Supporting information

Evaluation of bacterial adhesion strength on phospholipid copolymer films with antibacterial ability by microfluidic shear devices

Yuta Kozuka^{a,†}, Zhou Lu^{a,†}, Tsukuru Masuda^a, Shintaro Hara^a, Toshihiro Kasama^a, Ryo Miyake^a, Norifumi Isu^b, and Madoka Takai^{a*}

^aDepartment of Bioengineering, School of Engineering, The University of Tokyo, 73-1, Hongo, Bunkyo-ku, 113-8656, Tokyo, Japan
**E-mail: takai@bis.t.u-tokyo.ac.jp*^bLIXIL Corporation, 2-1-1, Ojima, Koto-ku, 136-8535, Tokyo, Japan
† These authors contributed equally to this work.

SI. 1 Calculations for shear stress at the bottom of the channel

The movement of fluid in the physical domain is directed by governing equations, among which the Navier–Stokes equations^{S1} are widely applied mathematical models to describe the flow properties during dynamic and/or thermal interactions.

$$\rho \left[\frac{\partial \Delta \vec{V}}{\partial t} + (V \cdot \nabla) \vec{V} \right] = -\nabla p + \mu \nabla^2 \vec{V} + \rho \vec{g}. \qquad Eq. S1$$

Eq. S1 is an expression of the Navier–Stokes equations, where ρ is the density of the fluid, *V* is the velocity of the fluid, *t* is the time, *p* is the pressure of the fluid, and *g* is the gravitational acceleration. In this study, the PBS was moved along a rectangular channel with a width (*w*) of 4.2 mm, a height (*h*) of 0.5 mm, and a length (*L*) of 42 mm. The Navier–Stokes equation can be extended as:

As fully developed fluid moves along the microchannel, $V_y = V_z = 0$ and $\frac{\partial p}{\partial y} = \frac{\partial p}{\partial z} = 0$. In this situation, V_x is the only velocity component. The velocity did not vary with time at a; hence, $\frac{\partial \Delta \vec{V}}{\partial t} = 0$. As $w \gg h$, the effect of the sidewall was negligible; therefore, the microchannel can

be regarded as a parallel plate flow chamber wherein the fluid velocity did not vary along z-

direction, $\frac{\partial V_x}{\partial z} = 0$, and V_x is the function of y only, $\frac{\partial V_x}{\partial x} = 0$. Finally, the channel was set horizontally such that the movement of the fluid is perpendicular to the gravitational force, $\vec{g} = 0$. Therefore, the equation can be simplified as,

$$0 = -\frac{\partial p}{\partial x} + \mu \frac{\partial^2 V_x}{\partial y^2}.$$
 Eq.S3

The medium moves along the x direction, $V_y = 0$; therefore $-\frac{\partial p}{\partial x} = -\frac{dp}{dx}$, and p = p(x).

$$V_x$$
 is a function of y only, implying $\frac{\partial p}{\partial x} = \mu \frac{\partial^2 V_x}{\partial y^2} = p'$ (constant).

After integrating twice, the resulting equation is,

$$V_x = \frac{p'}{2\mu}y^2 + c_1y + c_2.$$
 Eq. S4

Based on the assumption that there is no slip condition at the solid surface, where

$$V_{x} = f(y) = 0 \ (y = \pm h/2),$$
$$V_{x} = \frac{p'h^{2}}{2\mu} \left(\frac{y^{2}}{h^{2}} - \frac{1}{4}\right). \qquad Eq. S5$$

The flow rate (Q) in the microchannel can be obtained as,

$$Q = \int_{-h/2}^{h/2} V_x(y) \cdot w = -\frac{P'h^3}{12\mu} \cdot w.$$

Eq.S6
Therefore, $p' = -\frac{12\mu Q}{h^3 w}$ is obtained.

Next, plot for Poiseuille flow to calculate the shear stress.

$$\tau = \mu \frac{dV_x}{dy} = \mu \frac{p'}{\mu} y = -\frac{12\mu Q}{h^3 w} y \qquad Eq. S7$$

At the bottom of the channel, substituting y = -h/2 into Eq. S7, the shear stress can be expressed as:

$$\tau = \frac{6\mu Q}{h^2 w}.$$
 Eq. S8

SI. 2 Quartz crystal microbalance (QCM) measurement for fibrinogen adsorption



Figure S1. Result of the Fg adsorption on SiO₂, C-20, and C-40 surfaces measured by

QCM. $\Delta f(f \text{ in Hz})$ is the normalized frequency change by 7th overtone.

References

S1 F. White, Viscous Fluid Flow, McGraw-Hill Mechanical Engineering: 1991.