

Electronic Supplementary Information

for

Symmetry Breaking Charge Transfer and Intersystem Crossing in Copper Phthalocyanine Thin Films

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The excitation density for TA measurements is calculated as follows:

$$\rho = \alpha \cdot \frac{P_{in}}{E_{\gamma}} [cm^{-3}] \quad (1)$$

Where the absorption coefficient, α (Equation 2), depends on the measured absorption, which is wavelength dependent, *Abs*, in O.D., see Figure 1, and on the thickness of the sample, in this case, 50 nm:

$$\alpha = \frac{Abs(\lambda)}{l} \cdot 10^7 [cm^{-1}] \quad (2)$$

The incident power per pulse, P_{in} (Equation 3) is calculated from the measurement of the incident pump power P_{pump} , adjusted with a gradient neutral density filter, the laser repetition rate = 1 kHz, and the measured pump-beam radius, r :

$$P_{in} = \frac{P_{pump}[\mu W]}{RR \cdot \pi r^2} [J \cdot cm^{-2} \text{ per pulse}] \quad (3)$$

The photon energy, E_{γ} (Equation 4), for each wavelength measured in nm, using Planck's constant and the speed of light:

$$E_{\gamma} = h \cdot \frac{c}{\lambda} [J/\gamma] \quad (4)$$

Table S1. Excitation density at each excitation wavelength for the α -CuPc thin film.

Wavelength (nm)	580	620	680	720
Ex. Density (cm ⁻³)	1·10 ¹⁸	1·10 ¹⁸	8·10 ¹⁷	9·10 ¹⁷

Table S2. Excitation density at each excitation wavelength for the β -CuPc thin film.

Wavelength (nm)	580	620	680	720
Ex. Density (cm^{-3})	$2 \cdot 10^{18}$	$2 \cdot 10^{18}$	$7 \cdot 10^{18}$	$4 \cdot 10^{18}$

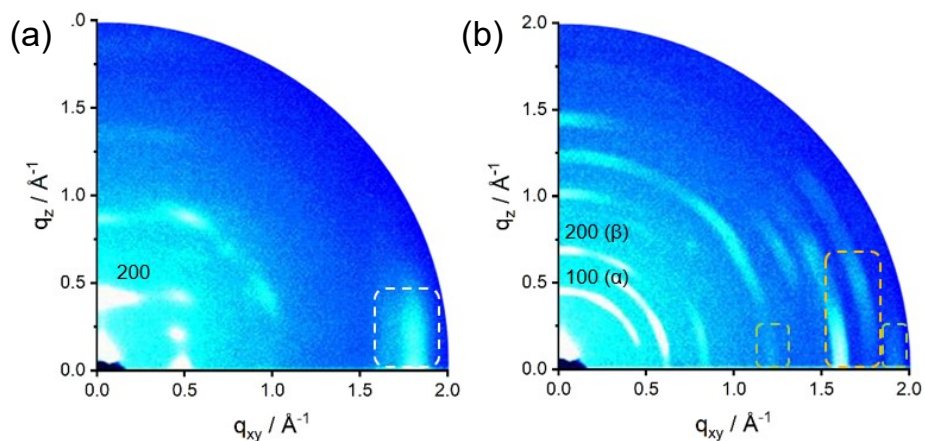


Figure S1. GIWAXS patterns of (a) α -CuPc (as-deposited) and (b) β -CuPc (annealed) thin films. Only the major reflection of the phases are assigned by the Miller index, while dashed boxes indicate reflections related to intracolumnar CuPc packing (white – α -herringbone, orange – α -brickstone, green – β -phase).

The as-deposited film reveals a characteristic α -herringbone structure (J. Am. Chem. Soc. 2012, 134, 14302–14305) with CuPc molecules arranged perpendicular to the substrate (Figure S1a). As indicated by the wide-angle reflections (white dashed box in Figure S1a), the molecules are in-plane shifted towards each other resulting in an interplanar distance of 3.45 \AA and Cu-Cu distance of 3.75 \AA . After annealing, the scattering intensities form arcs due to a broader orientation of the crystallites towards the surface (Figure S1b). Furthermore, the coexistence of two phases is identified based on the assignment of the reflections. New wide-angle reflections (orange dashed box in Figure S1b) are characteristic for the α -brickstone structure with interplanar and Cu-Cu distances of 3.49 \AA and 3.85 \AA , respectively, as well as a CuPc tilt angle of 75° with respect to the surface. Additionally, reflections for the β -phase are found in which the CuPc adapt a larger molecular displacement in the stacks resulting in Cu-Cu distance of 4.95 \AA and interplanar distance 3.30 \AA (green boxes in Figure S1b).

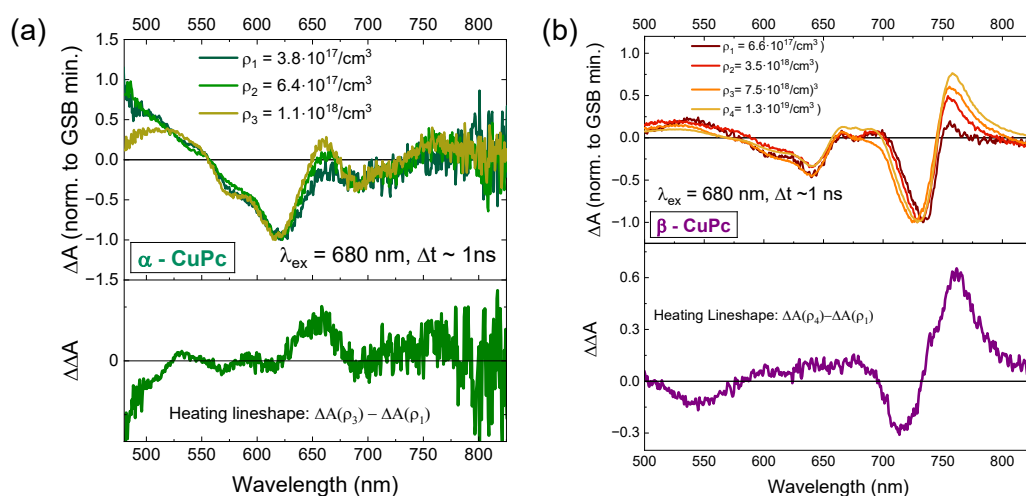


Figure S2. TA measurements with increasing excitation density are used to elucidate the ΔA signal that arises due to local heating effects for (a) α -CuPc and (b) β -CuPc thin films.

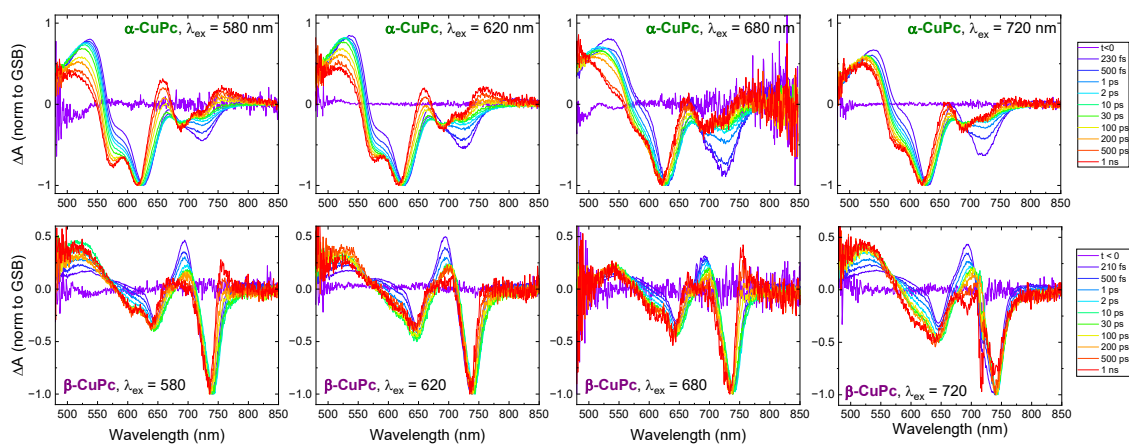


Figure S3. Normalised spectral traces at selected delay times for α -CuPc and β -CuPc thin films indicate that heating artefacts, evidenced by the blue shifting of the GSB and appearance of other features shown in Figure S1, appear later delay times (> 100 ps).

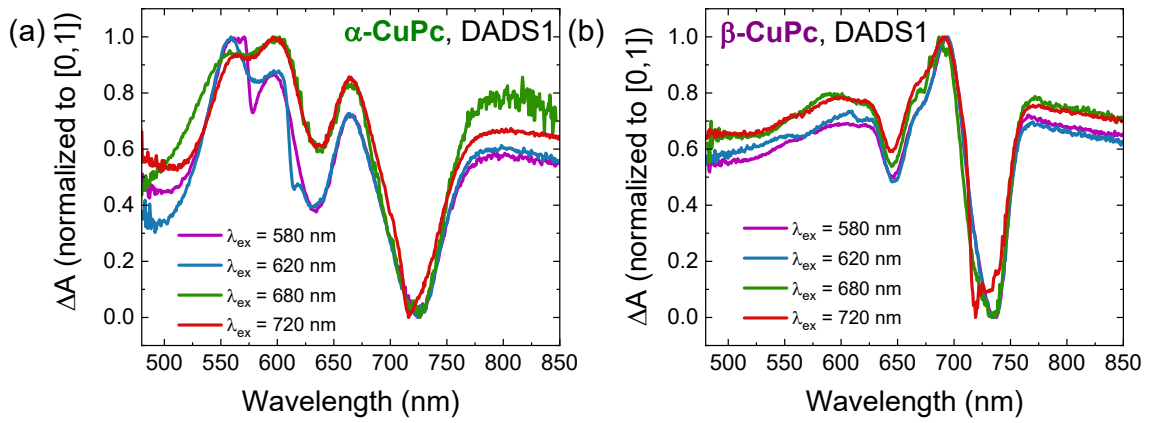


Figure S4. Comparison of Decay-Associated Difference Spectra for (a) α -CuPc and (b) β -CuPc thin films show an excitation dependent lineshape between 550-600 nm for α -CuPc that does not appear for β -CuPc

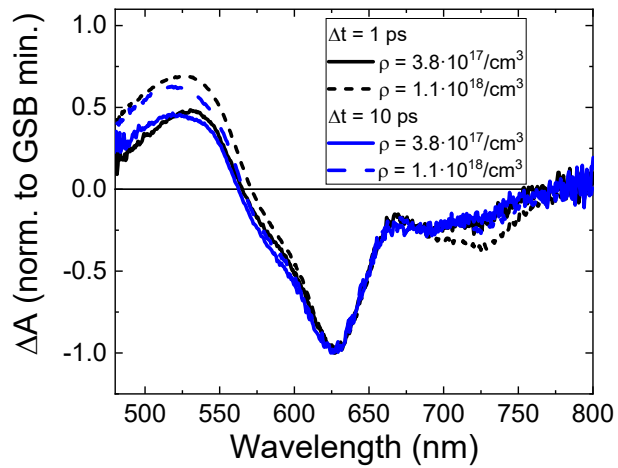


Figure S5. The spectral lineshape at $\Delta t = 1$ ps and 10 ps for α -CuPc at different excitation densities are similar, further confirming that the excitation wavelength dependent feature is not a heating artefact.