# Pseudo-tetrahedral vs pseudo-octahedral Er'II single molecule magnets and the disruptive role of coordinated TEMPO radicals 

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## Experimental section

## General synthesis considerations

All manipulations were performed under inert gas atmosphere inside an argon filled glovebox (Inert PureLab HE). All solvents used in this study (HPLC grade) were dried under argon using Inert PureSolv EN7 solvent purification system. Anhydrous $\mathrm{ErCl}_{3}$ (at least 99.9\%) was purchased from Alfa Aesar and used as received. $\mathrm{Li}_{3}\left(\mathrm{~N}_{3} \mathrm{~N}\right)$ was prepared according to the previously published method. ${ }^{1}$ Sodium 2,4,6-tri-tert-butyl-phenolate (NaTTBP•Et ${ }_{2}$ O) and sodium 2,6-di-tert-butyl-4methylphenolate ( $\mathrm{NaBHT} \cdot \mathrm{Et}_{2} \mathrm{O}$ ) were prepared by slowly adding dry sodium hydride to the $\mathrm{Et}_{2} \mathrm{O}$ solution of the respective phenol (Sigma Aldrich, $98 \%$ ) inside the glovebox. See below for a detailed procedure. $\mathrm{NaTTBP} \cdot \mathrm{Et}_{2} \mathrm{O}$ obtained in this way is identical with that obtained using a slightly different method reported in the literature. ${ }^{2}$

## Preparation of sodium 2,4,6-tri-tert-butyl-phenolate ( $\mathrm{NaTTBP}^{2}$-Et ${ }_{2} \mathrm{O}$ )

2,4,6-tri-tert-butyl-phenol ( $10.7 \mathrm{~g}, 40 \mathrm{mmol}$ ) was dissolved in 40 ml of $\mathrm{dry}^{\mathrm{Et}} \mathrm{F}_{2} \mathrm{O}$ (bright yellow solution) and heated to boiling. $\mathrm{NaH}(0.912 \mathrm{~g}, 38 \mathrm{mmol})$ was added in small portions while stirring. Each portion causes vigorous gas evolution. Stirring was continued for another 1.5 h after the last portion of NaH was added. The off-white crystalline precipitate was collected by filtration (the filtrate is yellow) and washed with three portions of $\mathrm{Et}_{2} \mathrm{O}$. Yield 10.95 g ( $80 \%$ ). NaTTBP•Et $\mathrm{E}_{2} \mathrm{O}$ obtained in this way is identical with that obtained using a slightly different method reported in the literature. ${ }^{2}$

## Preparation of [Er"'(TTBP) ${ }_{3}$ (THF)] (1)

$\mathrm{NaTTBP} \cdot \mathrm{Et}_{2} \mathrm{O}(3.89 \mathrm{~g}, 10.9 \mathrm{mmol})$ and $\mathrm{ErCl}_{3}(1.03 \mathrm{~g}, 3.65 \mathrm{mmol})$ were suspended in 40 ml of THF. The pink reaction mixture was stirred for 3 days and then the solvent was removed in vacuo. The pink residue was suspended in 50 ml of pentane, stirred for 30 min and then left to settle. Then pink solution was decanted and filtered through a P4 fritted funnel (gravitational filtering). The extraction/filtration was repeated two more times and the pink clear filtrates were combined and left at $-40^{\circ} \mathrm{C}$ overnight. Large pink crystals were collected from the mixture. Yield: $2.25 \mathrm{~g}(61 \%)$. The identity and purity of the compound was confirmed by powder X-ray diffraction (PXRD) measurements (Figure S1a).

## Preparation of [Er"'(TTBP) $)_{3}($ TEMPO $\left.)\right]$ (2)

Pink solution of $1(0.213 \mathrm{~g}, 0.21 \mathrm{mmol})$ dissolved in 3.0 ml of $n$-pentane was mixed with an orange solution of TEMPO ( $0.183 \mathrm{~g}, 1.17 \mathrm{mmol}$ ) dissolved in 1.4 ml of the same solvent. The clear dark violet solution was left in an open vial for crystallization. Within 1-2 h the mixture volume was reduced to half of the initial value and violet crystals formed on the bottom. The crystals of $\mathbf{2}$ were filtered, washed with a quantity of cold $n$-pentane and dried inside the glovebox for 10 min. Yield: $0.040 \mathrm{~g}(16 \%)$. The identity and purity of the compound was confirmed by powder X-ray diffraction (PXRD) measurements (Figure S1b).

## Preparation of [Er"'(BHT) $\left.\mathbf{3}_{3}(\mathrm{THF})\right]$ ( 3 )

Compound $\mathbf{3}$ was obtained in an analogous way as $\mathbf{1}$ by reacting $\mathrm{NaBHT} \cdot \mathrm{Et}_{2} \mathrm{O}(7.60 \mathrm{~g}, 24 \mathrm{mmol})$ instead of TTBP with $\mathrm{ErCl}_{3}$ $(2.00 \mathrm{~g}, 7.3 \mathrm{mmol})$ in 100 ml of THF. Yield: $5.01 \mathrm{~g}(77 \%)$. The identity and purity of the compound was confirmed by powder X-ray diffraction (PXRD) measurements (Figure S1c).

## Preparation of $\left[\mathrm{Li}(\mathrm{THF})_{2}\right]_{2}\left[\mathrm{Er}^{11 \prime}\left(\mathrm{~N}_{3} \mathrm{~N}\right) \mathrm{Cl}_{2}\right]$ (4)

Solid $\mathrm{ErCl}_{3}(0.191 \mathrm{~g}, 0.70 \mathrm{mmol})$ was added in portions to the THF solution ( 6.8 ml ) of $\mathrm{Li}_{3}\left(\mathrm{~N}_{3} \mathrm{~N}\right)(0.267 \mathrm{~g}, 0.71 \mathrm{mmol})$. The pink turbid reaction mixture became almost clear after stirring overnight. The solution was evaporated to dryness and extracted three times with $n$-pentane ( 12 ml in total). The pink solution was filtered (P3 fritted glass funnel, gravitational), concentrated down to $3-4 \mathrm{ml}$ and left in the freezer at $-40^{\circ} \mathrm{C}$ overnight. Pink crystals were collected by filtration. Yield: 49 $\mathrm{mg}(8 \%)$. The identity and purity of the compound was confirmed by powder X-ray diffraction (PXRD) measurements (Figure S1d).

## Single crystal X-ray diffraction (SCXRD)

SCXRD experiments were performed for [Er $\left.{ }^{\prime \prime \prime \prime}(\mathrm{TTBP})_{3}(\mathrm{THF})\right]$ (1) and [Er"'I(TTBP) $)_{3}$ (TEMPO)] (2) using Rigaku XtaLAB SynergyS (HyPix) and for $\left[E r^{\prime \prime \prime}(\mathrm{BHT})_{3}(\mathrm{THF})\right]$ (3) and $\left[\mathrm{Li}(\mathrm{THF})_{2}\right]_{2}\left[E r^{\prime \prime \prime}\left(\mathrm{N}_{3} \mathrm{~N}\right) \mathrm{Cl}_{2}\right]$ (4) using Bruker D8 Quest Eco (Photon50) diffractometers (Mo K $\alpha$ radiation, PhotonJet-S microsource in the case of Rigaku and sealed tube Triumph ${ }^{\circledR}$ in the case of Bruker machine). Details of these measurements are gathered in Tables S1 and S2. Single crystals were transferred from the mother liquor into the Paratone-N oil to avoid decomposition and potential solvent loss in contact with ambient atmosphere and mounted using MiTeGen cryomounts. The measurements were performed first at 100 K (full data collection, Table S1) and then near room temperature (fast data collection to avoid crystal decomposition, Table S2) for each crystal. Data processing was performed using CrysAlisPro 1.171.40.67a or Apex3 suite of programs, respectively. The structures were solved using direct methods and refined anisotropically (weighted full-matrix least-squares on $F^{2} 3,4$ ). Hydrogen atoms were placed in the calculated positions and refined as riding on the parent atoms. Structural diagrams were prepared using Mercury 2020.2.0 software (CCDC). CCDC 2025204 ( $\mathbf{1}$ at 100 K ), 2027594 ( $\mathbf{2}$ at 100 K ), 2025203 ( $\mathbf{3}$ at 100 K ), 2025202 ( 4 at 100 K ), 2025205 ( 1 at 250 K ), 2027595 ( $\mathbf{2}$ at 250 K ), 2025200 ( $\mathbf{3}$ at 296 K ), 2025201 ( 4 at 296 K ) contain the supplementary crystallographic data for this paper which can be obtained free of charge from the Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/data request/cif.

## Powder X-ray diffraction (PXRD)

PXRD measurements were performed using Bruker D8 Advance Eco diffractometer equipped with CuK $\alpha$ radiation source and Lynxeye silicon strip detector. The samples were ground using an agate mortar inside the glovebox and loaded into 0.7 mm glass capillaries. The capillaries were broken in half inside the glovebox and their open end was sealed with silicon grease before they were moved to the PXRD instrument. The capillaries were mounted on the goniometer head using wax. The PXRD of each sample was collected in at least three runs in the $2-502 \vartheta$ range to exclude the possibility of the decomposition of the sample, if the silicon grease seal is not air tight. In each case there were no signs of decomposition within the experiment duration. The experimental PXRD patterns are presented in Figure S1 (colored lines) and compared against the simulated curves (gray lines) obtained from the near-room temperature SCXRD data (the simulated curves were exported using the relevant option of the Mercury 2020.2.0 software).

## Magnetic measurements

Magnetic susceptibility measurements were performed using a Quantum Design MPMS-3 Evercool magnetometer in the magnetic fields up to 7 T and 1.8-300 K temperature range. The samples (typically $25-30 \mathrm{mg}$ ) was loaded into the custommade Delrin sample holders ${ }^{5}$ inside the glovebox and closed tightly. The experimental data were corrected for the diamagnetism of the sample and the sample holder.

Table S1. Crystal structure solution and refinement parameters for [Er"'(TTBP) $\left.)_{3}(\mathrm{THF})\right]$ (1), [Er"'(TTBP) $\left.)_{3}(\mathrm{TEMPO})\right](\mathbf{2})$, $\left[\mathrm{Er}{ }^{\prime \prime \prime}(\mathrm{BHT})_{3}(\mathrm{THF})\right]$ (3) and $\left[\mathrm{Li}(\mathrm{THF})_{2}\right]_{2}\left[\mathrm{Er}^{\prime \prime \prime}\left(\mathrm{N}_{3} \mathrm{~N}\right) \mathrm{Cl}_{2}\right](4)$ at 100 K .

| Compound | [Er'II(TTBP) $\left.)_{3}(\mathrm{THF})\right]$ (1) | [Er ${ }^{\text {III }}$ (TTBP) ${ }_{3}$ (TEMPO) ${ }^{\text {(2) }}$ | [Er'I'(BHT) ${ }_{3}(\mathrm{THF}) \mathrm{l}^{(3)}$ | $\left[\mathrm{Li}(\mathrm{THF})_{2}\right]_{2}\left[\mathrm{Er}^{1 \prime \prime}\left(\mathrm{~N}_{3} \mathrm{~N}\right) \mathrm{Cl} 2_{2}\right](4)$ |
| :---: | :---: | :---: | :---: | :---: |
| CCDC number | 2025204 | 2027594 | 2025203 | 2025202 |
| Formula | $\mathrm{C}_{58} \mathrm{H}_{95} \mathrm{ErO}_{4}$ | $\mathrm{C}_{63} \mathrm{H}_{105} \mathrm{ErNO}_{4}$ | $\mathrm{C}_{98} \mathrm{H}_{154} \mathrm{Er}_{2} \mathrm{O}_{8}$ | $\mathrm{C}_{31} \mathrm{H}_{71} \mathrm{Cl}_{2} \mathrm{ELLi}_{2} \mathrm{~N}_{4} \mathrm{O}_{4} \mathrm{Si}_{3}$ |
| $M_{r} / \mathrm{g} \mathrm{mol}^{-1}$ | 1023.59 | 1107.73 | 1794.72 | 900.23 |
| T/K | 100(1) | 100(1) | 100(1) | 100(1) |
| Crystal system | triclinic | monoclinic | triclinic | monoclinic |
| Space group | P-1 | $P 2_{1} / \mathrm{n}$ | P-1 | $P 2_{1} / \mathrm{c}$ |
| $a / \AA$ | 10.3908(1) | 14.12440(10) | 14.483(4) | 22.9459(19) |
| b/A | 13.7538(2) | 27.2504(3) | 18.193(10) | 10.0173(8) |
| c/ $\AA$ | 21.5780(3) | 15.9837(2) | 19.256(7) | 18.5794(15) |
| $\alpha /{ }^{\circ}$ | 98.962(1) | 90 | 71.753(13) | 90 |
| b/ ${ }^{\circ}$ | 99.581(1) | 90.2240(10) | 86.585(10) | 89.998(2) |
| $v /{ }^{\circ}$ | 109.287(1) | 90 | 81.656(17) | 90 |
| $V /{ }^{3}$ | 2795.55(7) | 6152.01(11) | 4767.(3) | 4270.6(6) |
| Z | 2 | 4 | 2 | 4 |
| $\rho_{\text {calc }} / \mathrm{g} \mathrm{cm}^{-3}$ | 1.216 | 1.196 | 1.250 | 1.400 |
| $\mu / \mathrm{mm}^{-1}$ | 1.542 | 1.407 | 1.799 | 2.210 |
| F(000) | 1086 | 2360 | 1884 | 1868 |
| Crystal size / mm ${ }^{3}$ | $0.10 \times 0.04 \times 0.03$ | $0.12 \times 0.06 \times 0.04$ | $0.52 \times 0.40 \times 0.35$ | $0.39 \times 0.20 \times 0.13$ |
| Instrument | Rigaku Synergy S | Rigaku Synergy S | Bruker D8 Quest Eco | Bruker D8 Quest Eco |
| Radiation | Mo K $\alpha$ ( $\lambda=0.71073$ Å) | Mo K $\alpha$ ( $\lambda=0.71073$ Å) | Mo K $\alpha$ ( $\lambda=0.71073$ Å) | Mo K $\alpha$ ( $\lambda=0.71073$ Å) |
| $2 \theta$ range/ ${ }^{\circ}$ | 2.241-25.242 | 2.434-25.242 | 2.26-25.75 | 2.31-25.24 |
| Reflections collected | 58260 | 197493 | 41119 | 39038 |
| Independent reflections | 19074 | 22437 | 17791 | 7719 |
| $R_{\text {int }}$ | 0.0585 | 0.0423 | 0.0762 | 0.0877 |
| restrains/parameters | 0/595 | 0/653 | 48/1015 | 48/461 |
| $R\left[F_{o}>2 \sigma\left(F_{o}\right)\right]$ | 0.0475 | 0.0284 | 0.0499 | 0.0415 |
| $w R\left(F^{2}\right)$ | 0.1107 | 0.0603 | 0.0960 | 0.0832 |
| GOF on $F^{2}$ | 1.032 | 1.041 | 1.011 | 1.030 |
| $\Delta \rho_{\text {max }}, \Delta \rho_{\text {min }} / \mathrm{e} \AA^{-3}$ | 4.648/-2.521 | 0.991/-0.907 | 1.268/-1.077 | 1.568/-1.247 |
| Completeness/\% | 99.9 | 99.9 | 97.4 | 99.9 |

Table S2. Crystal structure solution and refinement parameters for $\left[\mathrm{Erll}^{\prime \prime \prime}(\mathrm{TTBP})_{3}(\mathrm{THF})\right](1)$, [Er"II(TTBP) $\left.{ }_{3}(\mathrm{THF})\right](2)$, $\left[\mathrm{Er}{ }^{\prime \prime \prime}(\mathrm{BHT})_{3}(\mathrm{TEMPO})\right](3)$ and $\left[\mathrm{Li}(\mathrm{THF})_{2}\right]_{2}\left[\mathrm{Er}^{\prime \prime \prime}\left(\mathrm{N}_{3} \mathrm{~N}\right) \mathrm{Cl}_{2}\right](4)$ near room temperature ( 250 or 296 K ) - structural models obtained for the sole purpose of comparison with the experimental powder X-ray diffraction patterns (see Figure S1 below).

| Compound | [Er $\left.{ }^{\prime \prime \prime}(\mathrm{TTBP})_{3}(\mathrm{THF})\right]$ (1) | [Er"II(TTBP) $\left.{ }_{3}(\mathrm{TEMPO})\right](2)$ | [Er $\left.{ }^{\text {II' }}(\mathrm{BHT})_{3}(\mathrm{THF})\right](3)$ | $\left[\mathrm{Li}(\mathrm{THF})_{2}\right]_{2}\left[\mathrm{Er}^{\prime \prime \prime}\left(\mathrm{N}_{3} \mathrm{~N}\right) \mathrm{Cl}_{2}\right](4)$ |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{array}{ll} \hline \text { CCDC } & \text { deposition } \\ \text { number } \end{array}$ | 2025205 | 2027595 | 2025200 | 2025201 |
| Formula | $\mathrm{C}_{58} \mathrm{H}_{95} \mathrm{ErO}_{4}$ | $\mathrm{C}_{63} \mathrm{H}_{105} \mathrm{ErNO}_{4}$ | $\mathrm{C}_{49} \mathrm{H}_{77} \mathrm{ErO}_{4}$ | $\mathrm{C}_{31} \mathrm{H}_{71} \mathrm{Cl}_{2} \mathrm{ELLi}_{2} \mathrm{~N}_{4} \mathrm{O}_{4} \mathrm{Si}_{3}$ |
| $M_{r} / \mathrm{g} \mathrm{mol}^{-1}$ | 1023.59 | 1107.73 | 897.36 | 900.22 |
| T/K | 250(1) | 250(1) | 296(1) | 296(1) |
| Crystal system | triclinic | monoclinic | triclinic | monoclinic |
| Space group | P-1 | $P 2_{1} / \mathrm{n}$ | P-1 | $P 2_{1} / \mathrm{n}$ |
| $a / \AA$ | 10.4504(2) | 14.2971(2) | 10.7671(13) | 10.2048(16) |
| b/A | 13.8690(3) | 27.4460(5) | 12.4225(14) | 18.930(3) |
| c/ $\AA$ | 21.7789(5) | 16.0132(2) | 19.496(3) | 23.147(4) |
| $\alpha /{ }^{\circ}$ | 98.439(2) | 90 | 74.976(5) | 90 |
| $6 /^{\circ}$ | 99.276(2) | 90.190(2) | 77.692(5) | 90.010(5) |
| V/ ${ }^{\circ}$ | 108.951(2) | 90 | 76.963(4) | 90 |
| $V /{ }^{3}$ | 2878.61(11) | 6283.52(16) | 2420.3(5) | 4471.3(11) |
| Z | 2 | 4 | 2 | 4 |
| $\rho_{\text {calc }} / \mathrm{g} \mathrm{cm}^{-3}$ | 1.181 | 1.171 | 1.231 | 1.337 |
| $\mu / \mathrm{mm}^{-1}$ | 1.498 | 1.378 | 1.772 | 2.111 |
| $F(000)$ | 1086 | 2360 | 942 | 1868 |
| Crystal size / mm ${ }^{3}$ | $0.10 \times 0.04 \times 0.03$ | $0.12 \times 0.06 \times 0.04$ | $0.52 \times 0.40 \times 0.35$ | $0.39 \times 0.20 \times 0.13$ |
| Instrument | Rigaku Synergy S | Rigaku Synergy S | Bruker D8 Quest Eco | Bruker D8 Quest Eco |
| Radiation | Mo K $\alpha(\lambda=0.71073$ Å) | Mo K ${ }^{(\lambda}=0.71073$ Å) | Mo K $\alpha$ ( $\lambda=0.71073$ Å) | Mo K $\alpha(\lambda=0.71073$ Å) |
| $2 \theta$ range/ ${ }^{\circ}$ | 2.230-25.242 | 2.416-25.242 | 2.26-25.75 | 2.267-17.982 |
| Reflections collected | 68249 | 199852 | 7981 | 6737 |
| Independent reflections | 19831 | 22878 | 5102 | 2742 |
| $R_{\text {int }}$ | 0.0657 | 0.0670 | 0.0626 | 0.0831 |
| restrains/parameters | 0/595 | 0/653 | 72/570 | 60/461 |
| $R\left[F_{o}>2 \sigma\left(F_{o}\right)\right]$ | 0.0518 | 0.0400 | 0.0609 | 0.0460 |
| $w R\left(F^{2}\right)$ | 0.1252 | 0.0885 | 0.1397 | 0.1034 |
| GOF on $F^{2}$ | 1.036 | 1.036 | 1.072 | 1.055 |
| $\Delta \rho_{\text {max }}, \Delta \rho_{\text {min }} / \mathrm{e} \AA^{-3}$ | 2.641/-1.686 | 1.114/-0.801 | 0.863/-0.861 | 0.648/-0.445 |
| Completeness/\% | 99.9 | 99.9 | 89.0 | 89.5 |

Table S3. Results of the CShM analysis performed using SHAPE software for all Er'I' central ions in 1-4. Values closer to zero indicate better match with the reference geometry.

|  | geometry | 1 | 2 | 3 (A) | 3 (B) | 4 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | D6h Hexagon |  |  |  |  | 33.142 |
|  | C5v Pentagonal pyramid |  |  |  |  | 16.834 |
|  | Oh Octahedron |  |  |  |  | 4.980 |
|  | D3h Trigonal prism |  |  |  |  | 9.956 |
|  | C5v Johnson pentagonal pyramid |  |  |  |  | 20.017 |
|  | D4h Square | 27.369 | 30.309 | 29.819 | 27.941 |  |
|  | Td Tetrahedron | 1.029 | 0.746 | 0.893 | 1.036 |  |
|  | C2v Seesaw | 7.722 | 5.846 | 8.583 | 7.857 |  |
|  | C3v Vacant trigonal bipyramid | 2.921 | 2.182 | 2.803 | 2.725 |  |



Figure S1. Experimental PXRD patterns for $\mathbf{1}$ (red line) (a), $\mathbf{2}$ (violet line) (b), $\mathbf{3}$ (blue line) (c) and $\mathbf{4}$ (green line) (d) recorded at room temperature. The grey lines represent the PXRD patterns simulated from room temperature (or 250 K ) scXRD data.


Figure S2. Packing diagrams of $\mathbf{1}$ (a), 2 (b), $\mathbf{3}$ (c) and $\mathbf{4}$ (d) highlighting the ' $\mathrm{ErO}_{4}$ ' pseudo-tetrahedra (1-3) and ' $\mathrm{ErN}_{4} \mathrm{Cl}_{2}$ ' pseudo-octahedra (4). Organic ligands as well as THF molecules and $\mathrm{Li}^{+}$cations are drawn as sticks.


Figure S3. DC magnetic properties for compound 1: $\chi_{M} T(T)$ recorded at 100 Oe (main plot) and $M(H)$ recorded at 1.8 K (inset). Black points - experimental data, red lines - CASSCF ab initio calculations.


Figure S4. DC magnetic properties for compound 2: $\chi_{M} T(T)$ recorded at 100 Oe (main plot) and $M(H)$ recorded at 1.8 K (inset).


Figure S5. DC magnetic properties for compound 3: $\chi_{M} T(T)$ recorded at 100 Oe (main plot) and $M(H)$ recorded at 1.8 K (inset). Black points - experimental data, blue lines - CASSCF ab initio calculations.


Figure S6. DC magnetic properties for compound 4: $\chi_{\mathrm{M}} T(T)$ recorded at 100 Oe (main plot) and $M(H)$ recorded at 1.8 K (inset). Black points - experimental data, green lines - CASSCF ab initio calculations.


Figure S7. Comparison of the DC magnetic properties for compounds 1 and 2: $\chi_{M} T(T)$ recorded at 100 Oe (main plot) and $M(H)$ recorded at 1.8 K (inset). The difference between both curves correspond to the presence of the coordinated TEMPO radical ( $S=1 / 2$ ) in compound 2. Experimental $\Delta \chi_{\mathrm{M}} T$ at 300 K is $0.4 \mathrm{~cm}^{3} \cdot \mathrm{~K} \cdot \mathrm{~mol}^{-1}$ and $\Delta M$ at 70 kOe amounts to $0.77 \mu_{\mathrm{B}}$ while the corresponding values for TEMPO are $0.375 \mathrm{~cm}^{3} \cdot \mathrm{~K} \cdot \mathrm{~mol}^{-1}$ and $1.00 \mu_{\mathrm{B}}$.

## AC magnetic properties of compound 1 (Section AC1)

Frequency dependence of the out-of-phase magnetic susceptibility ( $\chi^{\prime \prime}$ ) of $\mathbf{1}$ at 1.8 K shows two maxima in the $1-1000 \mathrm{~Hz}$ range under 100 and 200 Oe applied DC field, but only the low frequency one is analysed further (Table S4 and Figure S8a). This maximum shifts with magnetic field from around 20 Hz at 100 Oe to around 9 Hz at 1000 Oe and then again to higher frequencies above 3000 Oe. Contributions from the QTM and the direct process have been identified in the lowand high-field ranges of the $\tau(H)$ dependence (Figure 3a), respectively. Fitting of the $\tau(H)$ dependence using Eq. 2 resulted in best-fit parameters $A_{1}, A_{2}, A_{3}$ and $A_{4}$ which were used to calculate parameters for the temperature dependence: $A=$ $2298 \mathrm{~s}^{-1}$ and $B_{2}=0 \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-1}$ at 0 Oe applied field and $A=0.46 \mathrm{~s}^{-1}$ and $B_{2}=3.7 \cdot 10^{-3} \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-1}$ at 1000 Oe applied field in Eq. 3. The temperature dependence of the AC magnetic susceptibility was measured at 0 Oe (Figure S 8 b ) and at 1000 Oe (optimum field; Figure S 8 c ) in order to analyze both temperature dependences simultaneously. The $\chi$ " maximum is visible up to 10.1 K at 0 Oe and up to 11.2 K at 1000 Oe in the $1-1000 \mathrm{~Hz}$ window. At 0 Oe the $\chi^{\prime \prime}(v)$ maximum does not move with temperature up to 6.8 K (Figure S8b). Simultaneous fitting of the temperature dependences of $\ln (\tau)$ at 0 and 1000 Oe using Eq. 3 (see main text) with $A$ and $B_{2}$ fixed (based on the $\tau(H)$ dependence fitting), shows a significant contribution of the Raman relaxation process ( $C=1.3(1) \cdot 10^{-4} \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-2}$ assuming fixed exponent $n=7$ which is not typical for Kramers ions and might indicate a more complex Raman relaxation mechanism involving a vibronic barrier; 6,7 note that reasonable fits could only be obtained when $n$ was close to 7 ) and Orbach relaxation process ( $\tau_{0}=1.4(5) \cdot 10^{-6} \mathrm{~s}$; $\left.U_{\text {eff }} / \mathrm{k}_{\mathrm{B}}=63(4) \mathrm{K}\right)$ (Figure 3 b ). All parameters obtained by fitting the respective relaxation rates vs magnetic field and temperature using Eqs 2 and 3 are collected in Table S5 and Table 3 in the main text.

Table S4. Values of $\alpha$ and $\tau$ parameters (uncertainties provided in parentheses) obtained from generalized Debye model fitting of the $\tau(v)$ dependences for 1 at 6.0 K measured in the $0-18000$ Oe magnetic field range (columns 1-5), at 0 Oe in the 1.8-10.1 K temperature range (columns $6-10$ ) and at 1000 Oe in the $4.2-11.7 \mathrm{~K}$ temperature range (columns 11-15).

| Figure S8a |  |  |  |  | Figure S8b |  |  |  |  | Figure S8c |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $T=6.0 \mathrm{~K}$ |  |  |  |  | $H=0 \mathrm{Oe}$ |  |  |  |  | $H=1000 \mathrm{Oe}$ |  |  |  |  |
| H (Oe) | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\alpha$ | $\begin{gathered} \tau \\ (\mathrm{s}) \end{gathered}$ | $\begin{gathered} T \\ (\mathrm{~K}) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\alpha$ | $\tau(\mathrm{s})$ | $\begin{gathered} \hline T \\ (K) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\alpha$ | $\tau(\mathrm{s})$ |
| 0 | 0.42(1) | 1.68(1) | 0.104(7) | 4.25(4) $10^{-4}$ | 1.8 | 1.41(8) | 5.33(4) | 0.094(20) | 6.13(22) $10^{-4}$ | 4.2 | 0.342(1) | 2.42(1) | 0.042(3) | $3.26(3) \cdot 10^{-1}$ |
| 100 | 1.14(2) | 1.66(1) | 0.024(20) | $1.07(3) \cdot 10^{-2}$ | 2.0 | 1.26(7) | 4.77(3) | 0.096(19) | $5.97(20) \cdot 10^{-4}$ | 4.4 | 0.331(1) | 2.30(1) | 0.035(2) | $2.16(1) \cdot 10^{-1}$ |
| 200 | 0.77(1) | 1.67(1) | 0.016(7) | $1.32(1) \cdot 10^{-2}$ | 2.2 | 1.16(6) | 4.35(3) | 0.097(17) | $5.86(19) \cdot 10^{-4}$ | 4.6 | 0.321(1) | 2.189(6) | 0.028(2) | $1.47(1) \cdot 10^{-1}$ |
| 300 | 0.45(1) | 1.71(2) | 0.128(17) | $1.34(4) \cdot 10^{-2}$ | 2.4 | 1.08(5) | 4.01(2) | 0.097(16) | $5.76(17) \cdot 10^{-4}$ | 4.8 | 0.313(1) | 2.096(4) | 0.024(2) | $1.03(1) \cdot 10^{-1}$ |
| 400 | 0.37(1) | 1.68(1) | 0.064(8) | $1.54(2) \cdot 10^{-2}$ | 2.6 | 0.99(4) | 3.71(2) | 0.102(14) | $5.57(15) \cdot 10^{-4}$ | 5.0 | 0.304(1) | 2.012(4) | 0.022(2) | $7.37(2) \cdot 10^{-2}$ |
| 500 | 0.324(3) | 1.66(1) | 0.038(5) | $1.65(1) \cdot 10^{-2}$ | 2.9 | 0.87(3) | 3.35(1) | 0.108(10) | $5.29(11) \cdot 10^{-4}$ | 5.2 | 0.295(1) | 1.935(3) | 0.019(2) | $5.38(2) \cdot 10^{-2}$ |
| 600 | 0.299(2) | 1.663(3) | 0.027(3) | 1.70 (1) $10^{-2}$ | 3.2 | 0.78(2) | 3.058(7) | 0.113(7) | $5.03(7) \cdot 10^{-4}$ | 5.4 | 0.287(1) | 1.865(2) | 0.018(2) | $3.99(1) \cdot 10^{-2}$ |
| 700 | 0.282(3) | 1.652(2) | 0.013(5) | $1.75(1) \cdot 10^{-2}$ | 3.5 | 0.70(1) | 2.809(5) | 0.119(6) | $4.83(6) \cdot 10^{-4}$ | 5.6 | 0.279(1) | 1.803(2) | 0.018(2) | $3.01(1) \cdot 10^{-2}$ |
| 1000 | 0.266(1) | 1.648(2) | 0.014(2) | $1.76(1) \cdot 10^{-2}$ | 3.8 | 0.64(1) | 2.597(4) | 0.121(5) | 4.70(5) $10^{-4}$ | 5.9 | 0.269(1) | 1.713(2) | 0.015(1) | $2.00(1) \cdot 10^{-2}$ |
| 1500 | 0.258(1) | 1.631(3) | 0.011(2) | $1.75(1) \cdot 10^{-2}$ | 4.1 | 0.60(1) | 2.415(3) | 0.122(4) | $4.63(4) \cdot 10^{-4}$ | 6.2 | 0.261(1) | 1.632(2) | 0.013(1) | $1.37(1) \cdot 10^{-2}$ |
| 2000 | 0.255(2) | 1.614(3) | 0.013(3) | $1.74(1) \cdot 10^{-2}$ | 4.4 | 0.56(1) | 2.255(3) | 0.123(5) | $4.56(5) \cdot 10^{-4}$ | 6.5 | 0.252(1) | 1.559(1) | 0.010(2) | $9.60(2) \cdot 10^{-3}$ |
| 2500 | 0.258(3) | 1.588(6) | 0.011(5) | $1.72(1) \cdot 10^{-2}$ | 4.7 | 0.53(1) | 2.116(3) | 0.122(5) | $4.51(5) \cdot 10^{-4}$ | 6.8 | 0.245(1) | 1.494(1) | 0.009(2) | $6.86(2) \cdot 10^{-3}$ |
| 3000 | 0.249(1) | 1.562(2) | 0.018(2) | $1.68(1) \cdot 10^{-2}$ | 5.0 | 0.50(1) | 1.993(3) | 0.121(5) | $4.47(5) \cdot 10^{-4}$ | 7.1 | 0.237(1) | 1.434(1) | 0.008(2) | $4.97(1) \cdot 10^{-3}$ |
| 3500 | 0.245(2) | 1.533(3) | 0.025(3) | $1.63(1) \cdot 10^{-2}$ | 5.3 | 0.48(1) | 1.884(3) | 0.119(5) | $4.41(5) \cdot 10^{-4}$ | 7.4 | 0.231(1) | 1.378(1) | 0.007(2) | $3.66(1) \cdot 10^{-3}$ |
| 4000 | 0.243(1) | 1.500(2) | 0.027(2) | $1.56(1) \cdot 10^{-2}$ | 5.6 | 0.46(1) | $1.787(3)$ | 0.117(5) | $4.34(5) \cdot 10^{-4}$ | 7.7 | 0.225(1) | 1.327(1) | 0.007(2) | $2.73(1) \cdot 10^{-3}$ |
| 4500 | 0.240(1) | 1.462(2) | 0.030(2) | $1.47(1) \cdot 10^{-2}$ | 5.9 | 0.45(1) | 1.698(3) | 0.111(5) | $4.30(5) \cdot 10^{-4}$ | 8.0 | 0.219(2) | 1.280(1) | 0.006(2) | $2.06(1) \cdot 10^{-3}$ |
| 5000 | 0.235(2) | 1.424(3) | 0.040(3) | $1.38(1) \cdot 10^{-2}$ | 6.2 | 0.43(1) | 1.618(2) | 0.105(5) | $4.21(5) \cdot 10^{-4}$ | 8.3 | 0.214(2) | 1.235(1) | 0.005(2) | $1.57(1) \cdot 10^{-3}$ |
| 5500 | 0.231(1) | 1.385(2) | 0.044(3) | $1.27(1) \cdot 10^{-2}$ | 6.5 | 0.42(1) | 1.545(2) | 0.098(5) | 4.10(4) $10^{-4}$ | 8.6 | 0.211(2) | 1.194(1) | 0.002(2) | $1.21(1) \cdot 10^{-3}$ |
| 6000 | 0.226(2) | 1.342(3) | 0.054(3) | $1.16(1) \cdot 10^{-2}$ | 6.8 | 0.40(1) | 1.479(2) | 0.091(5) | $3.95(4) \cdot 10^{-4}$ | 8.9 | 0.210(5) | 1.156(2) | 0.000(6) | $9.50(8) \cdot 10^{-4}$ |
| 7000 | 0.215(2) | 1.253(3) | 0.075(4) | $9.32(6) \cdot 10^{-3}$ | 7.1 | 0.395(7) | 1.418(2) | 0.082(5) | $3.79(4) \cdot 10^{-4}$ | 9.2 | 0.203(3) | 1.119(1) | 0.000(3) | $7.45(4) \cdot 10^{-4}$ |
| 8000 | 0.202(3) | 1.163(4) | 0.101(5) | $7.23(7) \cdot 10^{-3}$ | 7.4 | 0.40(1) | 1.365(4) | 0.077(10) | $3.63(7) \cdot 10^{-4}$ | 9.7 | 0.206(4) | 1.064(1) | 0.000(4) | $5.14(3) \cdot 10^{-4}$ |
| 9000 | 0.188(4) | 1.073(4) | 0.130(7) | $5.48(7) \cdot 10^{-3}$ | 7.7 | 0.380(7) | 1.309(2) | 0.057(5) | $3.44(4) \cdot 10^{-4}$ | 10.2 | 0.212(7) | 1.014(2) | 0.000(7) | $3.63(4) \cdot 10^{-4}$ |
| 10000 | 0.173(5) | 0.986(4) | 0.161(9) | $4.09(7) \cdot 10^{-3}$ | 8.0 | 0.373(6) | 1.261(2) | 0.045(5) | $3.22(3) \cdot 10^{-4}$ | 10.7 | 0.22(1) | 0.969(2) | 0.000(9) | $2.62(4) \cdot 10^{-4}$ |
| 11000 | 0.157(6) | 0.902(5) | 0.191(11) | $3.02(7) \cdot 10^{-3}$ | 8.3 | 0.367(7) | 1.217(2) | 0.033(5) | $2.98(3) \cdot 10^{-4}$ | 11.2 | 0.25(2) | 0.928(2) | 0.000(13) | $1.95(6) \cdot 10^{-4}$ |
| 12000 | 0.140(7) | 0.826(5) | 0.223(13) | $2.24(6) \cdot 10^{-3}$ | 8.6 | 0.362(7) | 1.175(1) | 0.021(5) | $2.75(3) \cdot 10^{-4}$ | 11.7 | 0.29(2) | 0.891(2) | 0.000(18) | $1.53(7) \cdot 10^{-4}$ |
| 13000 | 0.14(1) | 0.774(6) | 0.252(18) | $1.69(7) \cdot 10^{-3}$ | 8.9 | 0.359(7) | 1.137(1) | 0.009(6) | $2.51(10) \cdot 10^{-4}$ |  |  |  |  |  |
| 14000 | 0.10(1) | 0.688(5) | 0.284(16) | $1.22(5) \cdot 10^{-3}$ | 9.2 | 0.34(2) | 1.099(4) | 0.000(19) | $2.32(3) \cdot 10^{-4}$ |  |  |  |  |  |
| 15000 | 0.08(1) | 0.633(5) | 0.318(21) | $9.11(6) \cdot 10^{-4}$ | 9.5 | 0.35(1) | 1.068(1) | 0.000(7) | $2.05(3) \cdot 10^{-4}$ |  |  |  |  |  |
| 16000 | 0.06(2) | 0.581(5) | 0.350(23) | $6.72(5) \cdot 10^{-4}$ | 9.8 | 0.34(1) | 1.037(1) | 0.000(8) | $1.82(4) \cdot 10^{-4}$ |  |  |  |  |  |
| 17000 | 0.05(2) | 0.527(5) | 0.369(24) | $4.96(5) \cdot 10^{-4}$ | 10.1 | 0.34(1) | 1.008(1) | 0.000(9) | $1.62(4) \cdot 10^{-4}$ |  |  |  |  |  |
| 18000 | 0.02(3) | 0.488(5) | 0.419(28) | $3.45(5) \cdot 10^{-4}$ |  |  |  |  |  |  |  |  |  |  |



Figure S8. In-phase ( $\chi^{\prime}$ ) and out-of-phase ( $\chi^{\prime \prime}$ ) AC susceptibilities for $\mathbf{1}$ at 6.0 K measured in the 0-18000 Oe magnetic field range (a), at 0 Oe in the 1.8-10.1 K temperature range (b) and at 1000 Oe in the 4.1-11.7 K temperature range (c). Values of $\alpha$ and $\tau$ parameters are gathered in Table S4. The solid lines are the best fits to generalized Debye model.

Table S5. Values of the best-fit parameters obtained by fitting the $\tau^{1}(H)$ using Eq. 2 (top part of the table) and the $\tau(T)$ using Eq. 3 (bottom part of the table) for compound 1. The relevant fits are shown in Figure 3a and 3b in the main text, respectively. This table is also included in Table 3 in the main text.
$\tau^{-1}(H)=A_{1} /\left(1+A_{2} H^{2}\right)+B_{1} H^{4}+D$
(Eq. 2)
$\ln \tau\left(T^{-1}\right)=\ln \left[\left(A+B_{2} T+C T^{n}+\tau_{0}^{-1} \exp \left(-U_{\text {eff }} / k_{B} T\right)\right)^{-1}\right]$

| $T(\mathrm{~K})$ | 6.0 |  |
| :--- | :---: | :---: |
| range (Hz) | $1-1000$ |  |
| field range (Oe) | $0-18000$ |  |
| $A_{1}\left(\mathrm{~s}^{-1}\right)$ | $2298(118)$ |  |
| $A_{2}\left(\mathrm{Oe}^{-2}\right)$ | $5(3) \cdot 10^{-3}$ |  |
| $B_{1}\left(\mathrm{Oe}^{-4}\right)$ | $2.23(6) \cdot 10^{-14}$ |  |
| $D\left(\mathrm{~s}^{-1}\right)$ | $56(5)$ |  |
| $R^{2}$ | 0.98511 | 1000 |
| $H_{\mathrm{DC}}(\mathrm{Oe})$ | 0 | $1-1000$ |
| range $(\mathrm{Hz})$ | $1-1000$ | $1.8-10.1$ |
| temp. range (K) | $1.8-10.1$ | 0.46 (fixed) |
| $A\left(\mathrm{~s}^{-1}\right)$ | 2298 (fixed) |  |
| $B_{2}\left(\mathrm{~s}^{-1} \mathrm{~K}^{-1}\right)$ | 0 (fixed) |  |
| $C\left(\mathrm{~s}^{-1} \mathrm{~K}^{-n}\right)$ |  |  |
| $n$ | 7 (fixed) | $1.3(1) \cdot 10^{-4}$ |
| $\tau_{0}(\mathrm{~s})$ |  |  |
| $\left(U_{\text {eff }} / \mathrm{k}_{\mathrm{B}}\right)(\mathrm{K})$ |  | $7.4(5) \cdot 10^{-3}$ (fixed) |
| $R^{2}$ | $63(4)$ |  |
|  |  | 0.99843 |

## AC magnetic properties of compound 2 (Section AC2)

The AC magnetic susceptibility for compound $\mathbf{2}$ does not show any frequency dependence as evidenced by the $\chi^{\prime}, \chi^{\prime \prime}(T)$ plots performed at 7, 70 and 700 Hz (Figure S9).


Figure S9. Temperature dependence of the in-phase ( $\chi^{\prime}$ ) and out-of-phase ( $\chi^{\prime \prime}$ ) AC susceptibilities using 7,70 and 700 Hz frequency for compound $\mathbf{2}$ measured at 0 (a) and $1000 \mathrm{Oe}(\mathrm{b})$ external magnetic field.

## AC magnetic properties of compound 3 (Section AC3)

Similarly to 1 , the frequency dependence of the out-of-phase magnetic susceptibility ( $\chi^{\prime \prime}$ ) of $\mathbf{3}$ at 1.8 K shows two maxima in the $1-1000 \mathrm{~Hz}$ range under 100 and 200 Oe applied DC field, but only the low frequency one is analysed further (Table S6 and Figure S10a). This maximum shifts with magnetic field from around 15 Hz at 100 Oe to around 6 Hz at 1500 Oe and then again to higher frequencies above 3500 Oe. Contributions from the QTM and the direct process have been identified in the low- and high-field ranges of the $\tau(H)$ dependence (Figure 3c), respectively. Fitting of the $\tau(H)$ dependence using Eq. 2 resulted in best-fit parameters $A_{1}, A_{2}, A_{3}$ and $A_{4}$ which were used to calculate parameters for the temperature dependence: $A=3096 \mathrm{~s}^{-1}$ and $B_{2}=0 \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-1}$ at 0 Oe applied field and $A=0.20 \mathrm{~s}^{-1}$ and $B_{2}=0.3 \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-1}$ at 1000 Oe applied field in Eq. 3. The temperature dependence of the AC magnetic susceptibility was measured at 0 Oe (Figure S10b) and at 1500 Oe (optimum field; Figure S10c) in order to analyze both temperature dependences simultaneously. The $\chi^{\prime \prime}$ maximum is visible up to 6.5 K at 0 Oe and up to 8.0 K at 1500 Oe in the $1-1000 \mathrm{~Hz}$ window. At 0 Oe the $\chi^{\prime \prime}(v)$ maximum does not move with temperature up to 5.8 K (Figure S 10 b ). Simultaneous fitting of the temperature dependences of $\ln (\tau)$ at 0 and 1500 Oe using Eq. 3 (see main text) with $A$ and $B_{2}$ fixed (based on the $\tau(H)$ dependence fitting), shows a significant contribution of the Raman relaxation process $\left(C=1.5(1) \cdot 10^{-3} \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-2}\right.$, assuming fixed exponent $n=7$ which is not typical for Kramers ions and might indicate a more complex Raman relaxation mechanism involving a vibronic barrier; ${ }^{6,7}$ note that reasonable fits could only be obtained when $n$ was close to 7 ) and Orbach relaxation process ( $\tau_{0}=7(3) \cdot 10^{-7} \mathrm{~s}$; $U_{\text {eff }} / \mathrm{k}_{\mathrm{B}}=48(3)$ K) (Figure 3d). All parameters obtained by fitting the respective relaxation rates vs magnetic field and temperature using Eqs 2 and 3 are collected in Table S7 and in Table 3 in the main text.

Table S6. Values of $\alpha$ and $\tau$ parameters (uncertainties provided in parentheses) obtained from generalized Debye model fitting of the $\tau(v)$ dependences for 3 at 4.0 K measured in the $0-16000$ Oe magnetic field range (columns 1-5), at 0 Oe in the 1.8-7.7 K temperature range (columns $6-10$ ) and at 1000 Oe in the $3-8.0 \mathrm{~K}$ temperature range (columns 1-15).

| Figure S10a |  |  |  |  | Figure S10b |  |  |  |  | Figure S10c |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $T=4.0 \mathrm{~K}$ |  |  |  |  | $H=0 \mathrm{Oe}$ |  |  |  |  | $H=1500 \mathrm{Oe}$ |  |  |  |  |
| H (Oe) | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\alpha$ | $\begin{gathered} \tau \\ (\mathrm{s}) \\ \hline \end{gathered}$ | $\begin{gathered} \mathbf{T} \\ (\mathrm{K}) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\alpha$ | $\begin{gathered} \tau \\ (\mathrm{s}) \\ \hline \end{gathered}$ | $\boldsymbol{T}$ (K) | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \\ \hline \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \end{gathered}$ | $\alpha$ | $\begin{gathered} \tau \\ (\mathrm{s}) \\ \hline \end{gathered}$ |
| 0 | 0.52(1) | 2.456(3) | 0.080(5) | 3.19(3).10-4 | 3.2 | 0.63(2) | 5.52(1) | 0.092(5) | $3.27(4) \cdot 10^{-4}$ | 3.0 | 0.116(1) | 3.07(4) | 0.090(4) | 2.64(5). $10^{-1}$ |
| 100 | 1.74(2) | 2.460(6) | 0.090(50) | 1.44(13) $10^{-2}$ | 3.5 | 0.58(2) | 4.91(1) | 0.091(6) | $3.24(4) \cdot 10^{-4}$ | 3.2 | 0.109(1) | 2.87(2) | 0.089(3) | $1.57(1) \cdot 10^{-1}$ |
| 200 | 1.10(1) | 2.462(6) | 0.090(12) | $1.69(3) \cdot 10^{-2}$ | 3.8 | 0.54(2) | 4.481(6) | 0.087(6) | $3.22(4) \cdot 10^{-4}$ | 3.4 | 0.099(2) | 2.75(1) | 0.097(3) | 9.86(8) $\cdot 10^{-2}$ |
| 300 | 0.64(1) | 2.471(6) | 0.090(6) | 2.01(2). $10^{-2}$ | 4.1 | 0.51(2) | 4.119(6) | 0.083(6) | $3.16(4) \cdot 10^{-4}$ | 3.6 | 0.094(3) | 2.65(1) | 0.105(3) | 6.35(5). $10^{-2}$ |
| 400 | 0.36(1) | 2.479(6) | 0.094(3) | $2.28(1) \cdot 10^{-2}$ | 4.4 | 0.48(2) | 3.808(5) | 0.077(6) | $3.09(4) \cdot 10^{-4}$ | 3.8 | 0.086(3) | 2.53(1) | 0.110(4) | 4.12(3). $10^{-2}$ |
| 500 | 0.211(4) | 2.49(1) | 0.120(4) | $2.45(2) \cdot 10^{-2}$ | 4.7 | 0.46(1) | 3.420(5) | 0.069(6) | $2.97(4) \cdot 10^{-4}$ | 4.0 | 0.080(4) | 2.43(1) | 0.116(4) | 2.73 (2). $10^{-2}$ |
| 600 | 0.162(3) | 2.480(7) | 0.113(3) | $2.58(1) \cdot 10^{-2}$ | 5.0 | 0.44(1) | 3.104(5) | 0.058(6) | $2.83(4) \cdot 10^{-4}$ | 4.2 | 0.075(4) | 2.316(8) | 0.118(4) | $1.85(1) \cdot 10^{-2}$ |
| 800 | 0.116(3) | 2.459 (7) | 0.109(3) | 2.70 (1) $\cdot 10^{-2}$ | 5.3 | 0.43(1) | 2.842(4) | 0.044(6) | $2.66(3) \cdot 10^{-4}$ | 4.4 | 0.071(4) | 2.216(6) | 0.118(4) | 1.28(1). $10^{-2}$ |
| 1000 | 0.098(3) | 2.440(8) | 0.109(3) | $2.74(2) \cdot 10^{-2}$ | 5.6 | 0.42(1) | 2.620(4) | 0.029(6) | $2.45(3) \cdot 10^{-4}$ | 4.6 | 0.066(4) | 2.124(6) | 0.121(4) | 9.00(6) $10^{-3}$ |
| 1500 | 0.080(4) | 2.38(1) | 0.111(4) | 2.70 (2) $\cdot 10^{-2}$ | 5.9 | 0.41(1) | 2.431(4) | 0.014(6) | $2.24(3) \cdot 10^{-4}$ | 4.8 | 0.062(5) | 2.042(6) | 0.126(4) | 6.47(5) $10^{-3}$ |
| 2000 | 0.084(7) | 2.30(2) | 0.110(8) | 2.60 (4) $\cdot 10^{-2}$ | 6.2 | 0.41(1) | 2.267(4) | 0.000(6) | 2.01(3) $10^{-4}$ | 5.0 | 0.058(5) | 1.964(5) | 0.127(4) | 4.73(4). $10^{-3}$ |
| 2500 | 0.069(4) | 2.24(1) | 0.130(5) | 2.43(2). $10^{-2}$ | 6.5 | 0.41(2) | 2.124(3) | 0.000(8) | $1.79(4) \cdot 10^{-4}$ | 5.3 | 0.050(7) | 1.860(6) | 0.134(5) | $3.07(3) \cdot 10^{-3}$ |
| 3000 | 0.064(4) | 2.14(1) | 0.137(5) | $2.19(2) \cdot 10^{-2}$ | 6.8 | 0.40(2) | 1.997(3) | 0.00(1) | $1.55(4) \cdot 10^{-4}$ | 5.6 | 0.043(8) | 1.766(6) | 0.139(6) | $2.05(2) \cdot 10^{-3}$ |
| 3500 | 0.049(6) | 2.04(1) | 0.157(7) | $1.89(3) \cdot 10^{-2}$ | 7.1 | 0.40(3) | 1.885(3) | 0.00(1) | $1.36(5) \cdot 10^{-4}$ | 5.9 | 0.04(1) | 1.681(5) | 0.146(7) | $1.41(2) \cdot 10^{-3}$ |
| 4000 | 0.053(6) | 1.91(1) | 0.020(33) | $8.57(67) \cdot 10^{-3}$ | 7.4 | 0.44(4) | 1.784(2) | 0.00(2) | $1.23(7) \cdot 10^{-4}$ | 6.2 | 0.03(1) | 1.603(5) | 0.150(7) | 9.92(14) $10^{-4}$ |
| 4500 | 0.044(7) | 1.80(1) | 0.020(36) | $7.05(43) \cdot 10^{-3}$ | 7.7 | 0.47(5) | 1.693(2) | 0.00(2) | $1.11(9) \cdot 10^{-4}$ | 6.5 | 0.04(1) | $1.533(5)$ | 0.153(8) | $7.21(11) \cdot 10^{-4}$ |
| 5000 | 0.03(1) | 1.67(1) | 0.049(36) | $5.86(35) \cdot 10^{-3}$ |  |  |  |  |  | 6.8 | 0.05(1) | 1.467(5) | 0.149(9) | $5.43(10) \cdot 10^{-4}$ |
| 5500 | 0.02(1) | 1.56(2) | 0.061(24) | $4.47(18) \cdot 10^{-3}$ |  |  |  |  |  | 7.1 | 0.08(2) | 1.406(4) | 0.143(9) | 4.23(9). $10^{-4}$ |
| 6000 | 0.00(1) | 1.45(2) | 0.074(24) | $3.41(13) \cdot 10^{-3}$ |  |  |  |  |  | 7.4 | 0.13(2) | 1.351(4) | 0.127(10) | 3.47 (8). $10^{-4}$ |
| 7000 | 0.00(2) | 1.23(2) | 0.096(15) | 1.94(5) $10^{-3}$ |  |  |  |  |  | 7.7 | 0.19(2) | 1.299(3) | 0.106(9) | $2.96(6) \cdot 10^{-4}$ |
| 8000 | 0.00(3) | 1.03(2) | 0.104(15) | $1.10(3) \cdot 10^{-3}$ |  |  |  |  |  | 8.0 | 0.25(1) | 1.250(2) | 0.083(8) | $2.56(5) \cdot 10^{-4}$ |
| 9000 | 0.00(4) | 0.88(2) | 0.107(38) | 6.51(33) $10^{-4}$ |  |  |  |  |  |  |  |  |  |  |
| 10000 | 0.00(6) | 0.75(2) | 0.133(30) | 4.19(16) $10^{-4}$ |  |  |  |  |  |  |  |  |  |  |
| 11000 | 0.00(8) | 0.65(2) | 0.149(54) | $2.84(22) \cdot 10^{-4}$ |  |  |  |  |  |  |  |  |  |  |
| 12000 | 0.00(11) | 0.58(3) | 0.221(31) | 1.91(12) $10^{-4}$ |  |  |  |  |  |  |  |  |  |  |



Figure S10. In-phase ( $\chi^{\prime}$ ) and out-of-phase ( $\chi^{\prime \prime}$ ) AC susceptibilities for 3 at 4.0 K measured in the 0-12000 Oe magnetic field range (a), at 0 Oe in the 3.2-7.7 K temperature range (b) and at 1500 Oe in the 3.0-8.0 K temperature range (c). Values of $\alpha$ and $\tau$ parameters are gathered in Table S6. The solid lines are the best fits to generalized Debye model.

Table S7. Values of the best-fit parameters obtained by fitting the $\tau^{1}(H)$ using Eq. 2 (top part of the table) and the $\tau(T)$ using Eq. 3 (bottom part of the table) for compound 3. The relevant fits are shown in Figure 3c and 3d in the main text, respectively. This table is also included in Table 3 in the main text.
$\tau^{-1}(H)=A_{1} /\left(1+A_{2} H^{2}\right)+B_{1} H^{4}+D$
(Eq. 2)
$\ln \tau\left(T^{-1}\right)=\ln \left[\left(A+B_{2} T+C T^{n}+\tau_{0}^{-1} \exp \left(-U_{\text {eff }} / k_{B} T\right)\right)^{-1}\right]$

| $T(\mathrm{~K})$ | 4.0 |  |
| :--- | :---: | :---: |
| range (Hz) | $1-1000$ |  |
| field range (Oe) | $0-12000$ |  |
| $A_{1}\left(\mathrm{~s}^{-1}\right)$ | $3096(104)$ |  |
| $A_{2}\left(\mathrm{Oe}^{-2}\right)$ | $7(3) \cdot 10^{-3}$ |  |
| $B_{1}\left(\mathrm{Oe}^{-4}\right)$ | $2.36(4) \cdot 10^{-13}$ |  |
| $D\left(\mathrm{~s}^{-1}\right)$ | $34(4)$ |  |
| $R^{2}$ | 0.99569 | 1500 |
| $H_{\text {DC }}(\mathrm{Oe})$ | 0 | $1-1000$ |
| range $(\mathrm{Hz})$ | $1-1000$ | $3.0-8.0$ |
| temp. range (K) | $3.2-7.7$ | 0.20 (fixed) |
| $A\left(\mathrm{~s}^{-1}\right)$ | 3096 (fixed) |  |
| $B_{2}\left(\mathrm{~s}^{-1} \mathrm{~K}^{-1}\right)$ | 0 (fixed) |  |
| $C\left(\mathrm{~s}^{-1} \mathrm{~K}^{-n}\right)$ |  |  |
| $n$ | 7 (fixed) | $1.5(1) \cdot 10^{-3}$ |
| $\tau_{0}(\mathrm{~s})$ |  | $7(3) \cdot 10^{-1}$ (fixed) |
| $\left(U_{\text {eff }} / \mathrm{k}_{\mathrm{B}}\right)(\mathrm{K})$ |  | 7 (fixed) |
| $R^{2}$ |  | 0.99818 |
|  |  |  |

## AC magnetic properties of compound 4 (Section AC4)

Frequency dependence of the out-of-phase magnetic susceptibility ( $\chi^{\prime \prime}$ ) of 4 shows a single maximum in the $18.6-1000 \mathrm{~Hz}$ range under small applied DC field ( $T=1.8 \mathrm{~K}$; Table S 8 and Figure S 11 a ). This $\chi^{\prime \prime}$ maximum shifts with the magnetic field from around 850 Hz at 200 Oe to around 100 Hz at 2000 Oe (optimum field) and then again to higher frequencies in the 2000-4500 Oe range. Contributions from the QTM and the direct process have been identified in the low- and high-field ranges of the $\tau(H)$ dependence (Figure 3e), respectively. Fitting of the $\tau(H)$ dependence using Eq. 2 resulted in best-fit parameters $A_{1}, A_{2}, A_{3}$ and $A_{4}$ which were used to calculate parameters: $A=157.6 \mathrm{~s}^{-1}$ and $B_{2}=63.1 \mathrm{~s}^{-1} \cdot \mathrm{~K}^{-1}$ in Eq. 3 . The temperature dependence of the AC magnetic susceptibility was measured at the optimum DC field of 2000 Oe in order to analyze the temperature dependence of the relaxation process in a possibly wide temperature range. The $\chi$ " maximum is visible up to 3.3 K in the $18.6-1000 \mathrm{~Hz}$ window and shows larger thermal shift above 2.5 K (Figure S11b and Table S8). Fitting of the temperature dependence of $\ln (\tau)$ using Eq. 3 (see main text) with $A$ and $B_{2}$ fixed based on the results of $\tau(H)$ dependence fitting, shows a significant contribution of the phonon bottle-neck process ( $C=93(33) \mathrm{s}^{-1} \cdot \mathrm{~K}^{-2}, n=2.0(5)$ ) and Orbach relaxation process ( $\left.\tau_{0}=6.0(15.2) \cdot 10^{-10} \mathrm{~s} ; U_{\text {eff }} / \mathrm{k}_{\mathrm{B}}=41(8) \mathrm{K}\right)$ (Figure 3 f ). All parameters obtained by fitting the respective relaxation rates vs magnetic field and temperature using Eqs 2 and 3 are collected in Table S9 and in Table 3 in the main text.

Table S8. Values of $\alpha$ and $\tau$ parameters with uncertainties provided in parentheses obtained from generalized Debye model fitting of the $\tau(v)$ dependences for 4 at 1.8 K in the $0.020-0.450 \mathrm{~T}$ magnetic field range (columns 1-5) and at 0.200 T measured at various temperatures in the 1.8-3.5 K range (columns 6-10).

| Figure S11a |  |  |  |  | Figure S11b |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $T=1.8 \mathrm{~K}$ |  |  |  |  | $H=2000$ Oe |  |  |  |  |
| $\begin{gathered} H \\ (\mathrm{Oe}) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \\ \hline \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \\ \hline \end{gathered}$ | $\alpha$ | $\begin{gathered} \tau \\ (\mathrm{s}) \\ \hline \end{gathered}$ | $\begin{gathered} \hline \boldsymbol{T} \\ (\mathrm{K}) \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{s}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \\ \hline \end{gathered}$ | $\begin{gathered} \chi_{\mathrm{t}} \\ \left(\mathrm{~cm}^{3} \cdot \mathrm{~mol}^{-1}\right) \\ \hline \end{gathered}$ | $\alpha$ | $\begin{gathered} \tau \\ (\mathrm{s}) \\ \hline \end{gathered}$ |
| 200 | 3.26(9) | 4.93(2) | 0.012(30) | 1.97(14) $10^{-4}$ | 1.8 | 0.21(1) | 2.80(2) | 0.306(7) | 16.9(3) $\cdot 10^{-4}$ |
| 300 | 2.19(6) | 4.78 (2) | 0.069(16) | 2.52(9) $10^{-4}$ | 1.9 | 0.21(2) | 2.73(2) | 0.319(9) | $15.8(3) \cdot 10^{-4}$ |
| 400 | 1.48(4) | 4.65(1) | 0.132(9) | 3.40(7) $10^{-4}$ | 2.0 | 0.23(2) | 2.67(2) | 0.310(10) | $15.0(3) \cdot 10^{-4}$ |
| 500 | 1.05(3) | 4.54(1) | 0.172(7) | 4.46(7) $10^{-4}$ | 2.1 | 0.23(1) | 2.70(1) | 0.318(6) | $15.0(2) \cdot 10^{-4}$ |
| 600 | 0.78(3) | 4.44(1) | 0.200(6) | 5.73(7) $10^{-4}$ | 2.3 | 0.24(2) | 2.60(2) | 0.296(8) | $12.7(2) \cdot 10^{-4}$ |
| 800 | 0.47(2) | 4.22(2) | 0.241(6) | $8.27(9) \cdot 10^{-4}$ | 2.5 | 0.26(2) | 2.51(1) | 0.261(7) | 10.0(1) $\cdot 10^{-4}$ |
| 1000 | 0.30(2) | 4.03(2) | 0.283(5) | 10.9(1) $\cdot 10^{-4}$ | 2.7 | 0.28(2) | 2.36(1) | 0.211(9) | $6.71(11) \cdot 10^{-4}$ |
| 1500 | 0.21(1) | 3.37(2) | 0.310(5) | 15.6(2) $10^{-4}$ | 2.9 | 0.25(4) | 2.28(1) | 0.194(15) | 4.12(14) $10^{-4}$ |
| 2000 | 0.19(2) | 2.77(3) | 0.319(10) | 16.9(4) $10^{-4}$ | 3.1 | 0.22(7) | 2.20(2) | 0.191(21) | $2.36(15) \cdot 10^{-4}$ |
| 2500 | 0.18(1) | 2.25(2) | 0.308(8) | 14.5(2) $10^{-4}$ | 3.3 | 0.19(17) | 2.13 (2) | 0.206(36) | 1.28(21) $10^{-4}$ |
| 3000 | 0.19(2) | 1.74(2) | 0.260(14) | 10.0(3) $10^{-4}$ | 3.5 | -0.05(50) | 2.06(2) | 0.263(54) | $0.54(25) \cdot 10^{-4}$ |
| 3500 | 0.15(2) | 1.43(1) | 0.277(17) | 6.94(25) $10^{-4}$ |  |  |  |  |  |
| 4000 | 0.022(82) | 1.22(3) | 0.371(46) | 4.08(63) $10^{-4}$ |  |  |  |  |  |
| 4500 | -0.011(116) | 1.00(3) | 0.385(60) | $2.41(68) \cdot 10^{-4}$ |  |  |  |  |  |



Figure S11. In-phase ( $\chi^{\prime}$ ) and out-of-phase ( $\chi^{\prime \prime}$ ) AC susceptibilities for 4 at 1.8 K measured in the 0.020-0.450 T magnetic field range (a) and at 0.200 T measured at various temperatures in the 1.8-3.5 K range (b). Values of $\alpha$ and $\tau$ parameters with uncertainties provided in parentheses are gathered in Table S8. The solid lines are the best fits to generalized Debye model.

Table S9. Values of the best-fit parameters obtained by fitting the $\tau(H)$ using Eq. 2 (top part) and the $\tau(T)$ using Eq. 3 (bottom part) for compound 4. The relevant fits are shown in Figure $3 e$ and $3 f$ in the main text, respectively. This table is also included in Table 3 in the main text.
$\tau^{-1}(H)=A_{1} /\left(1+A_{2} H^{2}\right)+B_{1} H^{4}+D$
$\ln \tau\left(T^{-1}\right)=\ln \left[\left(A+B_{2} T+C T^{n}+\tau_{0}^{-1} \exp \left(-U_{\text {eff }} / \mathrm{k}_{\mathrm{B}} T\right)\right)^{-1}\right]$

| $T(\mathrm{~K})$ | 1.8 |
| :--- | :--- |
| range $(\mathrm{Hz})$ | $18.6-1000$ |
| field range (Oe) | $200-4500$ |
| $A_{1}\left(\mathrm{~s}^{-1}\right)$ | $7281(348)$ |
| $A_{2}\left(\mathrm{Oe}^{-2}\right)$ | $1.13(7) \cdot 10^{-5}$ |
| $B_{1}\left(\mathrm{Oe}^{-4}\right)$ | $7.1(2) \cdot 10^{-12}$ |
| $D\left(\mathrm{~s}^{-1}\right)$ | $322(9)$ |
| $\mathrm{R}^{2}$ | 0.99893 |
| $H_{\mathrm{DC}}(\mathrm{T})$ | 2000 |
| range $(\mathrm{Hz})$ | $18.6-1000$ |
| temp. range $(\mathrm{K})$ | $1.8-3.5$ |
| $A\left(\mathrm{~s}^{-1}\right)$ | $157.6($ fixed $)$ |
| $B_{2}\left(\mathrm{~s}^{-1} \mathrm{~K}^{-1}\right)$ | $63.1($ fixed $)$ |
| $C\left(\mathrm{~s}^{-1} \mathrm{~K}^{-n}\right)$ | $93(33)$ |
| $n$ | $2.0(5)$ |
| $\tau_{0}(\mathrm{~s})$ | $6.0(15.2) \cdot 10^{-10}$ |
| $\left(U_{\text {eff }} / \mathrm{k}_{\mathrm{B}}\right)(\mathrm{K})$ | $41(8)$ |
| $\mathrm{R}^{2}$ | 0.98322 |

## Ab initio calculations

Ab initio calculations were carried out using the OpenMolcas quantum chemistry software package, ${ }^{8}$ and were performed on experimental geometries taken from the single crystal X-ray diffraction analysis. For 3 two non-equivalent molecules denoted as 3A and 3B were considered, while for $\mathbf{1}$ and $\mathbf{4}$ only one molecule was taken for analysis. The fragments of crystal structure employed for analysis together with calculated main quantization axes (z) are presented in Figure 4. Two models with different basis sets were used: $\mathbf{S}$ - small with VDZP basis function quality for Er"'I central ion and VDZ for other atoms and $\mathbf{L}$ - large with VTZP basis for Er"II and VDZP for atoms in the first coordination sphere together with Li ions. Table S10 contains contractions and labels of basis sets $\mathbf{S}$ and $\mathbf{L}$ for all atoms. The performed calculations were of CASSCF/RASSI/SINGLE_ANISO type. ${ }^{9}$ Scalar relativistic effects were taken into account by employing second order DKH (Douglas-Kroll-Hess) Hamiltonian and relativistic basis sets of an ANO-RCC type. In order to save disk space, Cholesky decomposition of ERI-s (electron repulsion integrals) was used with the $1.0 \cdot 10^{-8}$ threshold. In the first step, a state average multi-configurational self-consistent field (SA-CASSCF) calculation for 35 quartets and 112 doublets rising from different Erl'I configurations was carried out. The active space was composed out of seven f-orbitals with eleven active valence electrons - CAS(11in7). In the next step, all quartets and doublets optimized as spin-free states in CASSCF step were mixed by the spin-orbit coupling within RASSI (Restricted Active Space State Interaction Program) using mean-field spin-orbit (SO) integrals (AMFI) resulting in 364 spin-orbit states. In the final step, a SINGLE_ANISO module was used to decompose the spin-orbit states into states with a definite projection of the total momentum on the located quantization axis and to calculate crystal-field parameters of the Zero Field Splitting (ZFS) Hamiltonian within the manifold of $J=15 / 2$ together with three components of the pseudo-g-tensor for eight ground Kramers doublets. ${ }^{10}$ The obtained energy splitting of the $J=15 / 2$ manifold, together with the $g_{x}, g_{y}, g_{z}$ components of the pseudo-g-tensors within the basis of each doublet ( $\tilde{S}=1 / 2$ ) and decomposition of the ground state into states with a definite angular momentum on quantization axis are presented in Tables S11-S14. All 27 crystal-field parameters of rank $k=2,4$ and 6 were gathered in Tables S15-S17 together with their weights on the crystal field splitting. Magnetic transition moments between states for $\mathbf{1 , 3}$ and $\mathbf{4}$ are visualized in Figure 4 in the article. Comparison of the calculated $M(H)$ and $\chi T(T)$ dependencies (solid lines) with experimental ones (points) can be found in Figures S3, S5 and S6. The agreement between the calculated and experimental data is satisfactory.

Attempts to calculate the energy diagram and magnetic properties of $\mathbf{2}$ were unsuccessful due to the additional spin on the TEMPO ligand. However, in order to better understand the crystal field in $\mathbf{2}$, we have replaced the $N$ atom in TEMPO with the $C$ atom so that the TEMPO has lost its radical character and the calculations were possible for such a modified model of 2 (Table S18 and S19).

Table S10. Description and contractions of the basis sets (two models: $\mathbf{S}$ - small, $\mathbf{L}$ - large) employed in ab initio calculations of the ErlII crystal field in Er_ttbp (compound 1), Er_bht_1 (compound 3), Er_bht_2 (compound 3) and Er_n3ncl2 Er_ttbp (compound 4).

| Basis set "S" | Basis set "L" |
| :--- | :--- |
| Er.ANO-RCC-VDZP 7S6P4D2F1G | Er.ANO-RCC-VTZP 8S7P5D3F2G1H |
| Li.ANO-RCC-VDZ 3S2P | Li.ANO-RCC-VDZP 3S2P1D |
| CI.ANO-RCC-VDZ 4S3P | Cl.ANO-RCC-VDZP 4S3P1D |
| Si.ANO-RCC_VDZ 4S3P | Si.ANO-RCC_VDZ 4S3P |
| N.ANO-RCC-VDZ 3S2P | N.ANO-RCC-VDZP 3S2P1D |
| O.ANO-RCC-VDZ 3S2P | O.ANO-RCC-VDZ 3S2P (distant atoms in Er_n3ncl2) <br> O.ANO-RCC-VDZP 3S2P1D (first coordination sphere in Er_bht) |
| C.ANO-RCC-VDZ 3S2P | C.ANO-RCC-VDZ 3S2P |
| H.ANO-RCC-VDZ 2S | H.ANO-RCC-VDZ 2S |

Table S11. Summary of the energy splitting of the ${ }^{4} \mathrm{I}_{15 / 2}$ multiplet of $\mathrm{Er}{ }^{\prime \prime \prime \prime}$ in compound 1 with pseudo- $g$-tensors of each Kramers doublet and the composition of the ground state in $\left|m_{j}\right\rangle$ basis.

| Compound 1 |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{S}$ |  |  |  | L |  |  |  |
| Energy and Pseudo- $g$-tensor components ( $\mathrm{g}_{\mathrm{x}}, \mathrm{g}_{\mathrm{y}}, \mathrm{g}_{2}$ ) of 8 ground Kramers doublets |  |  |  |  |  |  |  |
| Energy / cm ${ }^{-1}$ | Pseudo-g-tensor components |  |  | Energy / cm ${ }^{-1}$ | Pseudo-g-tensor components |  |  |
|  | $\mathrm{g}_{\mathrm{x}}$ | $\mathrm{g}_{\mathrm{y}}$ | $\mathrm{g}_{2}$ |  | $\mathrm{g}_{\mathrm{x}}$ | $\mathrm{g}_{\mathrm{y}}$ | $\mathrm{g}_{2}$ |
| 0.000 | 0.0166 | 0.0261 | 17.6511 | 0.000 | 0.0070 | 0.0129 | 17.6548 |
| 99.740 | 0.2645 | 0.6156 | 14.2647 | 104.431 | 0.0670 | 0.1459 | 14.8648 |
| 140.763 | 0.5799 | 1.1535 | 15.9697 | 156.372 | 0.5801 | 1.3121 | 15.2402 |
| 162.438 | 1.2964 | 3.3779 | 10.0513 | 176.619 | 0.6328 | 2.5562 | 10.1457 |
| 226.866 | 7.7119 | 6.3032 | 2.6093 | 241.634 | 7.0440 | 6.5546 | 2.9839 |
| 317.568 | 2.1591 | 4.0788 | 8.2183 | 330.470 | 0.9440 | 4.0624 | 8.7756 |
| 371.244 | 2.3020 | 3.6540 | 11.2772 | 388.376 | 2.1676 | 2.7734 | 11.3834 |
| 423.222 | 0.4566 | 1.6293 | 16.2343 | 439.286 | 0.4359 | 1.5891 | 15.8730 |
| Composition of ground Kramers doublet in the ${ }^{\left\|m_{j}\right\rangle}$ basis on the quantization axis of the ground doublet within $\mathrm{J}=$ 15/2 manifold |  |  |  |  |  |  |  |
| $\begin{gathered} 90.0 \%\|+15 / 2\rangle \\ 5.9 \%\|-15 / 2\rangle \\ 2.0 \%\|+11 / 2\rangle \\ 1.3 \%\|+9 / 2\rangle \\ 0.3 \%\|+7 / 2\rangle \\ 0.2 \%\|+13 / 2\rangle \end{gathered}$ |  |  |  | $\begin{gathered} 65.4 \%\|+15 / 2\rangle \\ 30.3 \%\|-15 / 2\rangle \\ 1.8 \%\|+11 / 2\rangle \\ 0.8 \%\|-11 / 2\rangle \\ 0.8 \%\|+9 / 2\rangle \\ 0.4 \%\|-9 / 2\rangle \end{gathered}$ |  |  |  |

Table S12. Summary of the energy splitting of the ${ }^{4}{ }_{15 / 2}$ multiplet of Er "II in compound $\mathbf{3}$ molecule $\mathbf{A}$ model with pseudo- $g$ tensors of each Kramers doublet and the composition of the ground state in $\left|m_{j}\right\rangle$ basis.

| 3A |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| S |  |  |  | L |  |  |  |
| Energy and Pseudo- $g$-tensor components ( $\mathrm{g}_{\mathrm{x}}, \mathrm{g}_{\mathrm{y}}, \mathrm{g}_{2}$ ) of 8 ground Kramers doublets |  |  |  |  |  |  |  |
| Energy / $\mathrm{cm}^{-1}$ | Pseudo- $g$-tensor components |  |  | Energy / $\mathrm{cm}^{-1}$ | Pseudo-g-tensor components |  |  |
|  | $\mathrm{g}_{\mathrm{x}}$ | gy | $\mathrm{g}_{2}$ |  | $\mathrm{g}_{\mathrm{x}}$ | gy | $\mathrm{g}_{2}$ |
| 0.000 | 0.0063 | 0.0101 | 17.4817 | 0.000 | 0.0047 | 0.0082 | 17.4996 |
| 89.941 | 0.1155 | 0.1614 | 14.8081 | 95.310 | 0.1209 | 0.1622 | 14.9565 |
| 143.673 | 0.2858 | 0.4257 | 15.4163 | 155.664 | 0.4379 | 0.6353 | 14.5597 |
| 166.980 | 0.5369 | 1.1325 | 10.6607 | 180.602 | 0.5652 | 1.2815 | 11.0978 |
| 226.748 | 3.8524 | 4.7912 | 7.3881 | 239.456 | 3.4272 | 4.6192 | 7.4037 |
| 301.051 | 0.7321 | 4.0036 | 9.1724 | 309.892 | 1.7007 | 3.8112 | 9.7762 |
| 361.458 | 1.5917 | 2.4556 | 11.9526 | 375.625 | 1.2780 | 2.0367 | 12.6439 |
| 419.942 | 0.3034 | 1.0719 | 16.0860 | 444.682 | 0.1960 | 0.6264 | 16.3898 |

Composition of ground Kramers doublet in ${ }^{\left|m_{j}\right\rangle}$ basis on the quantization axis within J=15/2 manifold
$85.9 \%|+15 / 2\rangle$
$5.7 \%|-15 / 2\rangle$
$6.2 \%|+11 / 2\rangle$
$1.0 \%|+9 / 2\rangle$
$0.5 \%|+13 / 2\rangle$
$92.0 \%|+15 / 2\rangle$
$6.5 \%|+11 / 2\rangle$
$1.0 \%|+9 / 2\rangle$
$0.4 \%|+13 / 2\rangle$

Table S13. Summary of the energy splitting of the ${ }^{4} 1_{15 / 2}$ multiplet of Er"II in compound $\mathbf{3}$ molecule $\mathbf{B}$ with pseudo- $g$ tensors of each Kramers doublet and the composition of the ground state in $\left|m_{j}\right\rangle$ basis.

| 3B |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| S |  |  |  | L |  |  |  |
| Energy and Pseudo- $g$-tensor components ( $\mathrm{g}_{\mathrm{x}}, \mathrm{g}_{\mathrm{y}}, \mathrm{g}_{\mathrm{z}}$ ) of 8 ground Kramers doublets |  |  |  |  |  |  |  |
| Energy / cm ${ }^{-1}$ | Pseudo-g-tensor components |  |  | Energy / $\mathrm{cm}^{-1}$ | Pseudo-g-tensor components |  |  |
|  | $\mathrm{g}_{\mathrm{x}}$ | $\mathrm{g}_{\mathrm{y}}$ | $\mathrm{g}_{2}$ |  | $\mathrm{g}_{\mathrm{x}}$ | $\mathrm{g}_{\mathrm{y}}$ | $\mathrm{g}_{2}$ |
| 0.000 | 0.0013 | 0.0049 | 17.2398 | 0.000 | 0.0008 | 0.0016 | 17.2563 |
| 94.418 | 0.0628 | 0.0913 | 14.2288 | 97.853 | 0.0806 | 0.1074 | 14.4333 |
| 158.601 | 0.1230 | 0.1888 | 17.4167 | 170.528 | 0.4235 | 0.6012 | 15.2183 |
| 172.526 | 0.5288 | 0.8592 | 10.4274 | 183.870 | 0.6178 | 1.2536 | 10.7061 |
| 236.249 | 3.9879 | 4.4280 | 7.7190 | 246.739 | 3.3865 | 4.4526 | 7.5997 |
| 312.298 | 0.3023 | 3.7431 | 9.2701 | 320.790 | 1.0493 | 3.7491 | 9.2548 |
| 373.865 | 1.6225 | 2.3838 | 12.5233 | 385.778 | 1.5180 | 2.2695 | 12.9284 |
| 430.487 | 0.3901 | 1.4739 | 16.0180 | 455.491 | 0.2613 | 0.8915 | 16.3942 |
| Composition of ground Kramers doublet in ${ }^{\left\|m_{J}\right\rangle}$ basis on the quantization axis within $\mathrm{J}=15 / 2$ manifold |  |  |  |  |  |  |  |
| $\begin{gathered} 69.6 \%\|+15 / 2\rangle \\ 17.0 \%\|-15 / 2\rangle \\ 6.9 \%\|+11 / 2\rangle \\ 1.7 \%\|+9 / 2\rangle \\ 1.8 \%\|+13 / 2\rangle \end{gathered}$ |  |  |  | $\begin{gathered} 85.3 \%\|+15 / 2\rangle \\ 1.6 \%\|-15 / 2\rangle \\ 9.0 \%\|+11 / 2\rangle \\ 1.6 \%\|+9 / 2\rangle \\ 1.8 \%\|+13 / 2\rangle \end{gathered}$ |  |  |  |

Table S14. Summary of the energy splitting of the ${ }^{4} I_{15 / 2}$ multiplet of Er ${ }^{\prime \prime \prime}$ in compound $\mathbf{4}$ with pseudo- $g$-tensors of each Kramers doublet and the composition of the ground state in ${ }^{\left|m_{j}\right\rangle}$ basis.

| Compound 4 |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| S |  |  |  | L |  |  |  |
| Energy and Pseudo- $g$-tensor components ( $\mathrm{g}_{\mathrm{x}}, \mathrm{g}_{\mathrm{y}}, \mathrm{g}_{\mathrm{z}}$ ) of 8 ground Kramers doublets |  |  |  |  |  |  |  |
| Energy / cm ${ }^{-1}$ | Pseudo-g-tensor components |  |  | Energy / cm ${ }^{-1}$ | Pseudo-g-tensor components |  |  |
|  | $\mathrm{g}_{\mathrm{x}}$ | gy | $\mathrm{g}_{2}$ |  | $\mathrm{g}_{\mathrm{x}}$ | $\mathrm{g}_{\mathrm{y}}$ | $\mathrm{g}_{\text {z }}$ |
| 0.000 | 0.0358 | 0.0824 | 16.4810 | 0.000 | 0.0480 | 0.1031 | 16.4440 |
| 39.535 | 0.0067 | 0.1101 | 12.8726 | 30.556 | 0.0296 | 0.1454 | 12.7906 |
| 101.488 | 2.3153 | 2.3826 | 9.0310 | 89.274 | 1.7998 | 1.9254 | 8.9212 |
| 155.009 | 4.0910 | 5.7441 | 8.3308 | 140.862 | 4.7443 | 5.2939 | 8.8112 |
| 217.760 | 0.0350 | 2.3700 | 13.5905 | 223.688 | 2.2751 | 5.3122 | 8.9437 |
| 256.080 | 0.8858 | 1.8978 | 13.8751 | 249.360 | 1.7824 | 5.1874 | 10.5539 |
| 308.850 | 0.0672 | 0.8419 | 16.5299 | 280.533 | 0.0887 | 0.9437 | 16.3007 |
| 338.370 | 0.0858 | 0.8477 | 16.7534 | 331.665 | 0.1312 | 0.4231 | 17.1844 |
| Composition of ground Kramers doublet in ${ }^{\left\|m_{j}\right\rangle}$ basis on the quantization axis within $\mathrm{J}=15 / 2$ manifold |  |  |  |  |  |  |  |
| $\begin{gathered} 40.5 \%\|+15 / 2\rangle \\ 23.6 \%\|-15 / 2\rangle \\ 12.1 \%\|+13 / 2\rangle \\ 7.3 \%\|+11 / 2\rangle \\ 7.0 \%\|-13 / 2\rangle \\ 4.1 \%\|-11 / 2\rangle \\ 2.6 \%\|+9 / 2\rangle \\ 1.8 \%\|-9 / 2\rangle \end{gathered}$ |  |  |  | $\begin{gathered} 31.5 \%\|+15 / 2\rangle \\ 31.5 \%\|-15 / 2\rangle \\ 9.9 \%\|+13 / 2\rangle \\ 9.9 \%\|-13 / 2\rangle \\ 6.1 \%\|-11 / 2\rangle \\ 5.8 \%\|-11 / 2\rangle \\ 2.3 \%\|-9 / 2\rangle \\ 1.8 \%\|+9 / 2\rangle \end{gathered}$ |  |  |  |

Table S15. Crystal-field parameters, $B_{k}^{q}(k=2,4,6)$ of Er'I' complexes defined in ref. 6 for $\mathbf{1}$ in the basis set $\mathbf{L}$ model obtained using ab initio calculations without symmetry constrains and their weights on crystal field splitting in descending order.

|  |  | $\mathbf{1}$ in the basis set $\mathbf{L}$ |  |
| :---: | :---: | :---: | ---: |
| $k$ | $q$ | $B_{k}^{q}$ | Weight $/ \%$ |
| 2 | 0 | $-2.07 \mathrm{E}+00$ | 20.78996 |
| 4 | -1 | $-4.76 \mathrm{E}-03$ | 8.669356 |
| 2 | -1 | $7.66 \mathrm{E}-01$ | 7.688352 |
| 4 | 3 | $3.20 \mathrm{E}-03$ | 5.8338 |
| 4 | -4 | $-3.10 \mathrm{E}-03$ | 5.657771 |
| 2 | 2 | $4.79 \mathrm{E}-01$ | 4.812641 |
| 6 | -4 | $2.48 \mathrm{E}-05$ | 4.230257 |
| 4 | -3 | $2.28 \mathrm{E}-03$ | 4.158291 |
| 6 | 2 | $-2.05 \mathrm{E}-05$ | 3.500833 |
| 4 | -2 | $-1.75 \mathrm{E}-03$ | 3.191822 |
| 6 | -6 | $1.83 \mathrm{E}-05$ | 3.122037 |
| 6 | -1 | $1.73 \mathrm{E}-05$ | 2.959918 |
| 2 | 1 | $2.85 \mathrm{E}-01$ | 2.861501 |
| 6 | 0 | $-1.61 \mathrm{E}-05$ | 2.758342 |
| 4 | 4 | $1.34 \mathrm{E}-03$ | 2.449922 |
| 6 | 5 | $1.38 \mathrm{E}-05$ | 2.360959 |
| 6 | 3 | $-1.35 \mathrm{E}-05$ | 2.311446 |
| 2 | -2 | $-2.23 \mathrm{E}-01$ | 2.240751 |
| 4 | 2 | $-1.11 \mathrm{E}-03$ | 2.017328 |
| 6 | 1 | $-1.14 \mathrm{E}-05$ | 1.942748 |
| 6 | -3 | $8.94 \mathrm{E}-06$ | 1.527337 |
| 4 | 0 | $5.77 \mathrm{E}-04$ | 1.051771 |
| 6 | 6 | $-5.48 \mathrm{E}-06$ | 0.935984 |
| 6 | 4 | $-4.65 \mathrm{E}-06$ | 0.795128 |
| 4 | 1 | $-3.95 \mathrm{E}-04$ | 0.720651 |
| 6 | -5 | $-3.93 \mathrm{E}-06$ | 0.671304 |
| 6 | -2 | $1.28 \mathrm{E}-06$ | 0.218606 |

Table S16. Crystal-field parameters, $B_{k}^{q}(k=2,4,6)$ of Erll' complexes defined in ref. 6 for $\mathbf{3 A}$ and $\mathbf{3 B}$ in the basis set $\mathbf{L}$ models obtained using ab initio calculations without symmetry constrains and their weights on crystal field splitting in descending order.

|  |  | 3A in the basis set $\mathbf{L}$ |  | 3B in the basis set $\mathbf{L}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $k$ | $q$ | $B_{k}^{q}$ | Weight / \% | $k$ | $q$ | $B_{k}^{q}$ | Weight / \% |
| 2 | 0 | $-2.05 \mathrm{E}+00$ | 21.58893 | 2 | 0 | -2.09E+00 | 20.62163 |
| 4 | 1 | 3.93E-03 | 7.506708 | 4 | 1 | $4.68 \mathrm{E}-03$ | 8.401507 |
| 2 | 1 | -6.50E-01 | 6.849065 | 4 | 3 | 4.01E-03 | 7.188348 |
| 4 | 3 | -3.54E-03 | 6.767693 | 6 | -2 | -3.96E-05 | 6.668036 |
| 2 | 2 | 6.17E-01 | 6.503988 | 2 | 1 | -6.73E-01 | 6.651333 |
| 4 | 4 | -3.00E-03 | 5.739744 | 4 | -4 | -2.79E-03 | 5.002434 |
| 2 | -1 | 5.23E-01 | 5.506405 | 6 | 3 | -2.55E-05 | 4.28814 |
| 4 | -1 | -2.63E-03 | 5.023797 | 2 | -2 | -4.14E-01 | 4.091854 |
| 4 | 2 | $2.24 \mathrm{E}-03$ | 4.291609 | 4 | -1 | -2.08E-03 | 3.728592 |
| 6 | 3 | $2.15 \mathrm{E}-05$ | 3.863424 | 6 | 4 | -1.98E-05 | 3.328073 |
| 6 | -2 | -2.15E-05 | 3.85813 | 6 | 2 | 1.75E-05 | 2.950623 |
| 6 | 0 | -1.67E-05 | 3.002886 | 4 | -3 | -1.56E-03 | 2.806738 |
| 6 | -6 | 1.52E-05 | 2.723784 | 6 | -3 | -1.55E-05 | 2.61578 |
| 6 | 1 | -1.31E-05 | 2.344341 | 4 | 0 | $1.27 \mathrm{E}-03$ | 2.285823 |
| 6 | -3 | -1.19E-05 | 2.141282 | 6 | 0 | -1.31E-05 | 2.21176 |
| 4 | 0 | 1.07E-03 | 2.051036 | 4 | -2 | -1.01E-03 | 1.816397 |
| 6 | 4 | 1.09E-05 | 1.954757 | 6 | -6 | -9.79E-06 | 1.648233 |
| 6 | 2 | 1.05E-05 | 1.874179 | 6 | 6 | -9.51E-06 | 1.599897 |
| 6 | 5 | 8.36E-06 | 1.498308 | 4 | 2 | -8.90E-04 | 1.596598 |
| 6 | 6 | -7.91E-06 | 1.417589 | 4 | 4 | $8.24 \mathrm{E}-04$ | 1.479102 |
| 4 | -3 | -5.50E-04 | 1.052509 | 6 | 5 | 8.61E-06 | 1.448428 |
| 6 | -4 | -3.96E-06 | 0.709288 | 2 | -1 | $1.35 \mathrm{E}-01$ | 1.337194 |
| 6 | -1 | -2.36E-06 | 0.422403 | 6 | -5 | 7.91E-06 | 1.330937 |
| 2 | -2 | -3.80E-02 | 0.400401 | 6 | -1 | 7.37E-06 | 1.240784 |
| 4 | -2 | 1.25E-04 | 0.239689 | 2 | 2 | -1.23E-01 | 1.220135 |
| 4 | -4 | -8.61E-05 | 0.164654 | 6 | -4 | 6.57E-06 | 1.105149 |
| 6 | -5 | -3.88E-07 | 0.069491 | 6 | 1 | -4.20E-06 | 0.706008 |

Table S17. Crystal-field parameters, $B_{k}^{q}(k=2,4,6)$ of Er'I' complexes defined in ref. 6 for $\mathbf{4}$ in the basis set $\mathbf{L}$ model obtained using ab initio calculations without symmetry constrains and their weights on crystal field splitting in descending order.

|  |  | 4 in the basis set $\mathbf{L}$ |  |
| :---: | :---: | :---: | :---: |
| $k$ | $q$ | $B_{k}^{q}$ | Weight / \% |
| 4 | -1 | -7.40E-03 | 14.47472 |
| 2 | 0 | $-1.01 \mathrm{E}+00$ | 10.86274 |
| 6 | 2 | -5.41E-05 | 9.924257 |
| 4 | -4 | 4.73E-03 | 9.246317 |
| 6 | 3 | -4.93E-05 | 9.045388 |
| 6 | -2 | -4.66E-05 | 8.550529 |
| 6 | 5 | -3.89E-05 | 7.131003 |
| 6 | -3 | 2.91E-05 | 5.342933 |
| 4 | 1 | $2.62 \mathrm{E}-03$ | 5.129553 |
| 4 | 4 | 1.75E-03 | 3.424274 |
| 4 | 0 | 1.33E-03 | 2.603503 |
| 4 | -2 | -1.09E-03 | 2.130787 |
| 4 | 2 | -7.79E-04 | 1.522168 |
| 2 | 2 | 1.30E-01 | 1.404634 |
| 6 | -4 | 6.89E-06 | 1.263064 |
| 6 | -1 | 6.22E-06 | 1.141012 |
| 4 | -3 | -5.22E-04 | 1.020338 |
| 2 | -1 | -9.38E-02 | 1.010148 |
| 6 | 1 | 4.71E-06 | 0.863375 |
| 6 | -5 | 4.26E-06 | 0.781522 |
| 6 | 0 | 3.71E-06 | 0.680811 |
| 2 | -2 | 4.78E-02 | 0.51449 |
| 4 | 3 | 1.95E-04 | 0.381768 |
| 6 | -6 | -1.83E-06 | 0.335396 |
| 2 | 1 | 9.92E-03 | 0.106921 |
| 6 | 6 | -1.72E-07 | 0.031494 |
| 6 | 4 | 1.71E-07 | 0.03137 |

Table S18. Summary of the energy splitting of the ${ }^{4} I_{15 / 2}$ multiplet of Er ${ }^{111}$ in a modified model of compound $\mathbf{2}$ ( N -atom of the TEMPO replaced by C-atom) with pseudo-g-tensors of each Kramers doublet and the composition of the ground state in $\left|m_{J}\right\rangle$ basis.

| Modified model of compound 2 |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| S |  |  |  | L |  |  |  |
| Energy and Pseudo- $g$-tensor components ( $\mathrm{g}_{\mathrm{x}}, \mathrm{g}_{\mathrm{y}}, \mathrm{g}_{2}$ ) of 8 ground Kramers doublets |  |  |  |  |  |  |  |
| Energy / $\mathrm{cm}^{-1}$ | Pseudo-g-tensor components |  |  | Energy / $\mathrm{cm}^{-1}$ | Pseudo-g-tensor components |  |  |
|  | $\mathrm{g}_{\mathrm{x}}$ | gy | $\mathrm{g}_{2}$ |  | $\mathrm{g}_{\mathrm{x}}$ | gy | $\mathrm{g}_{2}$ |
| 0.000 | 0.0363 | 0.1047 | 16.5251 | 0.000 | 0.0362 | 0.0611 | 17.2833 |
| 28.678 | 0.0286 | 0.1235 | 15.7704 | 43.882 | 0.0359 | 0.0905 | 16.4814 |
| 93.649 | 0.2365 | 0.3003 | 15.4267 | 101.200 | 0.3429 | 0.4377 | 14.8384 |
| 120.374 | 1.1035 | 1.3877 | 10.4748 | 138.069 | 1.1025 | 1.3779 | 11.0316 |
| 170.974 | 3.9208 | 4.8550 | 6.9869 | 185.379 | 4.0245 | 4.9822 | 6.4596 |
| 241.144 | 0.5377 | 3.3337 | 8.1177 | 257.340 | 0.2519 | 3.1021 | 8.4893 |
| 285.345 | 1.2909 | 4.4076 | 11.8418 | 308.092 | 1.1303 | 3.5089 | 12.5626 |
| 426.301 | 0.0306 | 0.0633 | 17.5141 | 459.700 | 0.0281 | 0.0602 | 17.5406 |
| Composition of ground Kramers doublet in the ${ }^{\left\|m_{J}\right\rangle}$ basis on the quantization axis of the ground doublet within J = 15/2 manifold |  |  |  |  |  |  |  |
| $\begin{gathered} 69.4 \%\|+15 / 2\rangle \\ 14.1 \%\|+13 / 2\rangle \\ 7.0 \%\|+11 / 2\rangle \\ 6.6 \%\|+9 / 2\rangle \\ 1.6 \%\|+7 / 2\rangle \\ 0.5 \%\|-15 / 2\rangle \end{gathered}$ |  |  |  | $\begin{gathered} \hline 85.1 \%\|+15 / 2\rangle \\ 4.1 \%\|-15 / 2\rangle \\ 4.1 \%\|+11 / 2\rangle \\ 2.9 \%\|+9 / 2\rangle \\ 1.8 \%\|+13 / 2\rangle \\ 1.1 \%\|+7 / 2\rangle \end{gathered}$ |  |  |  |

Table S19. Crystal-field parameters, $B_{k}^{q}(\mathrm{k}=2,4,6)$ of ErlII complexes defined in ref. 6 for a modified model of 2 ( $N$-atom of the TEMPO replaced by C-atom) in the basis set L obtained using ab initio calculations without symmetry constrains and their weights on crystal field splitting in descending order.

|  |  | modified 2 in the basis set L |  |
| :---: | :---: | :---: | :---: |
| $k$ | $q$ | $B_{k}^{q}$ | Weight / \% |
| 2 | 0 | $-1.63 \mathrm{E}+00$ | 14.10893 |
| 2 | 2 | $1.40 \mathrm{E}+00$ | 12.11048 |
| 4 | -1 | -5.73E-03 | 8.999707 |
| 2 | -1 | $9.82 \mathrm{E}-01$ | 8.494804 |
| 4 | 4 | $4.68 \mathrm{E}-03$ | 7.351983 |
| 2 | -2 | -8.40E-01 | 7.267856 |
| 6 | 2 | -3.48E-05 | 5.125603 |
| 4 | -4 | -2.62E-03 | 4.110489 |
| 6 | -4 | $2.64 \mathrm{E}-05$ | 3.887397 |
| 6 | -6 | -2.07E-05 | 3.04814 |
| 4 | 3 | $1.77 \mathrm{E}-03$ | 2.774343 |
| 4 | -3 | $1.75 \mathrm{E}-03$ | 2.741114 |
| 6 | 1 | -1.71E-05 | 2.513217 |
| 2 | 1 | $2.79 \mathrm{E}-01$ | 2.408999 |
| 6 | 5 | $1.37 \mathrm{E}-05$ | 2.011658 |
| 6 | 0 | -1.34E-05 | 1.979419 |
| 6 | 4 | -1.29E-05 | 1.894171 |
| 6 | -1 | $1.25 \mathrm{E}-05$ | 1.839754 |
| 4 | -2 | $1.01 \mathrm{E}-03$ | 1.578217 |
| 6 | 6 | $9.61 \mathrm{E}-06$ | 1.414427 |
| 4 | 2 | -7.80E-04 | 1.223706 |
| 6 | -3 | $4.85 \mathrm{E}-06$ | 0.71395 |
| 4 | 0 | $3.81 \mathrm{E}-04$ | 0.598206 |
| 6 | -5 | $3.41 \mathrm{E}-06$ | 0.501326 |
| 6 | 3 | $2.59 \mathrm{E}-06$ | 0.380612 |
| 6 | -2 | 1.17E-06 | 0.171642 |
| 4 | 1 | $6.51 \mathrm{E}-05$ | 0.102138 |

## References

1. C. C. Cummins, R. R. Schrock, W. M. Davis, Inorg. Chem. 1994, 33, 1448-1457
2. F. Yuan, J. Yang, L. Xiong, J. Organomet. Chem. 2006, 691, 2534-2539
3. G. M. Sheldrick, Acta Crystallogr. C, 2015, 71, 3-8
4. O. V Dolomanov, L. J. Bourhis, R. J. Gildea, J. A. K. Howard and H. Puschmann, J. Appl. Crystallogr., 2009, 42, 339-341
5. M. Arczyński, J. Stanek, B. Sieklucka, K. R. Dunbar and D. Pinkowicz, J. Am. Chem. Soc., 2019, 141, 19067-19077
6. L. Gu, R. Wu, Phys. Rev. B, 2021, 103, 014401
7. K. N. Shrivastava, Phys. Stat. Sol. B, 1983, 117, 437-458
8. I. F. Galvam, M. Vacher, A. Alavi, C. Angeli, F. Aquilante, J. Autschbach, J. J. Bao, S. I. Bokarev, N. A. Bogdanov, R. K. Carlson, L. F. Chibotaru, J. Creutzberg, N. Dattani, M. G. Delcey, S. S. Dong, A. Dreuw, L. Freitag, L. M. Frutos, L. Gagliardi, F. Gendron, A. Giussani, L. Gonzalez, G. Grell, M. Guo, C. E. Hoyer, M. Johansson, S. Keller, S. Knecht, G. Kovacevic, E. Kallman, G. L. Manni, M. Lundberg, Y. Ma, S. Mai, J. P. Malhado, P. A. Malmqvist, P. Marquetand, S. A. Mewes, J. Norell, M. Olivucci, M. Oppel, Q. M. Phung, K. Perloot, F. Plasser, M. Reiher, A. M. Sand, I. Schapiro, P. Sharma, C. J. Stein, L. K. Sorensen, D. G. Truhlar, M. Ugandi, L. Ungur, A. Valentini, S. Vancoillie, V. Veryazov, O. Weser, T. A. Wesołowski, P.-O. Widmark, S. Wouters, A. Zech, J. P. Zobel, R. Lindh, J. Chem. Theory Comput. 2019, 15, 5925-5964
9. L. F. Chibotaru, L. Ungur, J. Chem. Phys. 2012, 137, 064112
10. L. Ungur, L. F. Chibotaru, Chem. Eur. J. 2017, 23, 3708-3718
