# Simple models for the quick estimation of ground state splittings from hydrogen tunneling in alcohols and other compounds

Electronic Supplementary Information

-

Robert Medel\*

### Contents



<sup>\*</sup> Institute of Physical Chemistry, University of Goettingen, Tammannstr. 6, 37077 Goettingen. Germany. E-mail: rmedel@gwdg.de

#### <span id="page-1-0"></span>1 Scaled Eckart Barrier Model

By close inspection of the performance of the Eckart Barrier Model in Fig. 4 of the main document one notices that isotopologs typically have similar over- or underestimations. This supports the plausible expectation that isotopologs have more similar barrier shapes than chemically distinct species and thus also more akin deviations from the Eckart idealization. More similar errors in their quantum chemical evaluation likely also contribute. One can exploit this observation with eqn [\(S1\)](#page-1-1). Herein ∆(Eckart) corresponds to the calculation through the Eckart Barrier Model, given by eqn  $(S2)$ .

<span id="page-1-2"></span><span id="page-1-1"></span>
$$
\Delta_{\rm D}(\text{calc}) = \Delta_{\rm H}(\text{exp}) \frac{\Delta_{\rm D}(\text{Eckart})}{\Delta_{\rm H}(\text{Eckart})}
$$
(S1)

$$
\Delta(\text{Eckart}) = \frac{\omega_0 hc}{\sqrt{\pi e}} \exp\left[-\frac{2\pi}{|\omega_i|} \left(V_0 - \sqrt{E_0 V_0}\right)\right]
$$
(S2)

<span id="page-1-3"></span>For deuterated alcohols the correlation between the experimental tunneling splittings and those estimated from their lighter isotopologs is shown in Fig. [S1.](#page-1-3)



Figure S1: Correlation between the experimental tunneling splittings for deuterated alcohols and those calculated from their lighter isotopologs by the Scaled Eckart Barrier Model through eqn [\(S1\)](#page-1-1). The diagonal line represents perfect agreement.

In all cases the splitting is slightly to moderately underestimated, showing that the Eckart Barrier Model overestimates the isotope effect. For methanol the effect from the smaller increase of the tunneling mass by deuteration of the methyl group is better anticipated than from deuteration of the hydroxy group. Interestingly, the effect of deuteration on benzyl alcohol is also reasonably well captured, making it the only simple model discussed in this work to do so. However, the very extensive error cancellation for this case might be at least partly accidental.

Taking only hydroxy deuteration into account and excluding BnOH, one obtains for the 10 data points MSDF = 1.39 and MAX = 1.7 for *i*PrOD. This can be compared to MSDF = 1.87

and  $MAX = 3.2$  for  $iProof$  for the same alcohols in the original Eckart Barrier Model, showing the improvement.

With all values deviating in the same direction, it is tempting to employ a further empirical correction. A linear regression for the hydroxy deuteration data in Fig. [S1](#page-1-3) with a fixed slope of unity yields an intercept of 0.14, which converts to a scaling by factor 1.38. This is equivalent to moving all data points by 0.14 to the right in Fig. [1](#page-1-3) for overall better agreement with the diagonal line. Using eqns [\(S3\)](#page-2-0) and [\(S4\)](#page-2-1) brings the deviations further down to  $MSDF = 1.14$  and  $MAX = 1.32$  for  $HCCCH<sub>2</sub>OD$ .

$$
\Delta_{\rm D}(\text{calc}) = 1.38 \Delta_{\rm H}(\text{exp}) \frac{\Delta_{\rm D}(\text{Eckart})}{\Delta_{\rm H}(\text{Eckart})}
$$
(S3)

<span id="page-2-1"></span><span id="page-2-0"></span>
$$
\Delta_{\rm H}(\text{calc}) = 1.38^{-1} \Delta_{\rm D}(\text{exp}) \frac{\Delta_{\rm H}(\text{Eckart})}{\Delta_{\rm D}(\text{Eckart})}
$$
(S4)

### <span id="page-3-0"></span>2 Performance overview

Table S1: Performance of the investigated models for the reproduction of tunneling splittings of alcohols. MSDF is the mean symmetric deviation factor, defined as the arithmetic average of the ratios between experiment and calculation taken as  $\geq$  1, and MAX is the maximum symmetric deviation factor.  $n = 27$  is the full set of splittings given in Table 1 of the main document, with the exclusion of only  $BnOH/D$ .  $n = 21$  is a subset of alcohols with more similar properties with the further exclusion of CH<sub>3</sub>OH/D, MeOCH<sub>2</sub>OH, CF<sub>3</sub>CH<sub>2</sub>OH/D and FCH<sub>2</sub>OH.  $n = 10$  are deuterated alcohols (expect BnOD) for which also the splitting of the protiated isotopolog is known, so that models for the isotope effect can be applied.  $n = 77$  are isotopologic pairs of chemically very diverse systems with hydrogens tunneling, with the 10 pairs of alcohols forming a subset.



## <span id="page-4-0"></span>3 List of tunneling splittings in the Universal Direct Correlation Model

<span id="page-4-1"></span>Table S2: Isotopologic pairs of experimental hydrogen tunneling splittings for diverse compounds reported in the literature and used in the Universal Direct Correlation Model. Given values are in  $cm^{-1}$ hc and were rounded to two leading digits. The list continues in Table 1 of the main document (excluding benzyl alcohol) and in Table [S3.](#page-5-0)



<span id="page-5-0"></span>

category	system	$\Delta_H(exp)$	$\Delta_{\rm D}({\rm exp})$
out-of-plane bending	$H_2O \cdots HF$	$6.4\cdot10^{+1}$ $^{42}$	$3.2 \cdot 10^{+1}$ 42
acceptor switching	water dimer	$9.4 \text{ }^{43}$	$1.8\ ^{43}$
interchange (lower)	water dimer	$7.5 \cdot 10^{-1}$ 43	$3.9 \cdot 10^{-2}$ 43
bifurcation (lower)	water dimer	$2.10^{-2}$ 43	$2.3 \cdot 10^{-4}$ 43
flip	water trimer	$4.4{\cdot}10^{+2}$ $^{43}$	$2.1 \cdot 10^{2}$ $^{43}$
bifurcation	water trimer	$9.6 \cdot 10^{-3}$ 43	$1.7 \cdot 10^{-4}$ 43
bifurcation	$water + carbonmonoxide$	$5.6 \cdot 10^{-1}$ 44	$3.4 \cdot 10^{-2}$ 44
solid solution in metal	manganese	$5.2 \cdot 10^{+1}$ 45	$1.3 \cdot 10^{1}$ $^{45}$
	niobium doped with oxygen	$1.9\frac{46}{ }$	$1.6 \cdot 10^{-1}$ 46
	niobium doped with nitrogen	$1.4 \, 46$	$1.1{\cdot}10^{-1}$ $^{46}$
rocking	vinyl radical	$5.4 \cdot 10^{-1}$ 47	$4.0 \cdot 10^{-2}$ 47
hydrogen transfer	malonaldehyde	$2.2 \cdot 10^{+1}$ 48	$2.9^{48}$
	tropolone	$9.7 \cdot 10^{-1}$ 49	$5.1 \cdot 10^{-2}$ 49
	6-hydroxy-2-formylfulvene	$1.2{\cdot}10^{+2}$ $^{50}$	$3.6 \cdot 10^{+1}$ 50
ammonium reorientation	$(NH_4)_2SnCl_6$	$2.4\!\cdot\!10^{-2}$ 51	$2.4 \cdot 10^{-4}$ 52
ammonium torsion	$(NH_4)_2IrCl_6$	$9.3 \cdot 10^{+1}$ 53	$3.1 \cdot 10^{+1}$ 53
	$(NH_4)_2RuCl_6$	$7.0 \cdot 10^{+1}$ 54	$3.9{\cdot}10^{+1}$ $^{\,54}$
	$(NH_4)_2SnCl_6$	$1.3{\cdot}10^{+2}$ $^{54}$	$4.2 \cdot 10^{+1}$ 54
methyl torsion	nitromethane crystal 1 bar	$2.8 \cdot 10^{-1}$ 55	$1.4 \cdot 10^{-2}$ 55
	nitromethane crystal 4.8 kbar	$4.4 \cdot 10^{-1}$ 55	$2.3 \cdot 10^{-2}$ 55
	acetamide crystal	$2.6 \cdot 10^{-1}$ 56	$9.5{\cdot}10^{-3}$ $^{56}$
	$(CH_3)_2$ SnCl <sub>2</sub> crystal	$2.6 \cdot 10^{-1}$ 57	$2.1 \cdot 10^{-2}$ 57
	$(CH_3COO)_2Cu·H_2O$ crystal	$2.1 \cdot 10^{-3}$ 58	$2.0 {\cdot} 10^{-5}$ $^{58}$
	acetylsalicylic acid crystal	$9.8\!\cdot\!10^{-3}$ $^{59}$	$9.0 \cdot 10^{-5}$ 59
	toluene crystal	$2.0 \cdot 10^{-1}$ 60	$8.9{\cdot}10^{-3}$ 60

Table S3: Continuation of Table [S2.](#page-4-1)

<span id="page-6-0"></span>4 Further examples for the application of the Direct Correlation Model



Figure S2: Correlation between the experimental inversion splittings of N-single-deuterated and N-protiated symmetric primary amines and amides. Used values with references are given in Table [S4.](#page-6-1) lg  $(\Delta_{HD}/cm^{-1}hc) = 1.135 \cdot$ lg  $(\Delta_{HH}/cm^{-1}hc) - 0.436$ . MSDF = 1.026 and MAX = 1.058 for cyanamide.

<span id="page-6-1"></span>Table S4: Experimental inversion splittings for N-protiated and N-single-deuterated symmetric primary amines, i.e. with both hydrogen atoms being equivalent in the inverting conformation. Values are rounded to two leading digits and given in  $cm^{-1}hc$ .

system	$\Delta_{HH}(\exp)$	$\Delta_{HD}(\exp)$
cyanamide	$5.0 \cdot 10^{+1}$ <sup>25</sup>	$3.2 \cdot 10^{+1}$ 25
aniline	$4.1{\cdot}10^{+1}$ $^{22}$	$2.4{\cdot}10^{+1}$ $^{22}$
methylamine	$9.7 \cdot 10^{-1}$ 21	$3.3{\cdot}10^{-1}$ 61
isocyanamide	$3.7 \cdot 10^{-1}$ 27	$1.2{\cdot}10^{-1}$ $^{27}$
chloroamine	$2.3 \cdot 10^{-3}$ 62	$3.8 \cdot 10^{-4}$ 63



Figure S3: Correlation between the experimental inversion splittings of N-double-deuterated and N-single-deuterated symmetric primary amines and amides. Used values with references are given in Tables [S2](#page-4-1) and [S4.](#page-6-1) lg  $(\Delta_{DD}/cm^{-1}hc) = 1.24 \cdot \lg (\Delta_{HD}/cm^{-1}hc) - 0.61$ . MSDF = 1.12 and  $MAX = 1.19$  for methylamine.



Figure S4: Correlation between the experimental tunneling splittings of deuterated and protiated thiols. Used values with references are given in Table [S2.](#page-4-1) Values for trifluoromethanethiol were multiplied by 2/3. lg  $(\Delta_D/cm^{-1}hc) = 1.44 \cdot \lg (\Delta_H/cm^{-1}hc) - 0.94$ . MSDF = 1.11 and MAX = 1.21 for ethanethiol.



Figure S5: Correlation between the experimental tunneling splittings of deuterated and protiated phenols. Values with references given in Table [S2](#page-4-1) were multiplied with  $1/2$ . lg  $(\Delta_D/cm^{-1}hc)$  = 1.419 · lg ( $\Delta_H / \text{cm}^{-1}$ *hc*) − 1.123. MSDF = 1.0028 and MAX = 1.0042 for 4-chlorophenol.



Figure S6: Correlation between the experimental tunneling splittings of solid solutions of hy-drogen in metals. Values with references are given in Table [S3.](#page-5-0) lg  $(\Delta_D/cm^{-1}hc) = 1.311$ .  $\log(\Delta_H/cm^{-1}hc) - 1.135$ . MSDF = 1.014 and MAX = 1.021 for H/D in niobium doped with oxygen.



Figure S7: Correlation between the experimental tunneling splittings of hydrogen transfer system. Values with references are given in Table [S3.](#page-5-0) lg  $(\Delta_D/cm^{-1}hc) = 1.35 \cdot$  lg  $(\Delta_H/cm^{-1}hc) - 1.29$ .  $MSDF = 1.07$  and  $MAX = 1.10$  for malonaldehyde.



Figure S8: Correlation between the experimental tunneling splittings of system with methyl group tunneling. Values with references are given in Table [S3.](#page-5-0) lg  $(\Delta_D/cm^{-1}hc) = 1.39 \lg (\Delta_H/cm^{-1}hc) -$ 1.10. MSDF = 1.23 and MAX = 1.44 for  $(CH_3)_2$ SnCl<sub>2</sub>.

### <span id="page-10-0"></span>5 Example inputs and outputs

Input and output files for the example methanol are provided as part of this supplementary information as separate files. Before calculating other alcohols it is recommended to do a quick rerun to check whether the expected result is reproduced as stated in the output file as well as in the included comment in the input file. Default settings might change in future versions of the software, so that adjustments to the input might become necessary. The inputs were tested with Gaussian 09 Rev. E.01 and Gaussian 16 Rev. A.03. The provided cartesian coordinates in the input files are the optimized values taken from the output files.

A spreadsheet for the application of the models is provided as a separate file as well, again with methanol as an example.

### <span id="page-11-0"></span>6 Discussion of metrics for the evaluation of model perfomance

Finding an appropriate metric to evaluate and compare performance between models is not trivial.

#### <span id="page-11-1"></span>6.1 Mean absolute error (MAE)

With tunneling splittings spanning several orders of magnitude, common metrics for absolute errors, such as the *mean absolute error* (MAE, eqn [\(S5\)](#page-11-5)), become meaningless.

<span id="page-11-5"></span>
$$
MAE = \frac{1}{n} \sum_{j}^{n} |\Delta_j(\exp) - \Delta_j(\text{calc})|
$$
 (S5)

An overestimation by 0.1 cm<sup>-1</sup>hc represents pinpoint accuracy for HCCCCH<sub>2</sub>OH ( $\Delta = 21.8 \text{ cm}^{-1}$ hc), but very poor performance for  $CF_3CH_2OD$  ( $\Delta = 0.007$  cm<sup>-1</sup>hc), so these two results should not contribute equally to the sum of errors.

#### <span id="page-11-2"></span>6.2 Mean absolute percentage error (MAPE)

A common metric for relative errors is the *mean absolute percentage error* (MAPE, eqn [\(S6\)](#page-11-6)).

<span id="page-11-6"></span>
$$
MAPE = \frac{100\%}{n} \sum_{j}^{n} \frac{|\Delta_j(\exp) - \Delta_j(\text{calc})|}{\Delta_j(\exp)} \tag{S6}
$$

It works well for small relative errors but for larger ones the metric shows pronounced asymmetry. An overestimation by a factor of two is twice as heavily penalized  $(|+100\%|)$  as an underestimation by the same factor  $(|-50\%|)$ , while they arguably should be valued the same, at least when preferring a multiplicative scale.

#### <span id="page-11-3"></span>6.3 Symmetric mean absolute percentage error (SMAPE)

One metric in use to solve the symmetry problem is the symmetric mean absolute percentage error (SMAPE, eqn [\(S7\)](#page-11-7)), which uses instead the average of the calculated and experimental value as the reference in the denominator. [64,](#page-16-7)[65](#page-16-8)

<span id="page-11-7"></span>
$$
\text{SMAE} = \frac{100\%}{n} \sum_{j}^{n} \frac{|\Delta_j(\exp) - \Delta_j(\text{calc})|}{[\Delta_j(\exp) + \Delta_j(\text{calc})]/2}
$$
(S7)

However, it is hard to interpret. The previous example of a model being constantly a factor two off target results in an odd value of  $SMAE = 67\%$ . This metric is also very forgiving to larger errors. A strong over- or underestimation by a factor of 10 yields a penalty of 164%, which is only about four times as large as for an over- or underestimation by a far more acceptable factor of 1.5 (40%). As a matter of opinion, this might be seen as 'robustness' or 'ignorance' of the metric to outliers.

#### <span id="page-11-4"></span>6.4 Mean absolute geometric error (MAGE)

Easier interpretable is the *mean absolute geometric error* (MAGE, eqn  $(S8)$ ). <sup>[66](#page-16-9)</sup>

<span id="page-11-8"></span>
$$
\text{MAGE} = \exp\left[\frac{1}{n}\sum_{j}^{n} \left| \ln\left(\frac{\Delta_j(\exp)}{\Delta_j(\text{calc})}\right) \right| \right]
$$
(S8)

It can be rationalized as the geometric mean of ratios between calculated and experimental splittings, when always the smaller one is used in the denominator, so that ratios  $\geq 1$  result. MAGE is also the MAE of the logarithmized tunneling splittings back converted to obtain a 'typical' ratio between calculation and experiment.

It is as well rather forgiving to larger errors due to the logarithmic scaling. For the discussed example, the deviation by a factor of 10 is penalized  $\ln(10)/\ln(1.5) \approx 5.7$  times as strongly as the one with a factor of 1.5.

#### <span id="page-12-0"></span>6.5 Mean symmetric deviation factor (MSDF) and mean normalized absolute factor error (MNAFE)

Arguably even easier interpretable and more intuitive than MAGE is the mean symmetric deviation factor (MSDF, eqn  $(S9)$ ).

<span id="page-12-2"></span><span id="page-12-1"></span>
$$
MSDF = \frac{1}{n} \sum_{j}^{n} \exp \left| \ln \left( \frac{\Delta_j(\exp)}{\Delta_j(\text{calc})} \right) \right| \tag{S9}
$$

This metric was used before in a mean and median version.<sup>[67](#page-16-10)</sup> It is also closely related to the *mean* normalized absolute factor error MNAFE (eqn  $(S10)$ ). <sup>[68](#page-16-11)</sup>

$$
\text{MNAFE} = \frac{1}{n} \sum_{j}^{n} \left[ \exp \left| \ln \left( \frac{\Delta_j(\exp)}{\Delta_j(\text{calc})} \right) \right| - 1 \right] \tag{S10}
$$

Back converting the logarithmic deviations before the averaging (and not after as with MAGE) has some advantages. MSDF can be rationalized as the arithmetic (instead of the geometric) mean of the ratios between calculated and experimental values. This is equivalent to arithmetically averaging relative or percentage errors of tunneling splittings (as it is done with MNAFE). E.g., ratios of 1.80 (80%) and 1.20 (20%) will arithmetically average to 1.50 (50%) instead geometrically to 1.47 (47%) with MAGE. MSDF and MNAFE are thus closely related to a symmetrized MAPE with a flexible denominator containing always the smaller value of either  $\Delta_j$  (exp) or  $\Delta_j$  (calc) as reference.

MSDF and MNAFE penalize larger relative errors more heavily and more intuitively than both SMAPE and MAGE. A deviation by a factor of 10 weights 18 times as much as a deviation by a factor of 1.5, as one would intuitively expect from the associated percentage errors of  $900\% / 50\% =$ 18.

So why prefer MSDF over MNAFE (optionally multiplied by 100%)? The relative errors contributing to MNAFE have the problem that they refer to either  $\Delta_i(\exp)$  or  $\Delta_i(\text{calc})$ , depending on which value is smaller, to reach the desired symmetry. This needs elaboration and special attention of the reader, as values of MNAFE could be easily confused with those of the commonly used MAPE. A stated underestimation by 200% is hard to understand and could be easily misinterpreted as the calculated value having the wrong sign. In a multiplicative system the desired symmetry arises more natural, which is reflected in the language. It is common to state an underestimation by a factor of 3, even though a factor of 1/3 would be arguably more correct.

# <span id="page-13-0"></span>7 Computed quantities for additional alcohols

Table S5: Calculated barrier heights with  $(V_0)$  and without  $(V_{el})$  harmonic vibrational zero-point correction, as well as calculated harmonic torsional wavenumbers for the gauche equilibrium geometry  $\omega_0$  and for the transition state  $\omega_i$  of alcohols with a protiated (index H) or deuterated (index D) hydroxy group. Splittings and barrier heights are given in cm<sup>-1</sup>hc, wavenumbers in cm<sup>-1</sup>. The B3LYP-D3(BJ)/may-cc-pVTZ level was used.

alcohol					$V_{0, \mathrm{H}}$ $V_{0, \mathrm{D}}$ $V_{\mathrm{el}}$ $\omega_{0, \mathrm{H}}$ $\omega_{0, \mathrm{D}}$ $\omega_{i, \mathrm{H}}$ $\omega_{i, \mathrm{D}}$ $\frac{\omega_{i, \mathrm{H}}^2}{\omega_{i, \mathrm{D}}^2}$
NCCMe <sub>2</sub> OH 471 467 452 313 244 299i 223i 1.80					
1PrOH Tg 421 407 398 275 193				302i 228i 1.76	
1BuOH TTg 411 396 387 277 195				300i 226i	- 1.76

### References

- <span id="page-14-0"></span>[1] J. Nakagawa, K. Kuwada and M. Hayashi, Bull. Chem. Soc. Jpn., 1976, 49, 3420–3432.
- <span id="page-14-1"></span>[2] J. Nakagawa and M. Hayashi, J. Mol. Spectrosc., 1981, 85, 327–340.
- <span id="page-14-2"></span>[3] Y. Kawashima, Y. Tanaka, T. Uzuyama and E. Hirota, J. Phys. Chem. A, 2021, 125, 1166– 1183.
- <span id="page-14-3"></span>[4] A. Mirri, F. Scappini, R. Cervellati and P. Favero, J. Mol. Spectrosc., 1976, 63, 509–520.
- <span id="page-14-4"></span>[5] J. H. Griffiths and J. E. Boggs, *J. Mol. Spectrosc.*, 1975, **56**, 257–269.
- <span id="page-14-5"></span>[6] A. Cox, C. Rego and R. Stevens, *J. Mol. Struct.*, 1990, **223**, 185–205.
- <span id="page-14-6"></span>[7] J. Nakagawa and Y. Miyake, *J. Mol. Spectrosc.*, 1986, **119**, 201–213.
- <span id="page-14-7"></span>[8] J. Nakagawa, H. Okutani and M. Hayashi, J. Mol. Spectrosc., 1982, 94, 410–425.
- <span id="page-14-8"></span>[9] J. Nakagawa, A. Nagayama and M. Hayashi, J. Mol. Spectrosc., 1983, 99, 415–430.
- <span id="page-14-9"></span>[10] D. Christen, L. Coudert, R. Suenram and F. Lovas, *J. Mol. Spectrosc.*, 1995, 172, 57–77.
- <span id="page-14-10"></span>[11] L. Evangelisti, Q. Gou, L. Spada, G. Feng and W. Caminati, *Chem. Phys. Lett.*, 2013, **556**, 55–58.
- <span id="page-14-11"></span>[12] N. Larsen and F. Nicolaisen, *J. Mol. Struct.*, 1974, **22**, 29–43.
- <span id="page-14-12"></span>[13] N. Larsen, *J. Mol. Struct.*, 1986, **144**, 83–99.
- <span id="page-14-13"></span>[14] N. W. Larsen and L. Schulz, *J. Mol. Struct.*, 2009, 10.
- <span id="page-14-14"></span>[15] P. Helminger, W. C. Bowman and F. C. De Lucia, J. Mol. Spectrosc., 1981, 85, 120–130.
- <span id="page-14-15"></span>[16] R. H. Hunt and R. A. Leacock, J. Chem. Phys., 1966, 45, 3141–3147.
- <span id="page-14-16"></span>[17] S. Willitsch, U. Hollenstein and F. Merkt, J. Chem. Phys., 2004, 120, 1761–1774.
- <span id="page-14-17"></span>[18] E. Fischer and I. Botskor, *J. Mol. Spectrosc.*, 1984, **104**, 226–247.
- <span id="page-14-18"></span>[19] Š. Urban, V. Špirko, D. Papoušek, J. Kauppinen, S. Belov, L. Gershtein and A. Krupnov, J. Mol. Spectrosc., 1981, 88, 274–292.
- <span id="page-14-19"></span>[20] L. Fusina, G. di Lonardo and J. Johns, J. Mol. Spectrosc., 1985, 112, 211–221.
- <span id="page-14-20"></span>[21] K. Takagi and T. Kojima, *J. Phys. Soc. Jpn.*, 1971, **30**, 1145–1157.
- <span id="page-14-21"></span>[22] R. Kydd and P. Krueger, Chem. Phys. Lett., 1977, 49, 539–543.
- <span id="page-14-22"></span>[23] R. Penn and J. E. Boggs, J. Mol. Spectrosc., 1973, 47, 340–346.
- <span id="page-14-23"></span>[24] J. E. Wollrab and V. W. Laurie, *J. Chem. Phys.*, 1968, **48**, 5058–5066.
- <span id="page-14-24"></span>[25] Z. Kisiel, A. Kraśnicki, W. Jabs, E. Herbst, B. P. Winnewisser and M. Winnewisser, J. Phys. Chem. A, 2013, 117, 9889–9898.
- <span id="page-14-25"></span>[26] L. B. Favero, B. Velino, A. Maris and W. Caminati, J. Mol. Struct., 2002, 612, 357–367.
- <span id="page-14-26"></span>[27] M. Winnewisser and J. Reinstaedtler, *J. Mol. Spectrosc.*, 1986, **120**, 28–48.
- <span id="page-14-27"></span>[28] J. Loreau, J. Liévin, Y. Scribano and A. van der Avoird, *J. Chem. Phys.*, 2014, 141, 224303.
- <span id="page-14-28"></span>[29] R. Kydd, Spectrochim. Acta, Part A Mol. Spectrosc., 1979, 35, 409–413.
- <span id="page-14-29"></span>[30] S. Tsunekawa and T. Kojima, J. Phys. Soc. Jpn., 1980, 49, 1957–1964.
- <span id="page-14-30"></span>[31] S. Tsunekawa and T. Kojima, J. Phys. Soc. Jpn., 1978, 44, 1925–1930.
- <span id="page-15-0"></span>[32] S. Yu, B. J. Drouin, J. C. Pearson and H. M. Pickett, Astrophys. J. Suppl. Ser., 2009, 180, 119–124.
- <span id="page-15-1"></span>[33] M. Araki, H. Ozeki and S. Sato, Mol. Phys., 1999, 97, 177–183.
- <span id="page-15-2"></span>[34] J. T. Farrell, M. A. Suhm and D. J. Nesbitt, J. Chem. Phys., 1996, 104, 9313–9331.
- <span id="page-15-3"></span>[35] M. D. Schuder, C. M. Lovejoy, R. Lascola and D. J. Nesbitt, J. Chem. Phys., 1993, 99, 4346–4362.
- <span id="page-15-4"></span>[36] M. D. Schuder, D. D. Nelson and D. J. Nesbitt, J. Chem. Phys., 1993, 99, 5045–5060.
- <span id="page-15-5"></span>[37] H. J. Wörner, X. Qian and F. Merkt, *J. Chem. Phys.*, 2007, 126, 144305.
- <span id="page-15-6"></span>[38] S. G. Kukolich, E. G. Mitchell, S. J. Carey, M. Sun and B. A. Sargus, J. Phys. Chem. A, 2013, 117, 9525–9530.
- <span id="page-15-7"></span>[39] M. Sun, Y. Wang, S. J. Carey, E. G. Mitchell, J. Bowman and S. G. Kukolich, J. Chem. Phys., 2013, 139, 084316.
- <span id="page-15-8"></span>[40] G. Feng, L. B. Favero, A. Maris, A. Vigorito, W. Caminati and R. Meyer, J. Am. Chem. Soc., 2012, 134, 19281–19286.
- <span id="page-15-9"></span>[41] L. Evangelisti, P. Écija, E. J. Cocinero, F. Castaño, A. Lesarri, W. Caminati and R. Meyer, J. Phys. Chem. Lett., 2012, 3, 3770–3775.
- <span id="page-15-10"></span>[42] Z. Kisiel, A. C. Legon and D. J. Millen, Proc. R. Soc. Lond. A, 1982, 381, 419–442.
- <span id="page-15-11"></span>[43] J. O. Richardson, S. C. Althorpe and D. J. Wales, J. Chem. Phys., 2011, **135**, 124109.
- <span id="page-15-12"></span>[44] R. Bumgarner, S. Suzuki, P. A. Stockman, P. G. Green and G. A. Blake, Chem. Phys. Lett., 1991, 176, 123–127.
- <span id="page-15-13"></span>[45] V. E. Antonov, B. Dorner, V. K. Fedotov, G. Grosse, A. S. Ivanov, A. I. Kolesnikov, V. V. Sikolenko and F. E. Wagner, J. Alloys Compd., 2002, 5.
- <span id="page-15-14"></span>[46] P. Gutsmiedl, M. Schiekhofer, K. Neumaier and H. Wipf, in: Quantum Aspects of Molecular Motions in Solids, Berlin, Heidelberg, 1987, pp. 158–162.
- <span id="page-15-15"></span>[47] K. Tanaka, M. Hayashi, M. Ohtsuki, K. Harada and T. Tanaka, J. Chem. Phys., 2009, 131, 111101.
- <span id="page-15-16"></span>[48] S. L. Baughcum, Z. Smith, E. B. Wilson and R. W. Duerst, J. Am. Chem. Soc., 1984, 106, 2260–2265.
- <span id="page-15-17"></span>[49] J. C. Keske, W. Lin, W. C. Pringle, S. E. Novick, T. A. Blake and D. F. Plusquellic, J. Chem. Phys., 2006, 124, 074309.
- <span id="page-15-18"></span>[50] Z. N. Vealey, L. Foguel and P. H. Vaccaro, J. Phys. Chem. Lett., 2018, 9, 4949–4954.
- <span id="page-15-19"></span>[51] M. Prager, W. Press, B. Alefeld and A. Hüller, *J. Chem. Phys.*, 1977, 67, 5126–5132.
- <span id="page-15-20"></span>[52] L. P. Ingman, E. Koivula, Z. T. Lalowicz, M. Punkkinen and E. E. Ylinen, Z. Für Phys. B Condens. Matter, 1990, 81, 175–181.
- <span id="page-15-21"></span>[53] Y. Kume and T. Asaji, Z. Für Naturforschung A, 2002, 57, 504–508.
- <span id="page-15-22"></span>[54] Y. Kume, D. Amino and T. Asaji, J. Mol. Struct., 2013, 1043, 1–6.
- <span id="page-15-23"></span>[55] S. F. Trevino, in: Quantum Aspects of Molecular Motions in Solids, Springer Berlin Heidelberg, Berlin, Heidelberg, 1987, vol. 17, pp. 9–18.
- <span id="page-15-24"></span>[56] A. Heidemann, M. Prager and M. Monkenbusch, Z. Physik B - Condensed Matter, 1989, 76, 77–88.
- <span id="page-16-0"></span>[57] J. Stanislawski, M. Prager and W. Häusler, Phys. B Condens. Matter, 1989, 156-157, 356-358.
- <span id="page-16-1"></span>[58] Z. T. Lalowicz, U. Werner and W. Müller-Warmuth, Z. Für Naturforschung A, 1988, 43, 219–227.
- <span id="page-16-2"></span>[59] M. Johnson, B. Frick and H. Trommsdorff, Chem. Phys. Lett., 1996, 258, 187–193.
- <span id="page-16-3"></span>[60] D. Cavagnat, A. Magerl, C. Vettier and S. Clough, J. Phys. C: Solid State Phys., 1986, 19, 6665–6672.
- <span id="page-16-4"></span>[61] N. Ohashi, M. Oda and K. Takagi, J. Mol. Spectrosc., 1991, 145, 180–191.
- <span id="page-16-5"></span>[62] G. Cazzoli and D. Lister, *J. Mol. Spectrosc.*, 1973, **45**, 467-474.
- <span id="page-16-6"></span>[63] E. Masuko, Y. Hamada, A. Mizoguchi, H. Fukushi, N. Kuze and T. Sakaizumi, J. Mol. Spectrosc., 2009, 253, 77–87.
- <span id="page-16-7"></span>[64] J. S. Armstrong, Long-range forecasting: from crystal ball to computer, Wiley, New York, 2nd edn, 1985.
- <span id="page-16-8"></span>[65] S. Makridakis, Int. J. Forecast., 1993, 9, 527–529.
- <span id="page-16-9"></span>[66] S. Jachner, G. Boogaart and T. Petzoldt, J. Stat. Soft., 2007, 22, 1–30.
- <span id="page-16-10"></span>[67] M. Hunger and P. Döll, *Hydrol. Earth Syst. Sci.*, 2008, **12**, 841–861.
- <span id="page-16-11"></span>[68] S. Yu, B. Eder, R. Dennis, S.-H. Chu and S. E. Schwartz, Atmos. Sci. Lett., 2006, 7, 26–34.