Electronic Supplementary Material (ESI) for Nanoscale. This journal is © The Royal Society of Chemistry 2023

Insulator-to-Metal Phase Transition in a Few-Layered MoSe₂ Field Effect Transistor

Nihar R. Pradhan^{1,2,*}, Carlos Garcia², Bhaswar Chakrabarti³, 4, Daniel Rosenmann³, Ralu Divan³,

Anirudha V. Sumant³, Suzanne Miller³, David Hilton⁵, Denis Karaiskaj⁶, and Stephen A. McGill^{2*} ¹Layered Materials and Device Physics Laboratory,

Department of Chemistry, Physics and Atmospheric Science,

Jackson State University, Jackson, MS 39217, USA

²National High Magnetic Field Laboratory, Tallahassee, FL 32310, USA

³Center for Nanoscale Materials, Argonne National Laboratory, 9700 S-Cass Avenue, Lemont, IL-60439, USA

⁴Department of Electrical Engineering, Indian Institute of Technology Madras, Chennai, Tamil Nadu- 600036, India

⁵Department of Physics, Baylor University, One Bear Place 97316, Waco, TX76798-7316, USA

⁶Department of Physics, University of South Florida, Tampa, FL 33620, USA

(Dated: December 24, 2022)

SUPPLEMENTARY INFORMATION Ι.

Figure S1: The conductivity data vs temperature in semi-logarithmic scale is plotted in Figure 1. This is the same data shown in the main body in Fig. 3(a). This semi-logarithmic plot clearly shows the insulating behavior of the MoSe₂ at low applied gate voltage, V_{bg} , where the conductivity decreases with decreasing temperature. When the gate voltage increased the slope of the conductivity changed, and above a certain threshold of applied voltage (marking a critical density of charge carriers induced by the applied gate voltage) the slope of the conductivity data changed sign compared to the insulating phase conductivity.



FIG. S1: Conductivity measured using 4-terminal method plotted as function of temperature in semi-logarithmic scale. The sample became too resistive to measure below 50 K for $V_{bq} = 65$ V

^{*}Electronic address: Corresponding_Email:nihar.r.pradhan@jsums.edu;Email:mcgill@magnet.fsu.edu



FIG. S2: (left) Conductivity, σ , curves normalized by the critical conductivity $\sigma_c(T, V_{bg} = 95 \text{ V})$, depicted by the dashed line) at the critical carrier density $n_c = n(V_{bg} = 95 \text{ V})$ as functions of the temperature T. (right) Plotting the rescaled conductivity, σ/σ_c as a function of the rescaled temperature, T/T_0 , showed no collapse of the conductivity curves to a single line that would indicate a quantum-critical aspect of the phase transition.

Figure S2: We plotted the temperature-dependent conductivity, σ , by normalizing all the curves with the critical conductivity value, σ_c , shown in Fig. 2 (left). Critical conductivity σ_c is represented by a dashed line which separates the insulating branch of the conductivity (below σ_c) and the metallic branch of the conductivity (above σ_c). In a quantum phase transition (QPT), the material undergoes a phase transition at T=0 K. The QPT occurs at the quantum critical point, which is where quantum fluctuations drive the transition to diverge and become scale-invariant in space and time. Experimentally, the two branches of the conductivity should scale with a temperature parameter, T_o , in such a way that

$$T_o \propto \delta n^{z\nu} \tag{1}$$

where $\delta n = \left(\frac{n-n_c}{n_c}\right)$. z and ν are the dynamic and correlation exponents, respectively. This suggests that the two branches of the conductivity should collapse into a single curve with appropriate T_o values. As demonstrated from the fitting of Fig. 2 (right), we have not established any collapse of the conductivity curves on either branch, which indicated that our phase transition did not exhibit quantum criticality.