## SUPPLEMENTARY MATERIALS

## Theoretical prediction of superconductivity in two-dimensional MXenes of molybdenum carbides

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The structures and energy comparison of  $Mo_3C_2$  and  $Mo_3C_3$  systems with different configurations are shown in Figs. S1 and S2, respectively. As we can see, the most stable configurations of  $Mo_3C_2$  and  $Mo_3C_3$  are both 'H-1', which can be seen as the 'AA' stacking of H-Mo<sub>2</sub>C.

In addition, we also investigate the effects of biaxial stretching on the superconducting related physical quantities of four our studied 2D  $Mo_x C_y$ , which can be seen in Fig. S3.

As shown in Fig. S4, we explore the biaxial stretching limit of T-Mo<sub>2</sub>C, H-Mo<sub>2</sub>C, H-Mo<sub>3</sub>C<sub>2</sub>, and H-Mo<sub>3</sub>C<sub>3</sub>. Herein, we can find that when the  $\varepsilon$  is beyond the limit of biaxial stretching, the phonon dispersions of T-Mo<sub>2</sub>C, H-Mo<sub>3</sub>C<sub>2</sub>, and H-Mo<sub>3</sub>C<sub>3</sub> emerge the features of CDW.

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FIG. S1: The considered different structures of  $Mo_3C_2$  and  $Mo_3C_3$ .



FIG. S2: Energy comparison of different configurations for (a)  $Mo_3C_2$  and (b)  $Mo_3C_3$ . The configurations with the lowest total energy in different systems are set to 0.0 eV/atom. The dash line are guided to the eye.



FIG. S3: The logarithmically averaged phonon frequency  $\omega_{log}$  (red),  $N(E_F)$  (blue),  $T_c$  (red), and EPC constant  $\lambda$  (blue) under different biaxial tensile strains for (a)-(b) T-Mo<sub>2</sub>C, (c)-(d) H-Mo<sub>2</sub>C, (e)-(f) H-Mo<sub>3</sub>C<sub>2</sub>, and (g)-(h) H-Mo<sub>3</sub>C<sub>3</sub>.



FIG. S4: The phonon spectra of (a) T-Mo<sub>2</sub>C, (b) H-Mo<sub>2</sub>C, (C) H-Mo<sub>3</sub>C2, and (d) H-Mo<sub>3</sub>C<sub>3</sub> beyond the limit of biaxial stretching.