

Supporting Information for

**Three-dimensional lattice deformation of blue phase
liquid crystals under electrostriction**

Yuxian Zhang,^a Hiroyuki Yoshida,^{b*} Fan Chu,^a Yu-Qiang Guo,^a Zhou Yang,^c Masanori Ozaki^b and Qiong-Hua Wang^{a*}

a. School of Instrumentation and Optoelectronic Engineering, Beihang University, Beijing 100191, China.

b. Division of Electrical, Electronic and Infocommunications Engineering, Graduate School of Engineering, Osaka University, 2-1 Yamada-oka, Suita, Osaka 565-0871, Japan.

c. Department of Materials Science and Engineering, University of Science and Technology Beijing, Beijing 100083, China

Corresponding authors

H. Yoshida: yoshida@eei.eng.osaka-u.ac.jp

Q.-H. Wang: qionghua@buaa.edu.cn

Optical texture of the intermediate “quasi-tetragonal” phase

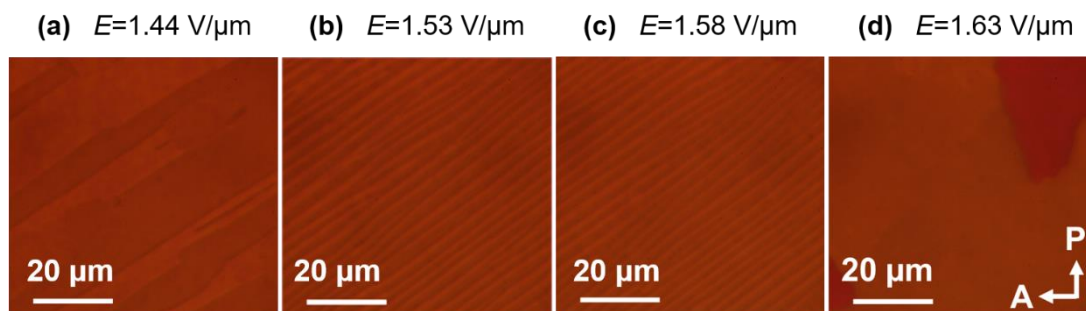


Fig. S1 POM images of the BP I₍₁₁₀₎ twinning sample, showing the phase transitions of orthorhombic ($E=1.44 \text{ V}/\mu\text{m}$)→quasi-tetragonal ($E=1.53\sim 1.58 \text{ V}/\mu\text{m}$)→tetragonal ($E=1.63 \text{ V}/\mu\text{m}$).

Discontinuous lattice elongation along the field

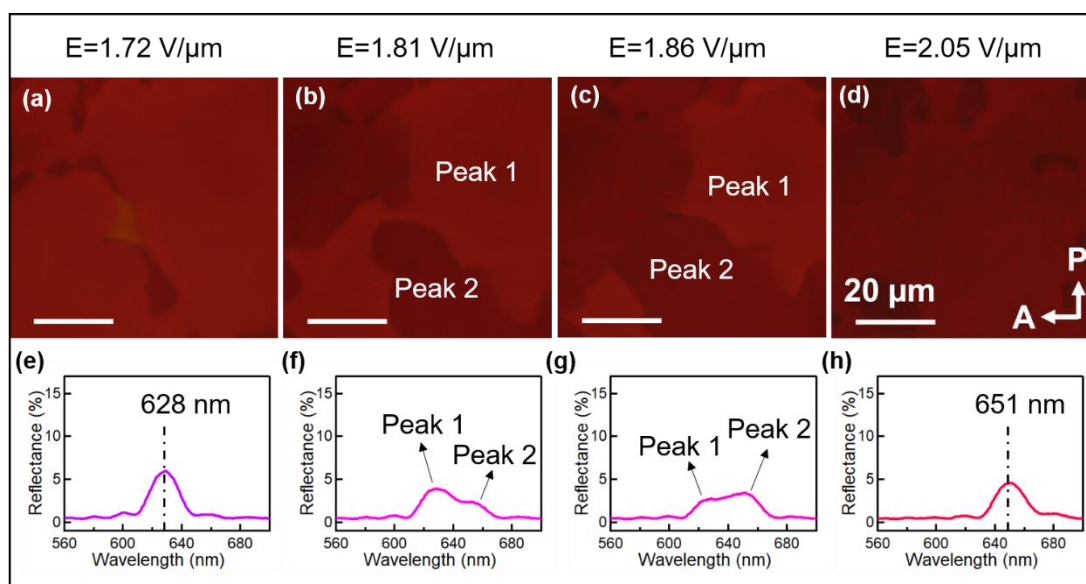


Fig. S2 POM images and reflection spectrum of BP I crystals at varied field intensities. The gradual transition of the POM texture and the discrete redshift of reflection peaks indicate the discontinuous elongation of BP I lattice along the field direction.

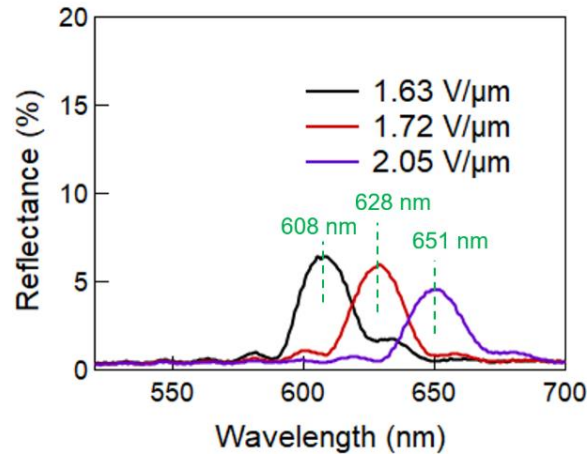


Fig. S3 Reflection spectrum of tetragonal BP I crystals respectively at $E=1.63$, 1.72 and 2.05 $\text{V}/\mu\text{m}$ (Cell thickness: $d=5.6$ μm). The Bragg peak wavelength at each field intensity is obtained by fitting Gaussian function.

Spectra of BP $I_{(110)}$ crystals confined in a 14.9 μm -thick cell

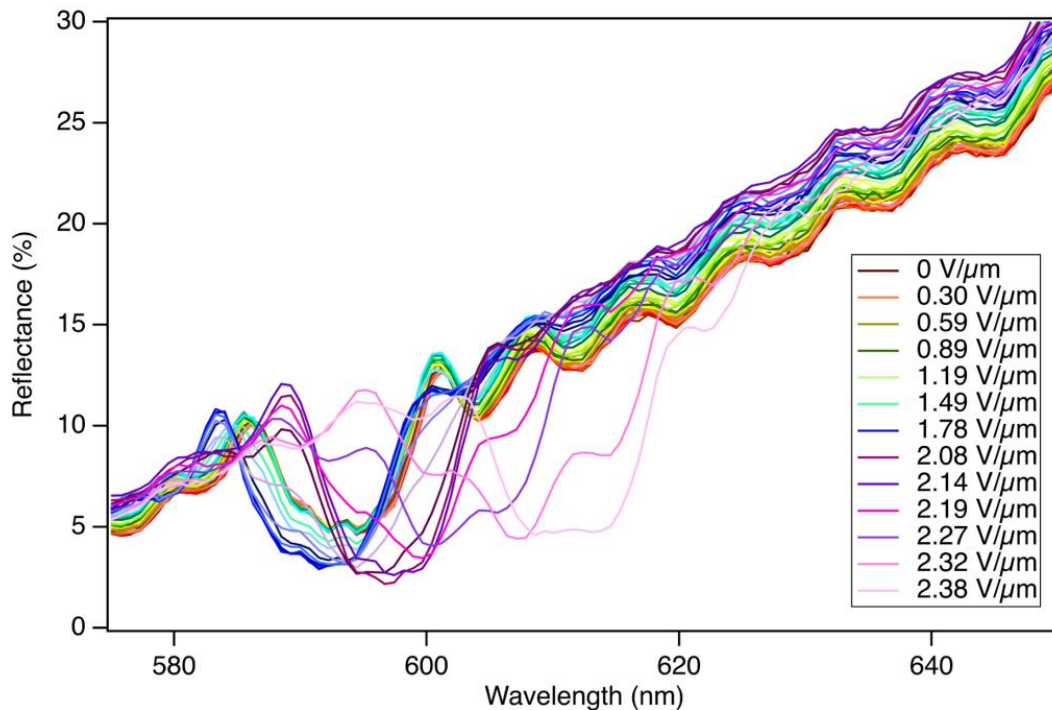
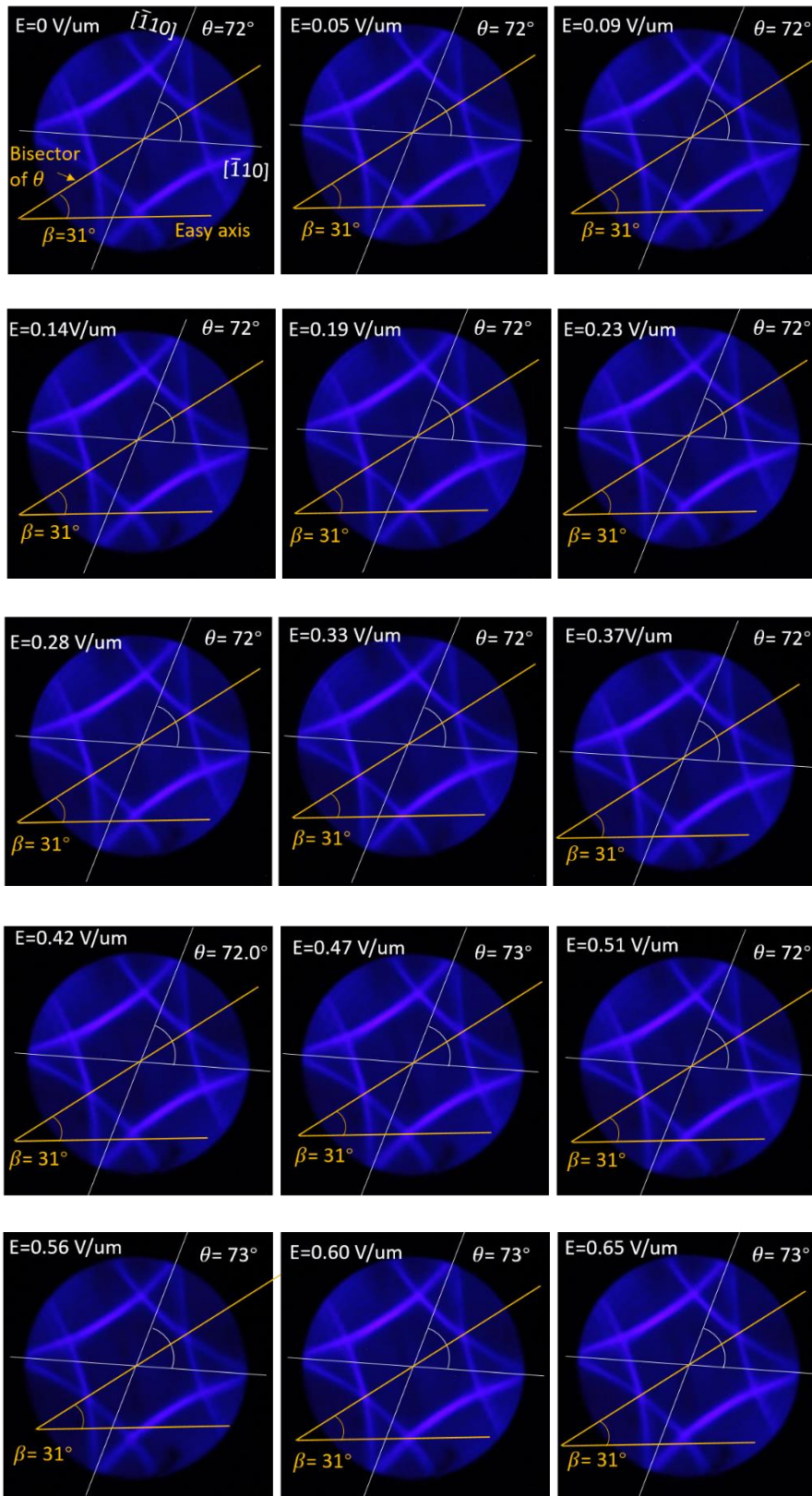
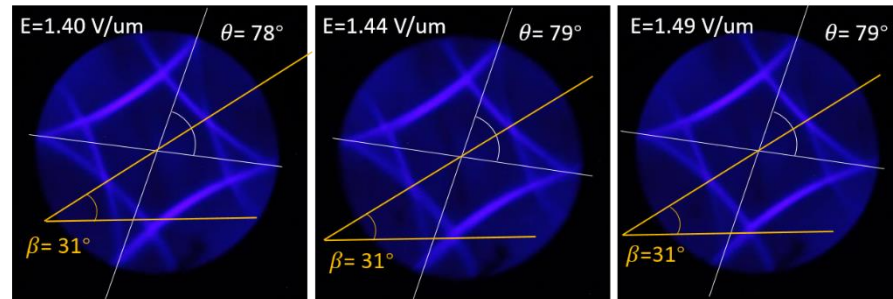
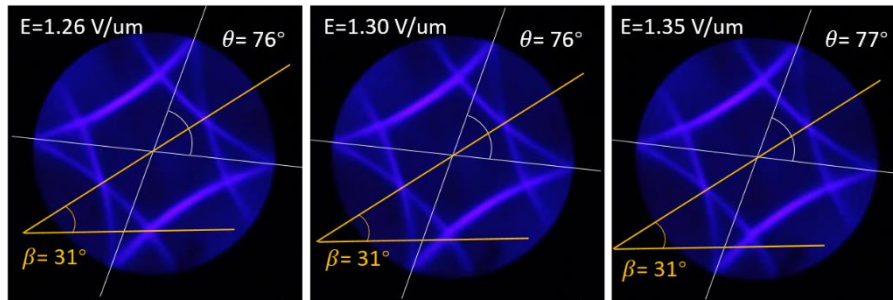
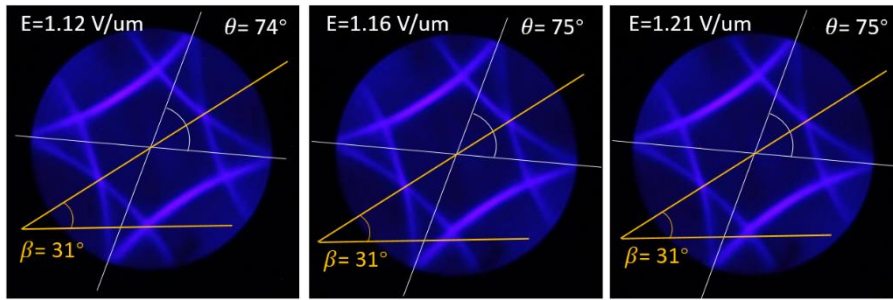
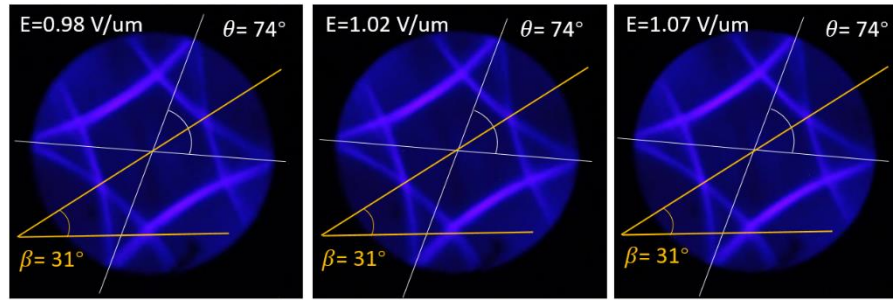
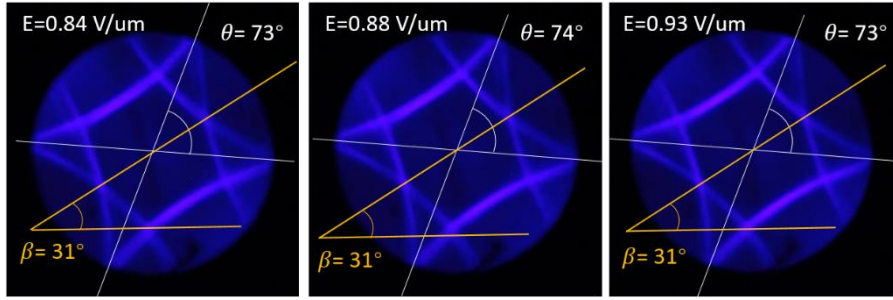


Fig. S4 Transmission spectra of the BP $I_{(110)}$ crystals confined in a 14.9 μm cell.

Raw Kossel pattern data under varying field intensities





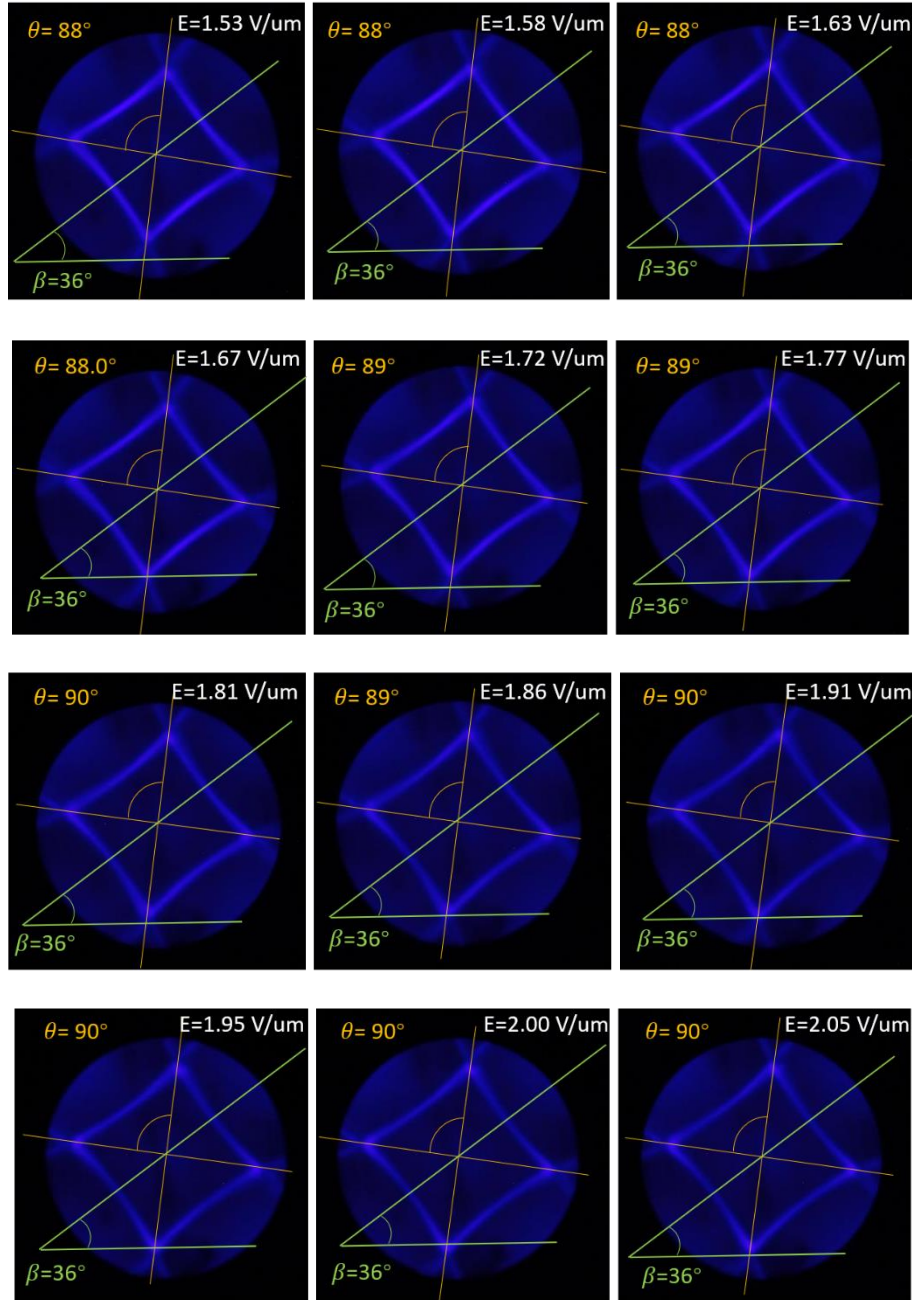


Fig. S5 Kossel diagrams of the BP $I_{(110)}$ twinned crystals under varied electric-field intensities, where the angle β indicates the lattice direction. The overlap angle θ is measured at each diagram to calculate a and b , the lattice constants perpendicular to the field.

Accuracy calculation of the lattice constant measurement

In this work, we use Equations (2)-(4) to calculate the three-dimensional lattice constants:

$$\left\{ \begin{array}{l} c = \lambda / n \quad (2) \\ b = \sqrt{\frac{V_0}{c \tan(\theta/2)}} \quad (3) \\ a = b \tan(\theta/2) \quad (4) \end{array} \right.$$

The main error comes from the overlap angle (θ) measurement that affect the accuracy of a and b in the field-perpendicular direction, and the field-induced variation of refractive index (n) of BPs that affect the accuracy of c along the field direction.

- (1) The overlap angles in Kossel diagram are measured by the screen protractor as follows in Figure S5, with the accuracy of 1' for each $[\bar{1}10]$ axis. Considering the $[\bar{1}10]$ axes are manually drawn, we set the error to $\pm 0.5^\circ$ for the $[\bar{1}10]$ axis and thus, the accuracy of measurement overlap angle θ is within $\pm 1^\circ$.

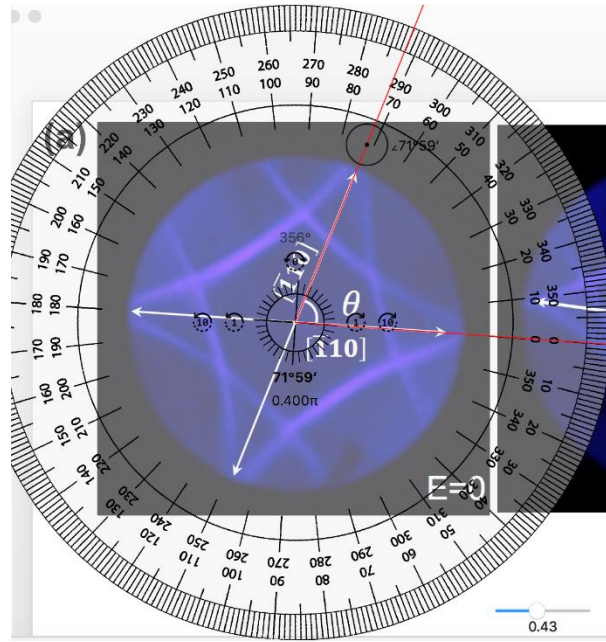


Fig. S6 Measuring the overlap angle θ in a Kossel diagram.

During the electrostriction, θ is measured as 72° to 90° and $V_E = V_0 = c_0^3 \tan(\theta_0/2)$.

For $\theta = 72^\circ$ at $E=0$, the estimated values perpendicular to the field are $a=282$ nm, $b=386$ nm.

If we use $\theta = 71^\circ$, the values turn to be $a=275$ nm, $b=386$ nm. ($\Delta_{\max} = -7$ nm)

If we use $\theta = 73^\circ$, the values turn to be $a=286$ nm, $b=386$ nm. ($\Delta_{\max} = 4$ nm)

For $\theta=89.5^\circ$ at $E=2.05 \text{ V}/\mu\text{m}$, the estimated values are $a=318 \text{ nm}$, $b=321 \text{ nm}$.

If we use $\theta=88.5^\circ$, the values turn to be $a=315 \text{ nm}$, $b=323 \text{ nm}$. ($\Delta_{\text{max}}= -3 \text{ nm}$)

If we use $\theta=90.5^\circ$, the values turn to be $a=320 \text{ nm}$, $b=317 \text{ nm}$. ($\Delta_{\text{max}}= -4 \text{ nm}$)

For $\theta=79^\circ$ at $E=1.44 \text{ V}/\mu\text{m}$, the estimated values are $a=301 \text{ nm}$, $b=362 \text{ nm}$.

If we use $\theta=78^\circ$, the values turn to be $a=296 \text{ nm}$, $b=366 \text{ nm}$. ($\Delta_{\text{max}} = -5 \text{ nm}$)

If we use $\theta=80^\circ$, the values turn to be $a=301 \text{ nm}$, $b=359 \text{ nm}$. ($\Delta_{\text{max}}= 3 \text{ nm}$)

Because c is calculated from Bragg peak wavelength, the error comes from the angle measurement is within $\pm 7 \text{ nm}$.

(2) For the refractive index, since the field variation is less than 0.01 and $n=1.58$ was used in the main text, here we use $n=1.59$ and $n=1.57$ for $c=\lambda/n$:

At $E=0$, $\lambda=610 \text{ nm}$, $n=1.58$ and $c=386 \text{ nm}$;

if $n=1.59$, $c=384 \text{ nm}$

if $n=1.57$, $c=389 \text{ nm}$

($\Delta_{\text{max}}= 3 \text{ nm}$)

At $E=1.72 \text{ V}/\mu\text{m}$, $\lambda=628 \text{ nm}$, $n=1.58$ and $c=397 \text{ nm}$;

if $n=1.59$, $c=395 \text{ nm}$

if $n=1.57$, $c=400 \text{ nm}$

($\Delta_{\text{max}}= 3 \text{ nm}$)

At $E=2.05 \text{ V}/\mu\text{m}$, $\lambda=651 \text{ nm}$, $n=1.58$ and $c=412 \text{ nm}$;

if $n=1.59$, $c=409 \text{ nm}$

if $n=1.57$, $c=415 \text{ nm}$

($\Delta_{\text{max}}= 3 \text{ nm}$)

When we calculate a and b , we use the value of c directly, so the error caused by refractive index is within ± 3 nm.

To sum up, the accuracy for the lattice constant measurement is at least within ± 10 nm.