

## Electronic Supplementary Information for “Topological bands in PdSe<sub>2</sub> pentagonal monolayer”

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### A. Topological invariant extraction from the symmetry indicators of space group #14

In this section we follow Ref. [1] applying the expressions for the particular case of 2D space group (SG) #14. The Smith normal form presented in main text can be decomposed as

$$\Delta = L \cdot EBR \cdot R. \quad (\text{S.1})$$

Here  $EBR$  is the matrix of the elementary band representations (EBR).  $L$  and  $R$  are unimodular matrices that can be obtained from the numerical computation of  $\Delta$ . Choosing a particular ordering of the irreducible representations (IR) in the EBR we have the  $EBR'$  matrix

$$EBR' = \begin{bmatrix} 2 & 2 & 0 & 0 \\ 1 & 0 & 1 & 0 \\ 0 & 0 & 2 & 2 \\ 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{bmatrix}. \quad (\text{S.2})$$

This particular choice renders all the elements in the  $\Delta$  matrix positive. Next a general band representation is defined by a vector  $B$  that comprises the multiplicities of the IR of the set of bands of interest, such that

$$B = \begin{bmatrix} n_{\Gamma_3\Gamma_4} \\ n_{D_3} \\ n_{\Gamma_5\Gamma_6} \\ n_{D_4} \\ n_{D_5} \\ n_{D_6} \\ n_{Z_2} \\ n_{B_2} \end{bmatrix}. \quad (\text{S.3})$$

The topological information of the bands is encoded in the vector (eq.(6) in [1])

$$C = L \cdot B. \quad (\text{S.4})$$

In the case of SG #14, for the particular  $EBR'$  matrix presented above we obtain for the L matrix

$$L = \begin{bmatrix} 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & -1 & 0 & 0 & 0 & 0 \\ 1 & -2 & 1 & 0 & -2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 & -1 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 & -1 & 0 & 1 & 0 \end{bmatrix}. \quad (\text{S.5})$$

Thereby, the C vector is given by

$$C = \begin{bmatrix} n_{D_5} \\ n_{D_3} \\ n_{\Gamma_5\Gamma_6} \\ -n_{D_3} + n_{D_4} \\ n_{\Gamma_3\Gamma_4} - 2n_{D_3} + n_{\Gamma_5\Gamma_6} - 2n_{D_5} \\ -n_{D_5} + n_{D_6} \\ -n_{D_3} - n_{D_5} \\ -n_{D_3} - n_{D_5} \end{bmatrix}. \quad (\text{S.6})$$

Additionally we set  $r$  as the range of the  $EBR'$  matrix, that is  $r = 3$ . The component  $C_r$  of vector  $C$  yields the definition of the topological index as stated by eq. (8) in [1]. For our case this results in the relation

$$\mathbb{Z}_2 = n_{\Gamma_5\Gamma_6} \pmod{2}, \quad (\text{S.7})$$

which defines the topological invariant only in terms of the multiplicities of the  $\Gamma$  point IR (An equivalent result holds for the other IR,  $\Gamma_3\Gamma_4$ , by choosing

a different ordering in the EBR matrix). A complementary result: using the components  $C_i$  and imposing the condition  $C_i = 0$  for  $i > r$  allows to construct the compatibility relations for the particular band representation.

## B. Spin Hall conductivity formula

The optical spin Hall conductivity, is calculated with the Kubo-Greenwood formula [2]

$$\sigma_{\alpha\beta}^{\gamma}(\omega) = \frac{\hbar}{V_c N_{\mathbf{k}}} \sum_{\mathbf{k}} \sum_n f_{n\mathbf{k}} \sum_{m \neq n} \frac{\text{Im} \left[ \langle n\mathbf{k} | \hat{j}_{\alpha}^{\gamma} | m\mathbf{k} \rangle \langle m\mathbf{k} | -e\hat{v}_{\beta} | n\mathbf{k} \rangle \right]}{(\epsilon_{n\mathbf{k}} - \epsilon_{m\mathbf{k}})^2 - (\hbar\omega - i\eta)^2}. \quad (\text{S.8})$$

Here  $\hat{j}_{\alpha}^{\gamma} = \frac{1}{2} \{ \hat{s}_{\gamma}, \hat{v}_{\alpha} \}$  is the spin current operator with the spin operator  $\hat{s}_{\gamma}$  defined as  $\hat{s}_{\gamma} = \frac{\hbar}{2} \hat{\sigma}_{\gamma}$  with  $\hat{\sigma}_{\gamma}$  as Pauli matrices. Indices  $\alpha, \beta$  represent Cartesian coordinates and  $\gamma$  appoints for the direction of spin (typically  $z$ -direction).  $V_c$  is the cell volume and  $N_{\mathbf{k}}$  is the number of sampled points in reciprocal space. Finally,  $f_{n\mathbf{k}}$  denotes the Fermi-Dirac distribution and  $\hbar\omega$  corresponds to the external perturbation frequency. The static case is calculated by taking the limit  $\omega \rightarrow 0$  in the above formula.

## C. Additional figures

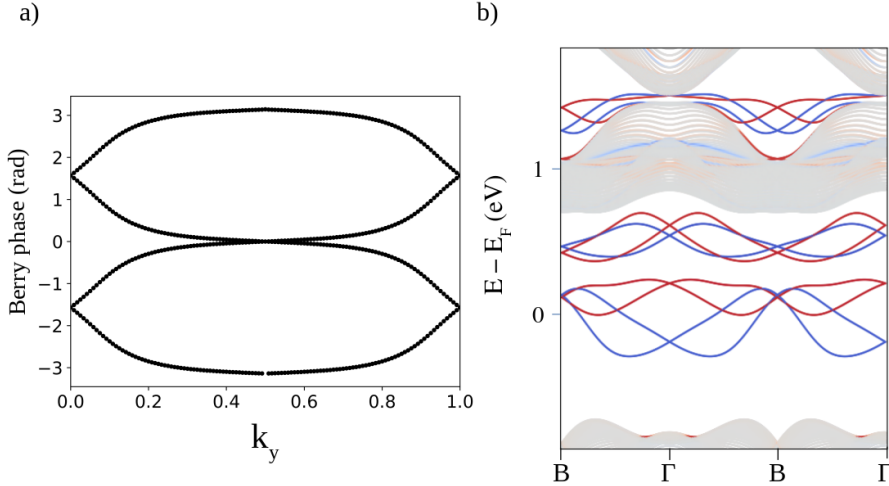


Figure S.1: a) Evolution of the Wannier charge center (Wilson loop) along the  $k_y$  direction. b) Energy dispersion of the ribbons for the Wannier interpolated 12B model for a cut along the  $y$  direction in real space. Red-colored and blue-colored bands represent states that are localized at the upper and lower edges of the ribbon, respectively.

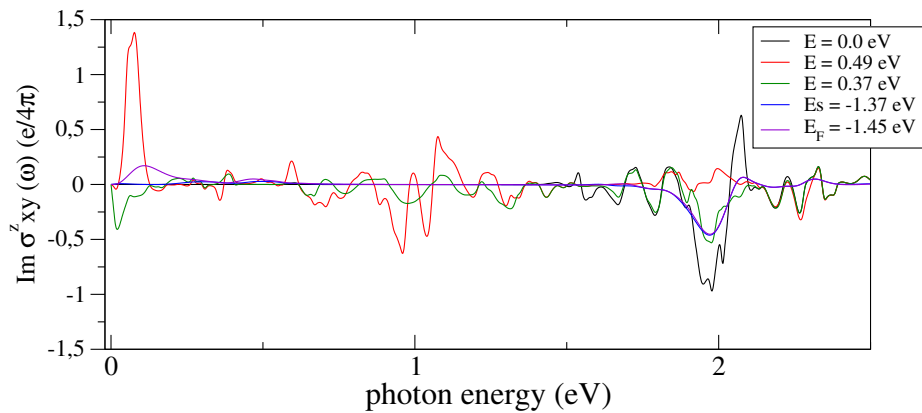


Figure S.2: Imaginary part of the spin Hall conductivity  $\sigma_{xy}^z$  as a function of external photon energy for different values of the chemical potential.

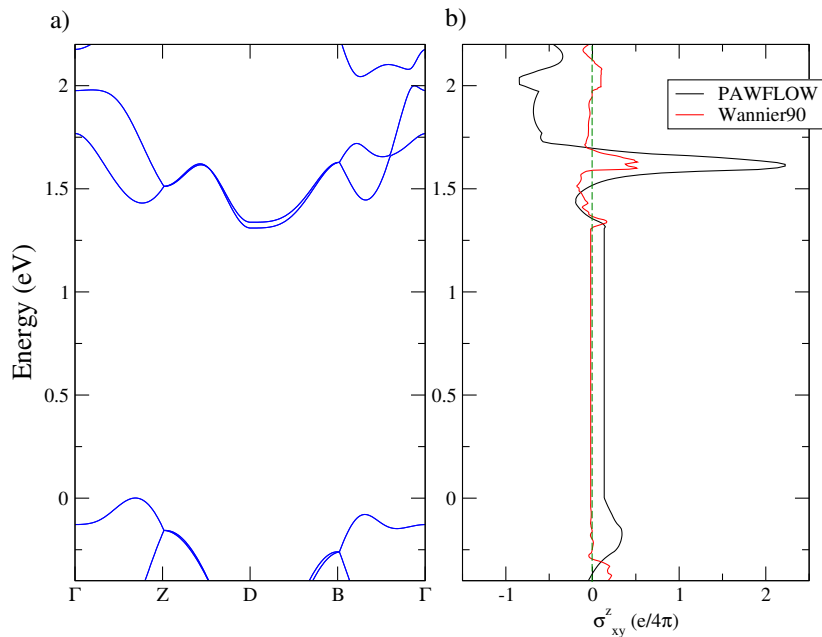


Figure S.3: a) Electronic bands of monolayer PdSe<sub>2</sub> from the PAWFLOW code. b) Static spin Hall conductivity  $\sigma_{xy}^z$  as a function of the Fermi level for two different codes.

## References

- [1] Luis Elcoro, Zhida Song, and B. Andrei Bernevig. Application of induction procedure and Smith decomposition in calculation and topological classification of electronic band structures in the 230 space groups. *Phys. Rev. B*, 102:035110, Jul 2020.
- [2] Junfeng Qiao, Jiaqi Zhou, Zhe Yuan, and Weisheng Zhao. Calculation of intrinsic spin Hall conductivity by Wannier interpolation. *Phys. Rev. B*, 98:214402, Dec 2018.