

Electronic Supplementary Information

**Large second harmonic generation in penta-CdO₂ sheet exfoliated
from its bulk phase**

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Table S1 The forty-five *d*¹⁰-TM oxides from the Materials Project (MP) Database. SG represents the space group.

Formula	MP-id	SG	Formula	MP-id	SG	Formula	MP-id	SG
PdO ₂	<i>P4₂/mnm</i>	1018886	CuO ₂	614565	<i>Fmmm</i>	ZnO	2133	<i>P6₃mc</i>
PdO	<i>P4₂/mmc</i>	1336	CuO ₂	600604	<i>Fmmm</i>	ZnO	1986	<i>F⁻⁴₃m</i>
PdO ₂	<i>Pn⁻³_m</i>	10729	CuO ₂	705439	<i>Pmmm</i>	ZnO ₂	8484	<i>Pa⁻³</i>
PdO	<i>Fm⁻³_m</i>	10728	Ag ₂ O	353	<i>Pn⁻³_m</i>	ZnO	2229	<i>Fm⁻³_m</i>
PdO	<i>Pm⁻³_m</i>	603251	Ag ₃ O ₄	1605	<i>P2₁/c</i>	ZnO ₂	1094003	<i>P⁻³_m1</i>
CuO	1692	<i>P4₂/mmc</i>	AgO	499	<i>Cccm</i>	CdO	1132	<i>Fm⁻³_m</i>
Cu ₂ O	361	<i>Pn⁻³_m</i>	AgO	1065190	<i>C2/c</i>	CdO ₂	2310	<i>Pa⁻³</i>
CuO	704645	<i>C2/c</i>	AgO	1066856	<i>C2/c</i>	CdO ₂	1096866	<i>P⁻³_m1</i>
CuO	1064456	<i>Cccm</i>	Ag ₂ O ₃	546190	<i>Fdd2</i>	CdO ₂	1096876	<i>Cm</i>
Cu ₄ O ₃	1478	<i>I4₁/amd</i>	AgO	1079720	<i>P2₁/c</i>	HgO	1224	<i>Pnma</i>
Cu ₈ O	31217	<i>Cmmm</i>	AgO	1288	<i>I4₁/a</i>	HgO	556859	<i>P2₁/m</i>

CuO	14549	$Fm\bar{3}m$	Ag ₂ O	7711	$P\bar{3}m1$	HgO	7826	$P3_121$
Cu ₈ O	704745	$Amm2$	Ag ₂ O ₃	11872	$Pn\bar{3}m$	HgO	1077107	$P3_221$
CuO ₂	1181499	$Cmcm$	AgO	8222	$F\bar{4}3m$	HgO ₂	557266	$Pbca$
CuO ₂	601195	$Cmcm$	Au ₂ O ₃	27253	$Fdd2$	Hg ₂ O	2278	$C2/m$

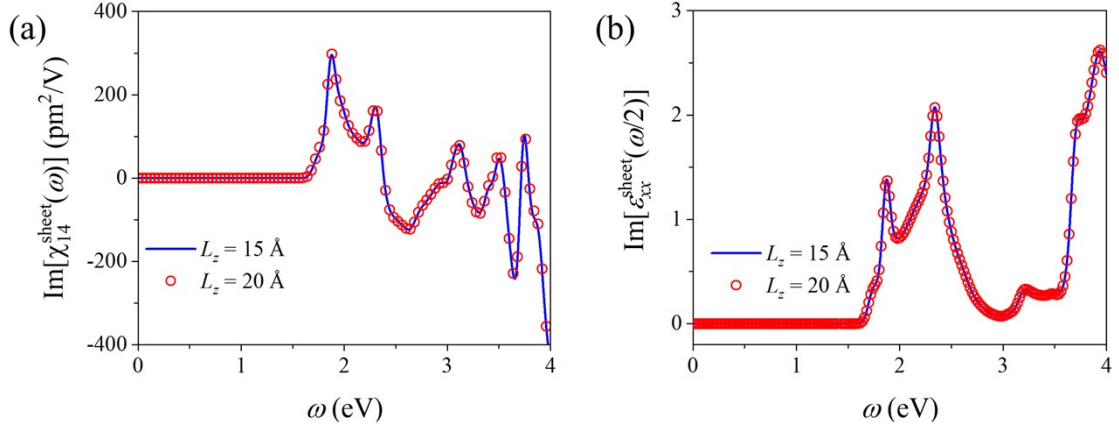


Fig. S1 Imaginary part of (a) the SHG susceptibility, and (b) the dielectric function of penta-CdO₂ obtained with different vacuum space L_z .

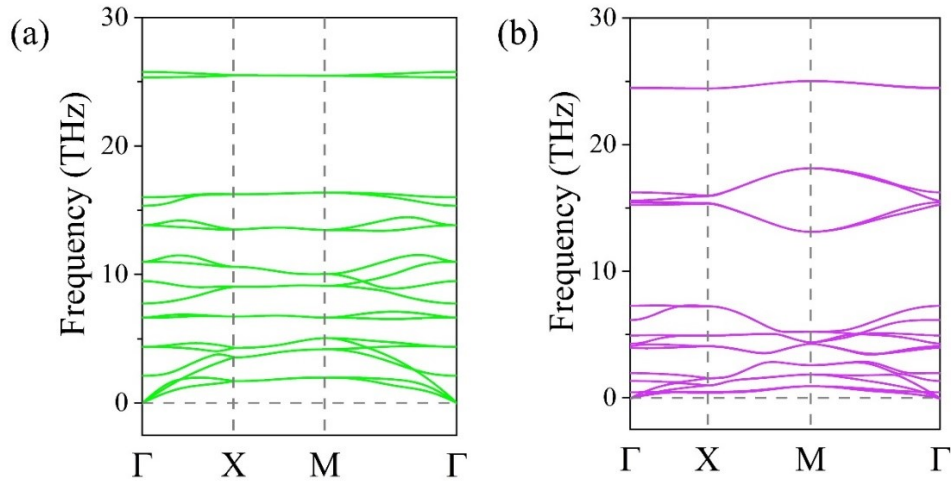


Fig. S2 Phonon spectra of (a) penta-ZnO₂ and (b) penta-HgO₂.

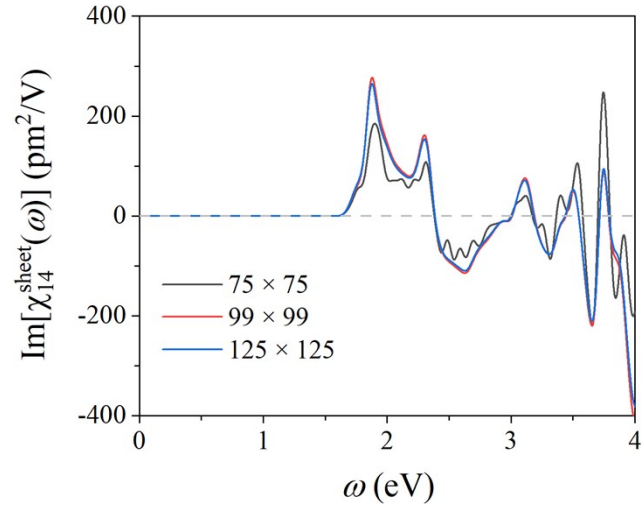


Fig. S3 Imaginary part $\text{Im}[\chi_{\text{sheet } 14}(\omega)]$ of the SHG susceptibility of penta- CdO_2 obtained with different k -mesh.

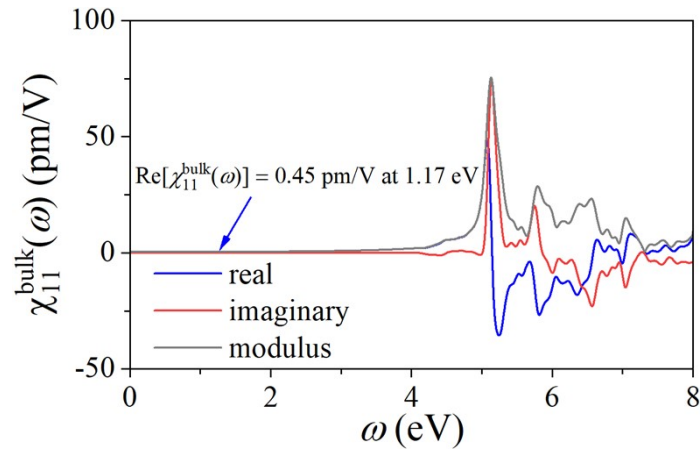


Fig. S4 SHG susceptibility $\chi_{\text{bulk } 11}(\omega)$ of KBBF.

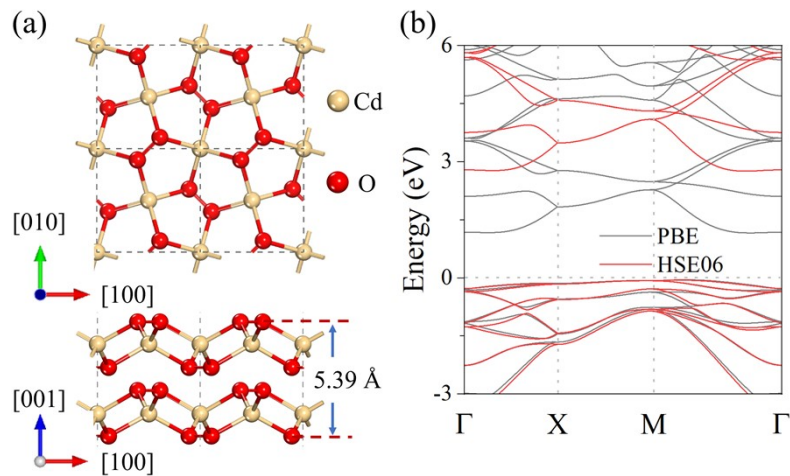


Fig. S5 (a) Top and side views of the optimized geometry, and (b) electronic band structure of bilayer penta-CdO₂.

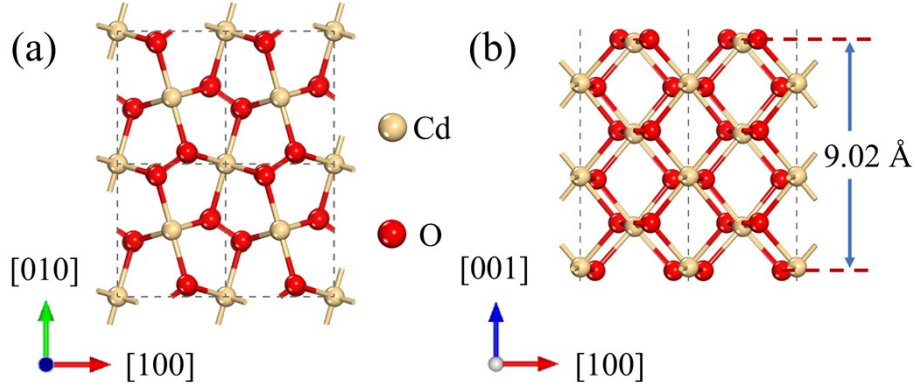


Fig. S6 (a) Top and (b) side views of the optimized geometry of trilayer penta-CdO₂.

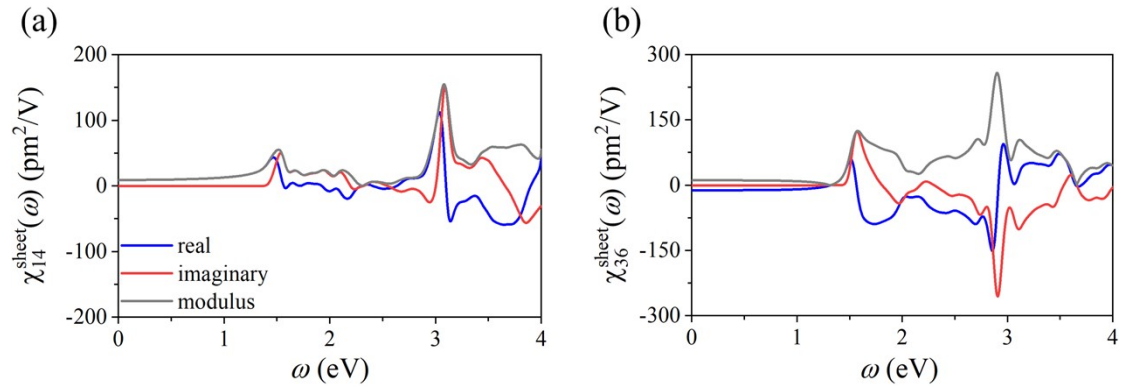


Fig. S7 SHG susceptibility of (a) $\chi_{14}^{\text{sheet}}(\omega)$ and (b) $\chi_{36}^{\text{sheet}}(\omega)$ of bilayer penta-CdO₂.

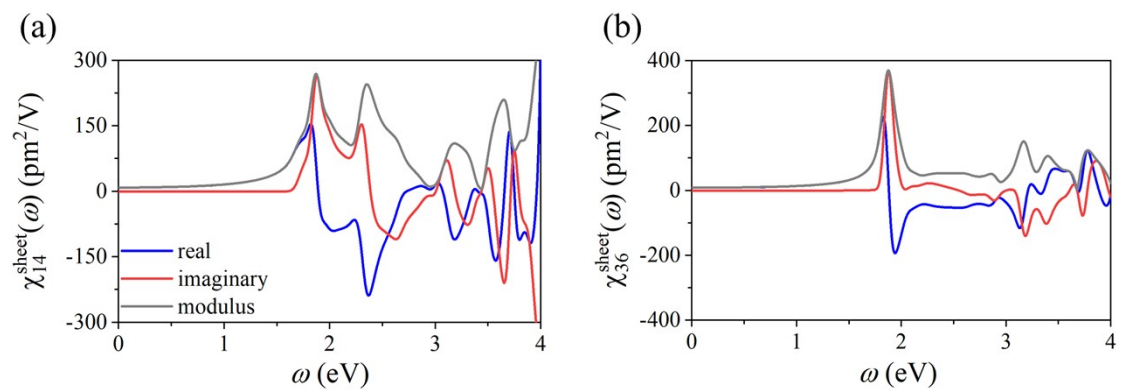


Fig. S8 SHG susceptibility (a) $\chi_{\text{sheet } 14}(\omega)$ and (b) $\chi_{\text{sheet } 36}(\omega)$ of penta-CdO₂.

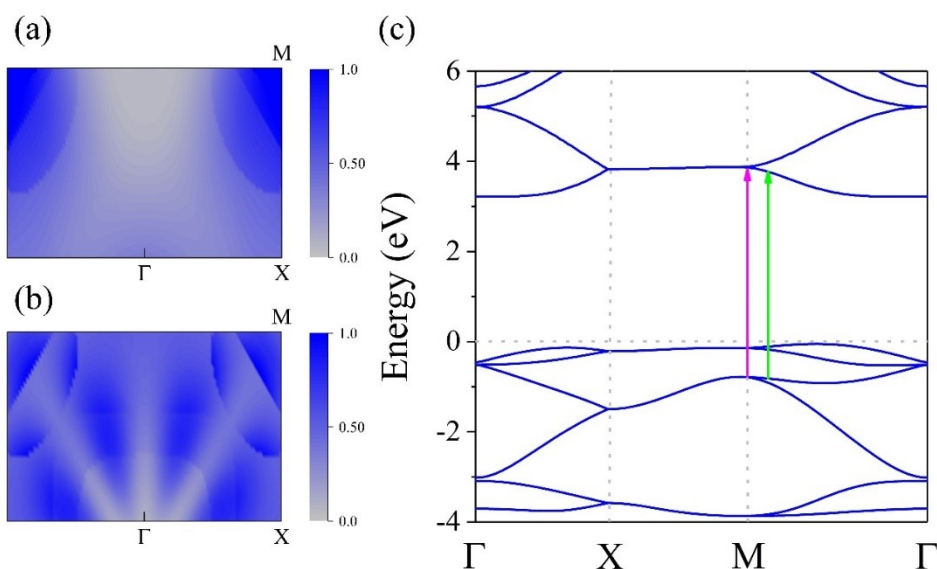


Fig. S9 Normalized k -point-dependent distributions of (a) $\text{Im}[\epsilon_{\text{sheet } xx}(\omega/2)]$ and (b) $\text{Im}[\epsilon_{\text{sheet } zz}(\omega/2)]$ at the peak value at $\omega/2 = 1.88$ eV, and (c) band structure of penta-CdO₂ at the PBE level with scissors correction.

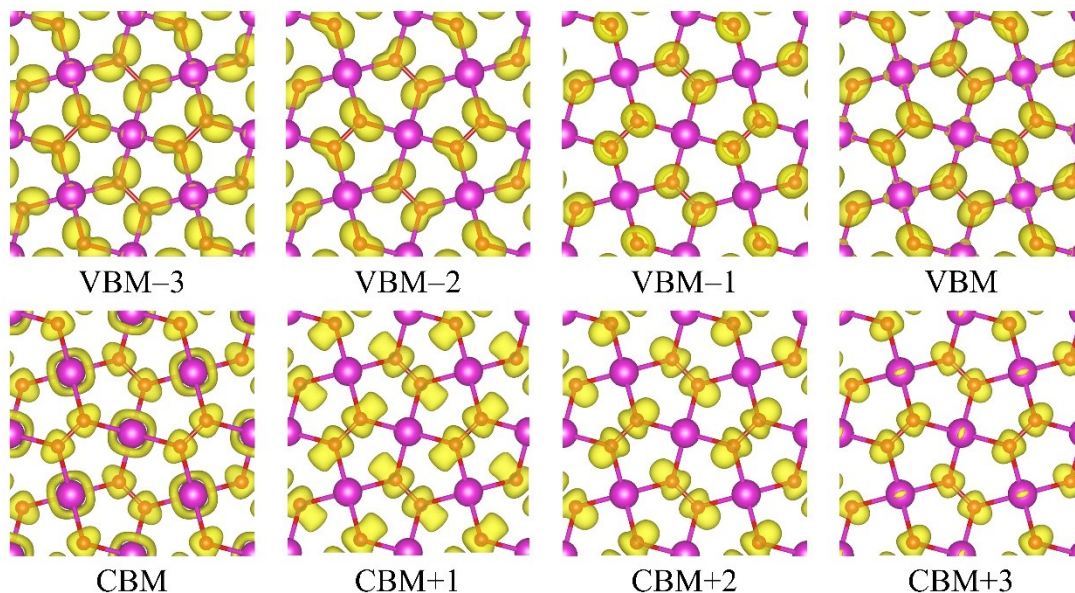


Fig. S10 Band-decomposed charge density distribution of the energy bands near the Fermi level of penta-CdO₂ (isosurface values = 0.005 \AA^{-3}).

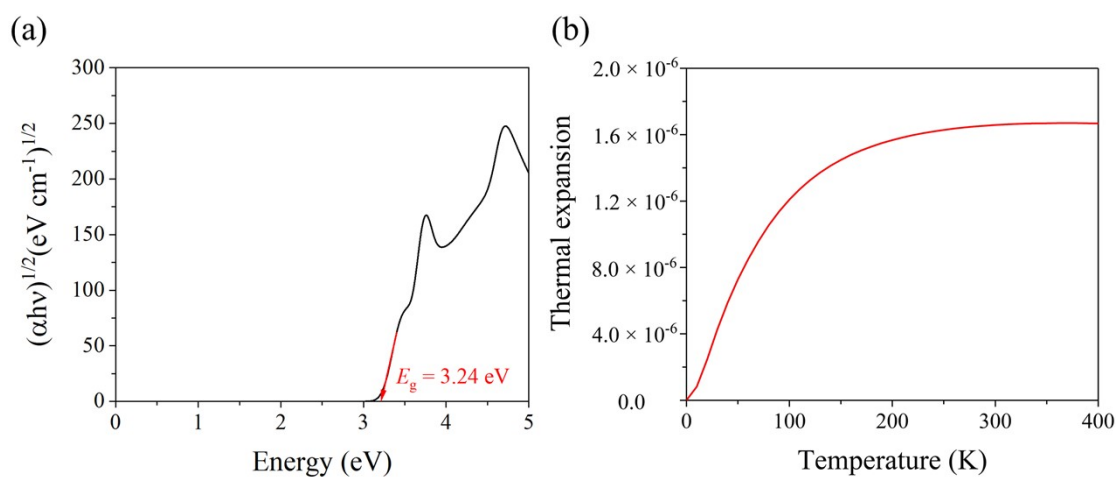


Fig. S11 (a) Optical band gap (E_g) determined from the Tauc plot, and (b) thermal expansion of penta- CdO_2 .

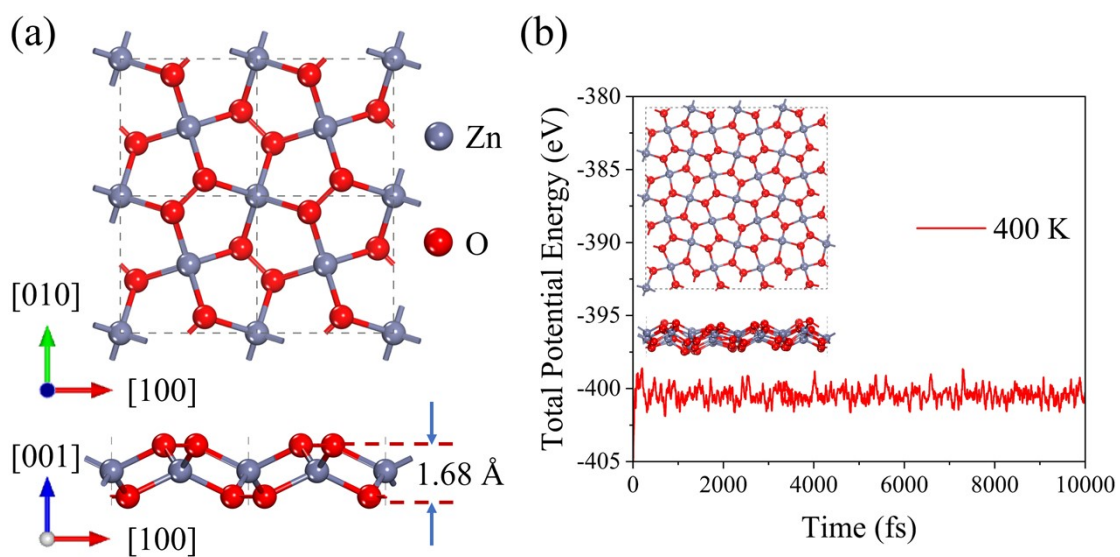


Fig. S12 (a) Top and side views of the optimized geometry of penta- ZnO_2 . (b) Total potential energy fluctuation with time during the AIMD simulation at 400 K.

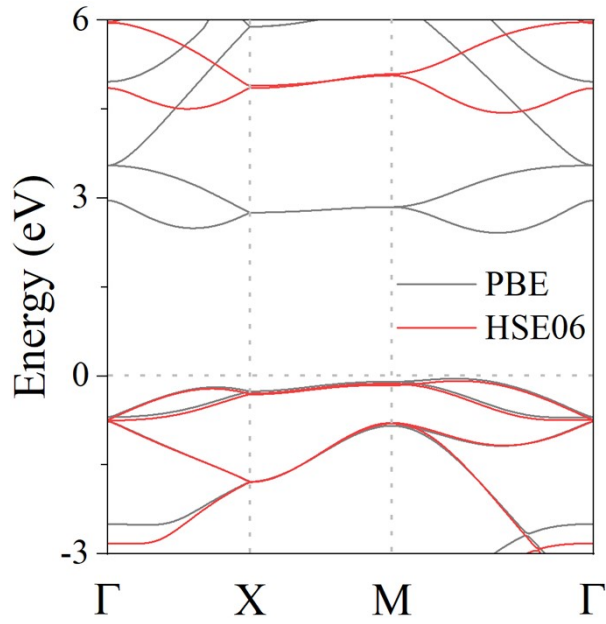


Fig. S13 Electronic band structure of penta-ZnO₂.

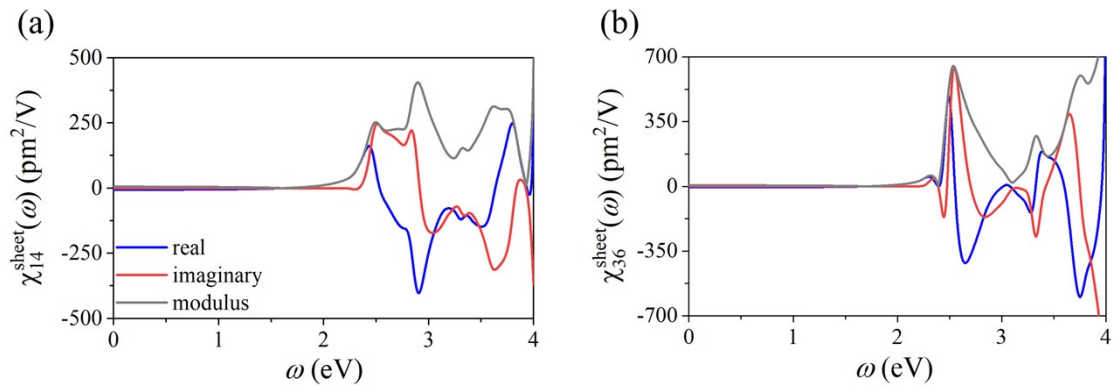


Fig. S14 SHG susceptibilities (a) $\chi_{14}^{\text{sheet}}(\omega)$ and (b) $\chi_{36}^{\text{sheet}}(\omega)$ of penta-ZnO₂.

We also find that another 2D d^{10} -TM oxide, penta-ZnO₂, can also exhibit large SHG response, as we expected. The fully optimized geometry of penta-ZnO₂ is presented in Fig. S12(a). Its unit cell contains two Zn and four O atoms, which occupy two nonequivalent Wyckoff positions, namely, 2b (0.500, 0.500, 0.500) and 4e (0.882, 0.618, 0.458), respectively. The lattice parameters are $a = b = 4.49$ Å. The penta-ZnO₂ sheet is also dynamically and thermally stable, as illustrated in Fig. S1(a) and Fig.

S12(b). The calculated linear elastic constants are $C_{11} = 42.50 \text{ N m}^{-1}$, $C_{12} = 3.00 \text{ N m}^{-1}$, and $C_{66} = 22.56 \text{ N m}^{-1}$, which meet the requirements of Born-Huang criteria ($C_{11} > 0$, $C_{11} > C_{12}$ and $C_{66} > 0$),⁵⁰ confirming that penta-ZnO₂ is mechanically stable. The band structures of penta-ZnO₂ at the HSE06 level and PBE level are given in Fig. S13, which show that penta-ZnO₂ is also semiconducting with an indirect band gap of 2.46 eV at the PBE level, and 4.52 eV at the HSE06 level, larger than the corresponding values of 1.45 and 3.28 eV, respectively, for penta-CdO₂. We further explore its potential for SHG performance because of the non-centrosymmetric and non-metallic features of penta-ZnO₂. We find that penta-ZnO₂ possesses strong in-plane and out-of-plane SHG response with large static SHG susceptibilities of $\chi_{\text{sheet } 14}(0) = \chi_{\text{sheet } 25}(0) = \chi_{\text{sheet } 36}(0) = -5.60 \text{ pm}^2/\text{V}$. We further calculate the frequency-dependent SHG susceptibility of penta-ZnO₂, especially consider its imaginary part ($\text{Im}[\chi_{\text{sheet } abc}(\omega)]$). The real parts, imaginary parts, and modulus of different nonzero independent SHG tensor components are plotted in Fig. S14. One can see that the $\text{Im}[\chi_{\text{sheet } 14}(\omega)]$ and $\text{Im}[\chi_{\text{sheet } 36}(\omega)]$ of penta-ZnO₂ have the most significant peak values of 236.85 pm²/V and 639.56 pm²/V at 2.54 eV, respectively. The most significant peak value of $\text{Im}[\chi_{\text{sheet } abc}(\omega)]$ in penta-CdO₂ is at a lower energy range due to its smaller band gap of 3.28 eV at the HSE06 level compared with that of penta-ZnO₂, leading to a larger SHG response in penta-CdO₂.