

Supplementary information of:

Interface contact and modulated electronic properties by in-plane strains in graphene-MoS₂ heterostructure

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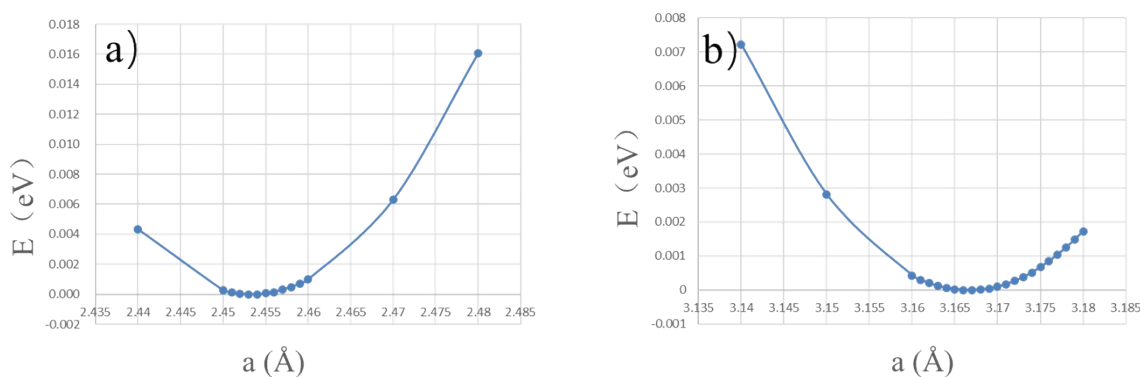


FIG. S1. Evolution of the total energy as a function of lattice parameter of (a) graphene and (b) MoS₂, taking the origin at the lowest energy configuration.

Fig. S1 shows the total energy curves of isolated graphene and MoS₂ monolayer. As shown, the lattice parameter of graphene and MoS₂ are 2.454 Å and 3.167 Å.

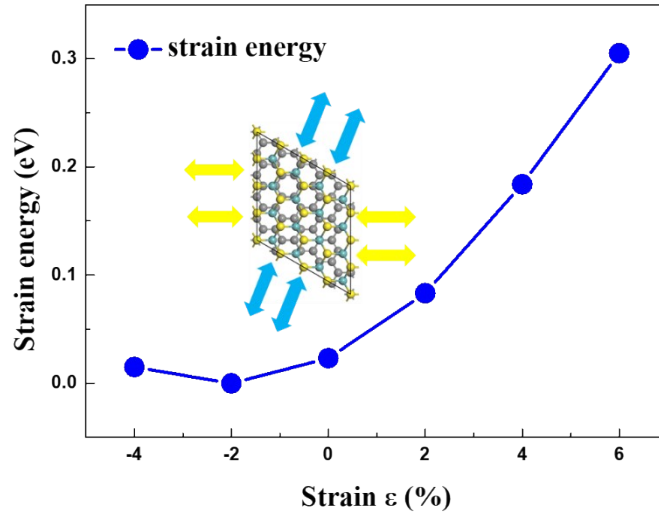


FIG. S2. The in-plane biaxial strain applied to Gr-MoS₂ heterojunction, and the relationship of strain energy with the biaxial strains.

The strain energy is checked here to ensure the strains considered are all within the elastic limit. The strain energy is given by the formula: $E = (E_S - E_U)/n$, where E_S and E_U are the total energy of strained and unstrained heterostructure, respectively, and n indicates the atom number in the supercell. Fig. S2 exhibits the relationship of the strain energy E with the imposed biaxial strain. In the biaxial strain range of $-4\% \sim 6\%$, it is obviously that the strain energy shows a characteristic of a quadratic function, which indicates that the mechanical strain falls within the elastic limit range that the deformed structure can restore to its initial state when the mechanical strains are removed.

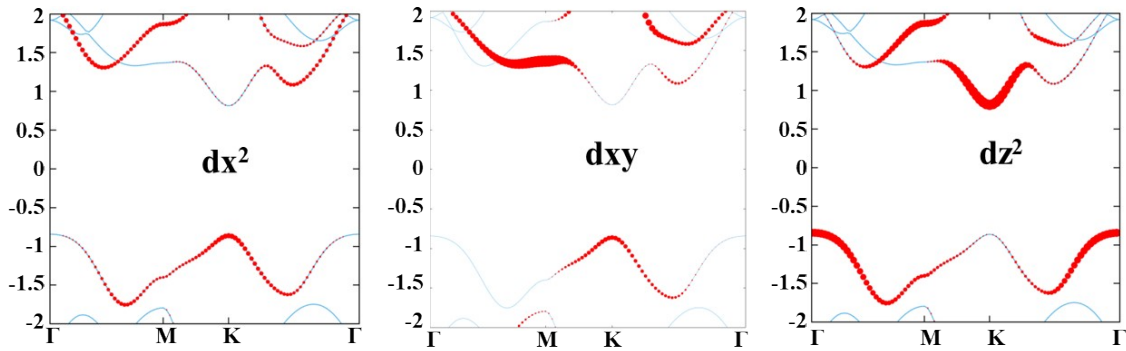


FIG. S3. The dx^2 , dxy , dz^2 orbitals band structures of MoS₂.

Fig. S3 shows the dx^2 , dxy , dz^2 orbitals band structures of MoS₂. It can be seen that the valence band minimum (VBM) of MoS₂ is located at the highly symmetric point K, which is contributed by the dz^2 orbital. Relatively, the conduction band maximum (CBM) is relatively complex, the CBM originates from the competition between the conduction band at K point contributed by dx^2 orbital, and the conduction band at Γ point contributed by dz^2 orbital. For MoS₂ monolayer with direct bandgap, CBM and VBM are both located at K.

While after adsorbed on graphene, the interaction in the direction perpendicular to the horizontal plane from graphene will enhance the MoS₂ dz² orbit, so that the CBM at Γ assumes a leading position and the direct bandgap in MoS₂ transforms to indirect.

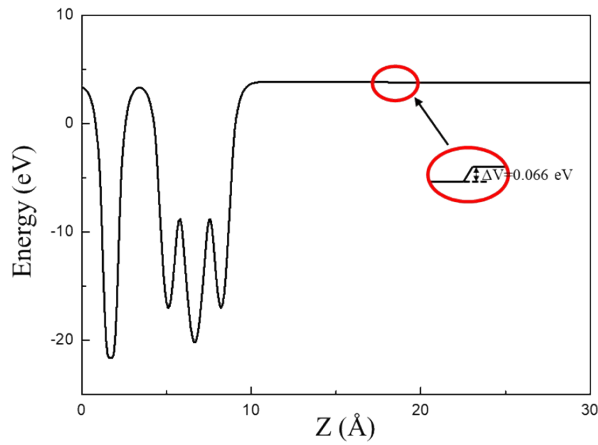


FIG. S4. The plane averaged electrostatic potentials of Gr-MoS₂ heterojunction along z-direction.

In Fig. S4, the average Coulomb potential for the Gr-MoS₂ system is shown. It can be seen that difference in work function in the case of Gr-MoS₂ and isolated graphene surface is defined as $\Delta V = W_{Gr-MoS_2} - W_{Gr}$, and it is about 0.066 eV. As $\Delta V > 0$, electrons will be transferred from graphene to MoS₂ surface, making the n-type contact.

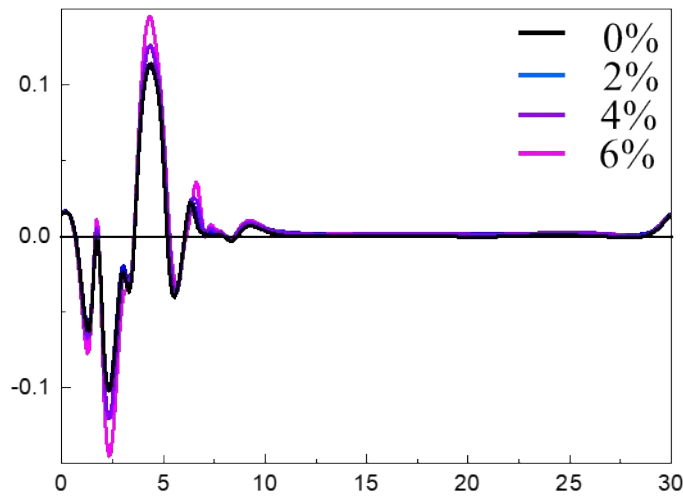


FIG. S5. Evolution of charge density differences under tensile strains.

Fig. S5 presents the charge density differences of Gr-MoS₂ bilayer under various in-plane tensile strains. It is clear that the electron redistribution mainly occurs within the strain. As mentioned, the electron accumulates and depletes on the side of MoS₂ and graphene layer, and it is observed that the degree of electron redistribution increases with increasing the strain, reaching a maximum value at 6%.

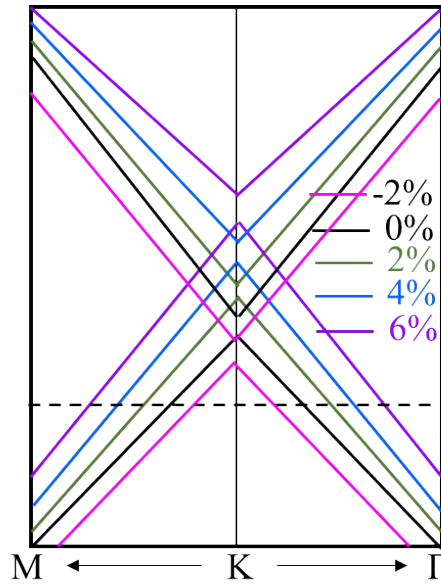


FIG. S6. Linear dispersion near the Fermi level at Dirac point of Gr-MoS₂ interface with tensile strain.

To describe the effect of the strain on the band gap of graphene in Gr-MoS₂ heterojunctions, in Fig. 6 we show the energy dispersion near the Fermi energy level at Dirac point for the VBM and the CBM in the band structures with different strain. As shown in Fig. S2, the total energy of Gr-MoS₂ junction is minimum at the equilibrium state of strain -2%, when graphene is subjected to the least external force. Therefore, for graphene, strain -2% is the most stable state. As mentioned above, the charge transfer from graphene to MoS₂, moves the Fermi level down below the Dirac point, leading to the hole doping in graphene in the Gr-MoS₂ heterojunction.

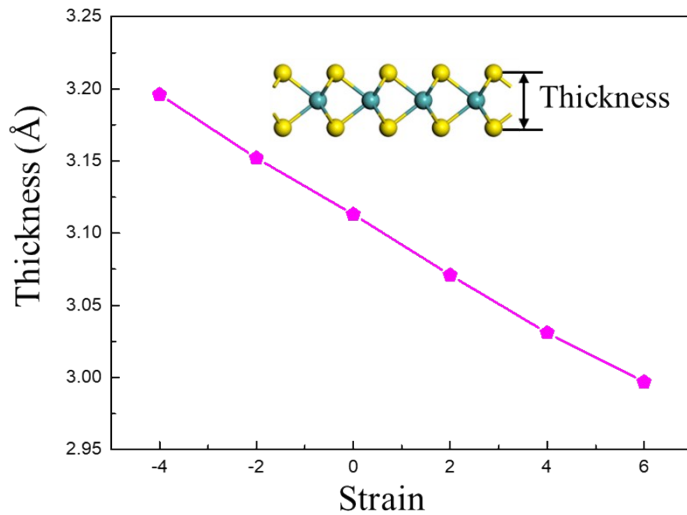


FIG. S7. The relationship of thickness of monolayer MoS₂ with the biaxial strains.

The thickness of layered MoS₂ changes with the application of strain, as shown in the Fig. S7. It can be seen that when compressive strain is applied, the thickness increases and Mo-S bond tends to the horizontal vertical direction, while when tensile strain is applied, its thickness decreases and Mo-S bond tends to the horizontal

direction. That is, the tensile strain in MoS₂ achieved by increasing lattice constant flattens the monolayer. Moreover, with tensile strain, the bond length of Mo-S is elongated, so that the ability of constraint valence electrons is weakened and then the band gap is reduced accordingly.