two simplifications can be applied leading to a function having only two unknown parameters, kI_{Ra} and kI_R (= $kI_{Ra} + kI_{Rb}$). The first is to recognise the general great disparity in the magnitude of rate constants for reactions involving, on the one hand, O(¹D₂) and C₂H, and, on the other, O(³P_J) with neutral, stable molecules. This leads to elimination of

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k3' from equation (iv). The second concerns the low yield of O(³P_J) generated by reactions of C₂H. O(³P_J) cannot be generated efficiently from the C₂H + N₂O reaction because the *overall* rate constant is low at room temperature and below. If this holds also for R – either because it is void of O-atoms, the rate of reaction of R with C₂H is low, or the branching fraction to O(³P_J) for R + C₂H is negligible – then the second term of equation (iv) may be neglected, because, $[O^1D_2]_0(k1_b'/k1') \gg [C_2H]_0(k2_b'/k2')$ – assuming $[O^1D_2]_0$ is of the same order as, or greater than, $[C_2H]_0$.

To a very good approximation,

$$I_{em} \propto B[C_2H]_0[O(^1D)]_0 \left\{ \exp(-(k1'+k2')t) + \left\{\frac{1}{3.0}\right\} \left\{\frac{k1_b'}{k1'}\right\} \left\{ \exp(-k2't) - \exp(-(k1'+k2')t) \right\} \right\}$$
(v)

Thus at long times, the decay rate is that of [C₂H], given by $A(kl_b'/3kl')\exp(-k2't)$, where A is a constant.

In order to determine the values kI_{aR} and kI_R , the following procedure may be followed. A chemiluminescence profile is recorded in the presence of N₂O, C₂H₂, and He only. Fitting equation (v) to this profile gives the values for kI', k2', kI_b'/kI' (*i.e.*, their values when [R] = 0) and A. In cases when kI_b'/kI' is less than about 0.03, a second decay is recorded under conditions in which a small fraction of [He] is replaced by either [Ar] or [N₂]. This promotes quenching of O(¹D₂), and therefore increases the intensity of the longtime portion of the decay, and in turn increases the precision to which k2' is known. Thus k2' may be determined from the second profile while k1' and $k1_b'$ are always determined from the first.

Once these parameters are known, [R] may be added and equation (v) fit to the new profile, with the only unknown parameters now being kI_{bR} , kI_R , and $(k2_{aR} + k2_{bR})$. Again $(k2_{aR} + k2_{bR})$ may be determined more precisely, if required, by replacement of a small fraction of [He] by [N₂]. In practice, several sets of decay profiles are taken with different [R] to determine the values of the desired parameters. If $kI_{bR} \ll kI_{aR}$ these parameters are not strongly coupled, since kI_{bR} affects the long-time *intensity* of the decay profile, whilst kI_{aR} affects its initial *slope*.

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