

Supporting Information for

**Putting Xenon and Nitrogen under Pressure: Towards New Layered and  
Two-Dimensional Nitrogen Allotropes with Crown Ether-like Nanopores**

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# S1. Methodology details

## S1.1 Evolutionary algorithm (USPEX)

We provide here the technical details of the USPEX evolutionary structural searches with a brief description of the fixed-composition evolutionary algorithm (EA). USPEX<sup>1-4</sup> (Universal Structure Predictor: Evolutionary Xtallography, version 10.4) is a global evolutionary search algorithm for structure prediction developed by the A.R. Oganov laboratory since 2004 (see <http://uspex-team.org/en/uspex/overview>). In this work, this EA code is interfaced with VASP<sup>5,6</sup> (Vienna Ab initio Simulation Package, version 5.4.4) for DFT structure relaxation (shape, volume, and atomic positions are optimized by VASP). To identify stable ground-state structures and compositions within the binary Xe-N system and other Ng-N binary systems (Ng = He, Ne, Ar, and Kr), a variable-composition evolutionary algorithm (VC-EA) implemented in the USPEX code is employed. It allows us to scan the configurational, structural, and composition spaces of XeN<sub>x</sub> (where x is positive integer) and seek the local minima on the potential energy surface (PES). The VC-EA is performed at given pressures, in this case, 25, 50, 75 and 100 GPa. Finally, fixed composition EA (FC-EA) structural searches are undertaken for the given (x, y) compositions and pressures to ensure that the proposed structures are indeed the lowest in energy on the PES. The total number of atoms in the primitive cell is up to 44 for FC-EA searches, unless stated otherwise. One-hundred candidate structures are randomly produced in the first generation, and all subsequent generations contain 80 structures.

A structure is removed from the pool according to its fitness, that is the computed free enthalpy derived from *ab initio* total energy calculations (VASP). The remaining structures form the parent structures which participate in producing the next generation. Thus, a new candidate structure is produced from parent structures using one of four operators: (i) heredity (50%), (ii) permutation (10%), (iii) lattice mutations (10%), and (iv) softmutation (10%). Thus, 80% of the new candidates are generated from these evolutionary operators, while 20% were produced randomly. The USPEX search is done when

the global structure minimum is found in the last 10 generations. For a USPEX EA search, the number of generations is about 15-50. As mentioned previously, structure relaxations and energy calculations are done by the external code VASP (5 INCAR files, 5 steps per phase). Each fixed-composition USPEX job is run at least twice to ensure convergence to a global minimum.

## S1.2 DFT computational details

**DFT framework of USPEX calculations.** First-principles calculations are performed using the projected-augmented-wave (PAW) method<sup>7,8</sup> as implemented in VASP (version 5.4.4). In USPEX calculation, exchange-correlation energy is treated using Perdew-Burke-Ernzerhof (PBE)<sup>9,10</sup>, within the generalized gradient approximation (GGA) PAW potentials for xenon and nitrogen atoms have a radius of 2.0 au for Xe and 1.5 au for N ([He] core), respectively. Pseudopotentials with  $4d^{10}5s^25p^6$  and  $2s^22p^3$  valence configurations are employed for Xe and N, respectively.

To optimize each crystalline structure, 5 successive steps of increasing convergence accuracy are usually required (*i.e.*, 5 INCAR files). An illustration of this local optimization procedure is given in the PhD manuscript of B. Huang, “Computational materials discovery: prediction of carbon dioxide and nitrogen-based compounds under pressure using density functional theory and evolutionary algorithm”, supervised by Pr. G. Frapper, IC2MP, Poitiers University (France), Dec. 2017 (<http://theses.univ-poitiers.fr/notice/view/59262>). The parameters and criteria associated with VASP calculations then correspond to the last (5th) step, with the highest accuracy. A kinetic cutoff energy of 1000 eV is used for the wavefunction expansion with a Monkhorst-Pack k mesh grid with a spacing of  $2\pi \times 0.03 \text{ \AA}^{-1}$ . All structures are optimized until the net forces on atoms are below  $1 \text{ meV/\AA}$ , resulting in enthalpies that converged to better than 1 meV per atom (lower than a chemical accuracy of 1 kcal/mol, *i.e.*, 0.04 eV/atom), and all forces on atoms are converged to less than  $0.005 \text{ eV/\AA}$ . Then, the lowest ground-state (meta)stable phases obtained at the PBE level of theory are selected. For the two-dimensional systems,



we set a sufficiently large vacuum space to avoid interactions between neighboring images (vacuum  $\sim 15$  Å).

### **Post-USPEX treatment: optimization**

To fully optimize the selected phases, we tested other functionals. We finally chose to treat the weak dispersion forces of molecular structures (and others) through the inclusion of the revised Vydrov–van Voorhis nonlocal correlation (rVV10)<sup>11,12</sup> and used the strongly constrained and appropriately normed (SCAN) meta-GGA functional.<sup>13</sup> In this work, all the structural and energy values are discussed based on SCAN+rVV10 functional. For phonon calculations and *ab initio* molecular dynamics simulations, DFT-D3<sup>14</sup> correction was used instead of SCAN+rVV10, to take the vdW interactions into account with affordable computational resources.

### **Post-USPEX treatment: electronic structure, band gap and chemical bonding analysis**

To perform chemical bonding analysis, we carried out single-point calculations at the SCAN+rVV10 level (using the geometries obtained from VASP) to calculate the density of states (DOS). Moreover, DFT-D3 correction was used in all phonon calculations and *ab initio* molecular dynamics simulations. SCAN has proven to be accurate for the description of structural properties of materials. In contrast, it may underestimate the value of the band gap. Therefore, we calculate the band gap ( $E_g$ ) at the Heyd–Scuseria–Ernzerhof (HSE06)<sup>15</sup> hybrid functional level of theory, using the optimized SCAN+rVV10 structure (single-point energy calculation). This level of theory is noted thereafter as HSE06//SCAN+rVV10.

To further probe the electronic interaction between the subunits constituting the crystal structures (*e.g.*, Xe atoms and N<sub>20</sub> matrix in XeN<sub>20</sub>), we conducted additional single point calculations for each separated subunit, in the geometry, they adopt in the optimized structures, and computed the electron

density reorganization associated to the formation of the XeN<sub>x</sub> crystal structures from the separated subunits,

$$\Delta\rho(r) = \rho(r) - \sum_{\text{subunits}} \rho_{\text{subunit}}(r)$$

We then plotted the associated isodensity maps (cf section S7). Non-zero values indicate the electron density in the XeN<sub>x</sub> structure deviates from the simple superimposition of subunit densities, hence may probe charge transfer between subunits as well as intra-subunit polarization.  $\Delta\rho$  calculations and plots were managed using VESTA.<sup>16</sup>

To complement this analysis, we additionally computed the Non-Covalent Interactions Index (NCI)<sup>17,18</sup> using the critic2 code,<sup>19,20</sup> working with CHGCAR files from VASP. NCI analysis relies on the fact that non-covalent interactions are expected to occur in regions where both the electron density and its (reduced) gradient are low valued; hence a study of regions displaying low values of both quantities should help in identifying interaction patterns. Inclusion of the second-order derivative of the electron density further helps in discriminating the type of interactions: depending on the sign of the second-highest Hessian eigenvalue, the interactions are stabilizing (<0, e.g. hydrogen bonds), repulsive (>0, e.g. steric clash) or weakly attractive (~0, e.g. van der Waals). NCI visualization (reduced electron density gradient isosurfaces multiplied by the sign of the second Hessian eigenvalue) was made using VMD scripts<sup>21</sup> (critic2 output).

**Table S1:** The different levels of theory employed and the properties computed in this work (√ stands for done, and x for undone).

Functionals	USPEX searches	Structural optimization	Single point energy	Electronic properties	Dynamical stability (Phonon)	Kinetic stability (AIMD)
PBE <sup>9</sup>	√	√	×	×	×	×
DFT-D3 <sup>14</sup>	×	√	√	×	√	√
SCAN+rVV10 <sup>11-13</sup>	×	√	√	√	×	×
HSE06 <sup>15</sup> // SCAN+rVV10	×	×	√	√	×	×

### S1.3 AIMD simulations

*Ab initio* molecular dynamics (AIMD) within canonical *NVT* ensemble using the Nosé heat bath scheme were performed to evaluate the thermal stability of specific phases up to 2000 K for 12 ps with a time step of 1 fs, and we allowed 2 ps for thermalization and then extracted data from the last 10 ps. In such AIMD simulations, the Brillouin zone integration is restricted to the  $\Gamma$  point of the supercell, due to a high calculation cost.

### S1.4 Phonon calculations

In this work, first-principles phonon calculations using the finite displacements method at a quasi-harmonic level are done using the open-source package PHONOPY<sup>22</sup> (<https://atztogo.github.io/phonopy/>). Supercell structures with or without displacements are created from a reference unit cell considering all possible crystal symmetry operations. In general, a 2 x 2 x 2 supercell is sufficient, but larger ones can be required to avoid unphysical imaginary frequencies. Force constants are calculated using the optimized structure (VASP).

## S2. Structural parameters and energies

**Table S2.** Structural parameters of dynamically stable Noble-nitride phases (distances in Å, angles in °, energy in eV/atom) at the SCAN+rVV10 level of theory.

Phase	Pressure (GPa)	Space group	Z	Lattice parameters	Atomic coordinates (fractional)
XeN <sub>20</sub>	100	<i>Fm-3</i>	4	a=b=c=8.037 α=β=γ=90.0	Xe (0.000, 0.000, 0.500); N (0.000, 0.217, 0.170), (0.141, 0.141, 0.141)
XeN <sub>20</sub>	0	<i>Fm-3</i>	4	a=b=c=8.673 α=β=γ=90.0	Xe (0.000, 0.000, 0.500); N (0.000, 0.171, 0.211), (0.138, 0.138, 0.138)
XeN <sub>18</sub>	100	<i>P-1</i>	1	a=5.452, b=5.446, c=5.009 α=110.1, β=94.9, γ=60.6	Xe (0.000, 0.000, 0.000); N (0.321, 0.350, 0.258), (-0.342, -0.322, 0.239), (0.078, -0.458, 0.214), (-0.378, -0.068, 0.247), (-0.466, 0.388, 0.258), (-0.099, 0.487, 0.291), (0.357, 0.102, 0.272), (0.442, -0.354, 0.245), (-0.060, 0.101, 0.466)
XeN <sub>16</sub>	100	<i>P-1</i>	1	a=4.671, b=5.418, c=5.469 α=119.2, β=100.0, γ=104.6	Xe (0.000, 0.000, 0.000); N (0.222, 0.804, 0.411), (0.245, 0.519, 0.999), (0.266, 0.072, 0.640), (0.267, 0.609, 0.468), (0.277, 0.249, 0.553), (0.291, 0.795, 0.185), (0.339, 0.516, 0.788), (0.343, 0.736, 0.750)
XeN <sub>8</sub>	100	<i>P-1</i>	1	a=5.452, b=5.230, c=2.599 α=80.6, β=78.4, γ=114.7	Xe (0.000, 0.000, 0.000); N (0.049, 0.373, 0.156), (0.119, 0.093, -0.485), (0.414, -0.472, -0.098), (0.333, 0.667, 0.222)
XeN <sub>8</sub>	0	<i>P6/m</i>	1	a=b= 5.886, c=4.796 α=β=90.0, γ=120.0	Xe (0.000, 0.000, 0.000); N (0.083, -0.379, 0.500), (0.333, 0.667, 0.500)
XeN <sub>4</sub>	25	<i>I4<sub>1</sub>/amd</i>	4	a=b=5.330, c=8.930 α=β=γ=90.0	Xe (0.000, 0.250, 0.375); N (0.000, -0.025, -0.060)
XeN <sub>4</sub> <sup>23</sup>	18.7	<i>I4<sub>1</sub>/amd</i>	4	a=b= 5.723, c= 9.213 α=β=γ=90.0	Xe (0.000, 0.250, 0.375); N (0.000, -0.025, -0.060); N (0.5, 0.721, 0.719)
XeN <sub>4</sub> <sup>24</sup>	10.3	<i>I4<sub>1</sub>/amd</i>	4	a=b= 5.982, c= 9.338 α=β=γ=90.0	Xe (0.000, 0.250, 0.375); N (0.000, -0.025, -0.060) [0.238, 0.252, 0.809] [0.272, 0.229, 0.687]
XeN <sub>4</sub>	5	<i>I4<sub>1</sub>/amd</i>	4	a=b=6.241, c=10.057 α=β=γ=90.0	Xe (0.000, 0.250, 0.375); N (0.000, 0.023, -0.446)
XeN <sub>4</sub> <sup>23,24</sup> Xe(N <sub>2</sub> ) <sub>2</sub>	5.6	<i>Fd-3m</i> Cu <sub>2</sub> Mg-type	2	a=b=c=9.236 α=γ=β=90.0	Xe (0.000, 0.000, 0.000); N <sub>2</sub> (0.625, 0.125, 0.125)

XeN<sub>4</sub> was identified as an *Fd-3m* structure Xe(N<sub>2</sub>)<sub>2</sub> at ~ 5 GPa in previous experimental studies,<sup>23,24</sup> the authors treated the freely rotating N<sub>2</sub> molecules as units during the symmetry analysis. It's convenient and easy to understand the crystal structure, however, the XeN<sub>4</sub> is not a real *Fd-3m* structure.

**Table S3.** Calculated enthalpies of the predicted XeN<sub>x</sub> compounds at different pressure.

Structure	Space Group	Pressure (GPa)	Enthalpy (eV/atom)			Formation Enthalpy (eV/atom)		
			PBE	D3	SCAN+rVV10	PBE	D3	SCAN+rVV10
<b>Stable</b>								
XeN <sub>4</sub>	<i>I4<sub>1</sub>/amd</i>	5	-5.929	-6.035	-13.816	-0.004	-0.017	-0.016
XeN <sub>8</sub>	<i>P-1</i>	100	-0.782	-0.932	-5.622	-0.286	-0.283	-0.236
XeN <sub>18</sub>	<i>P-1</i>	100	-1.873	-2.007	-4.688	-0.267	-0.260	-0.216
XeN <sub>20</sub>	<i>Fm-3</i>	100	-1.949	-2.085	-4.598	-0.248	-0.244	-0.204
<b>Metastable</b>								
XeN <sub>16</sub>	<i>P-1</i>	100	-1.705	-1.859	-4.731	-0.216	-0.229	-0.173
<b>Quenchable</b>								
XeN <sub>20</sub>	<i>Fm-3</i>	0	-6.291	-6.408	-8.867	1.643	1.573	1.577
XeN <sub>16</sub>	<i>P-1</i>	0	-6.530	-6.916	-9.729	1.312	0.973	1.012
XeN <sub>8</sub>	<i>P6/m</i>	0	-6.589	-6.551	-11.163	0.823	0.910	0.960

**Table S4.** Structural parameters and relative enthalpy of several nitrogen allotropes @ 0 GPa (distances in Å, angles in °, energy in eV/atom) at the SCAN+rVV10 level of theory. N indicates the number of atoms in each unit cell.

Phase	Relative E	Space group	N	Lattice parameters	Atomic coordinates (fractional)
$\epsilon$ -N <sub>2</sub>	0.000	<i>R-3c</i>	48	a=b=8.595, c=12.104 $\alpha=\beta=90.0$ , $\gamma=120.0$	N (0.000, 0.000, 0.454), (0.058, 0.790, 0.721)
$\alpha$ -N <sub>2</sub>	-0.016	<i>Pa-3</i>	8	a=b=c=5.222 $\alpha=\beta=\gamma=90.0$	N (0.061, 0.061, 0.061)
cg-N <sup>25,26</sup>	1.368	<i>I2<sub>1</sub>3</i>	8	a=b=c=5.222 $\alpha=\beta=\gamma=90.0$	N (0.166, 0.166, 0.166)
N <sub>8</sub> <sup>27</sup>	0.902	<i>P1</i>	16	a=4.406, b=6.436, c=7.467 $\alpha=99.6$ , $\beta=99.9$ , $\gamma=91.1$	N (0.037, 0.619, 0.871), (0.057, 0.011, 0.150), (0.149, 0.474, 0.981), (0.228, 0.062, 0.728), (0.274, 0.010, 0.256), (0.410, 0.542, 0.072), (0.450, 0.072, 0.831), (0.513, 0.391, 0.177), (0.516, 0.045, 0.375), (0.564, 0.504, 0.670), (0.610, 0.856, 0.434), (0.769, 0.440, 0.277), (0.787, 0.546, 0.771), (0.847, 0.871, 0.553), (0.989, 0.076, 0.613), (0.994, 0.448, 0.379)
LP-N <sup>28</sup>	1.573	<i>Pba2</i>	16	a=b=4.618, c=5.257 $\alpha=\beta=\gamma=90.0$	N (0.005, 0.329, 0.236), (0.171, 0.995, 0.763), (0.239, 0.743, 0.631), (0.241, 0.757, 0.369)
<i>t</i> -N <sub>4</sub> <sup>29</sup>	1.532	<i>I4<sub>1</sub>/a</i>	16	a=b=5.369, c=3.818 $\alpha=\beta=\gamma=90.0$	N (0.129, 0.295, 0.057)
N <sub>10</sub> <sup>30</sup>	0.807	<i>P2<sub>1</sub></i>	10	a=5.015, b=5.934, c=8.268 $\alpha=\gamma=90.0$ , $\beta=103.3$	N (0.018, 0.387, 0.270), (0.077, 0.717, 0.654), (0.103, 0.929, 0.022), (0.121, 0.174, 0.470), (0.128, 0.986, 0.869), (0.272, 0.450, 0.345), (0.289, 0.046, 0.119), (0.328, 0.133, 0.870), (0.336, 0.317, 0.470), (0.427, 0.170, 0.026)
N <sub>10</sub> <sup>31</sup>	1.388	<i>P6<sub>3</sub>/m</i>	10	a=b=5.687, c=5.366 $\alpha=\beta=90.0$ , $\gamma=120.0$	N (0.086, 0.812, 0.523), (0.115, 0.704, 0.750), (0.333, 0.667, 0.750)
BP-N <sup>32,33</sup>	1.470	<i>Cmce</i>	8	a=2.220, b=7.610, c=3.263 $\alpha=\beta=\gamma=90.0$	N (0.000, -0.411, -0.377)
N <sub>20</sub> <sup>a</sup>	1.384	<i>Fm-3</i>	40	a=b=c=8.417 $\alpha=\beta=\gamma=90.0$	N (0.000, 0.221, 0.166), (0.140, 0.140, 0.140)
3D-N <sub>8</sub> <sup>a</sup>	1.031	<i>P-1</i>	8	a=4.093, b=4.961, c=4.208 $\alpha=79.3$ , $\beta=89.7$ , $\gamma=80.2$	N (0.174, -0.323, -0.326), (-0.026, -0.075, - 0.368), (-0.431, 0.374, -0.474), (0.132, 0.468, - 0.073)
2D-N <sub>8</sub> <sup>a</sup>	1.080	<i>P-3</i>	8	a=b=5.880, c=15.425 $\alpha=\beta=90.0$ , $\gamma=120$	N (0.333, 0.667, 0.500), (-0.463, -0.381, 0.500)

<sup>a</sup>This work.

**Table S5.** Enthalpies and structural parameters of Nitrogen and Xenon elements at selected pressures (distances in Å, angles in °, energy in eV/atom) at the SCAN+rVV10 level of theory.

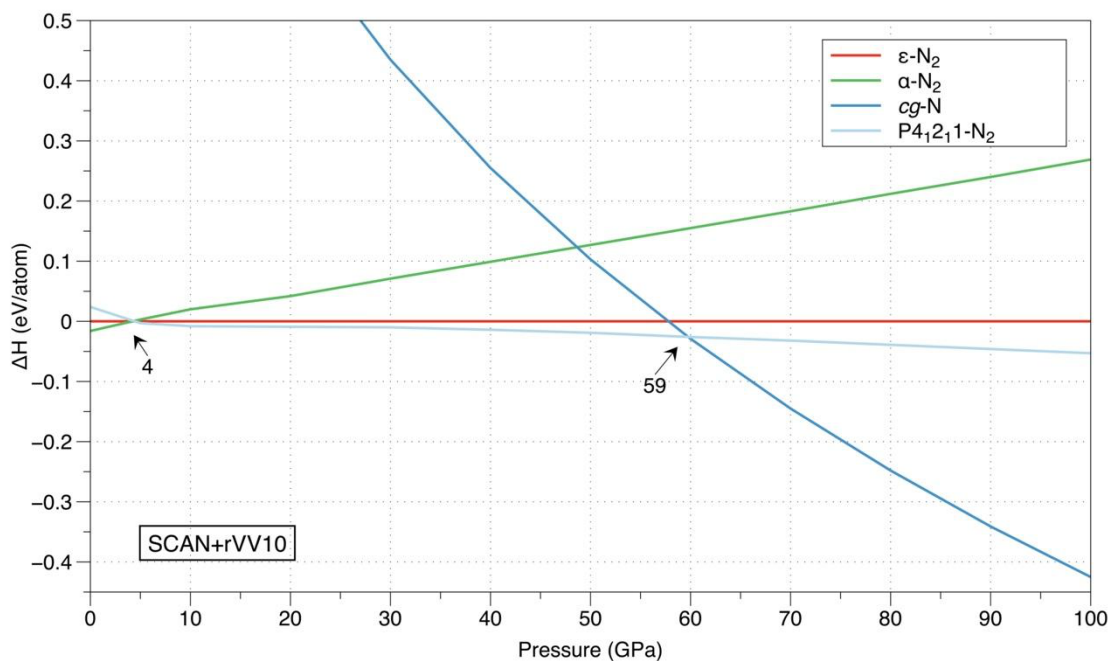
Phase	E	P	Space Group	Lattice Parameters	Atomic coordinates (fractional)
Xe	-35.617	0	<i>Fm-3m</i>	a=b=c=6.037 $\alpha=\beta=\gamma=90.0$	Xe (0.000, 0.000, 0.000)
Xe	-34.195	5	<i>Fm-3m</i>	a=b=c=5.426 $\alpha=\beta=\gamma=90.0$	Xe (0.000, 0.000, 0.000)
Xe	-30.082	25	<i>Fm-3m</i>	a=b=c=4.879 $\alpha=\beta=\gamma=90.0$	Xe (0.000, 0.000, 0.000)
Xe	-26.012	50	<i>Fm-3m</i>	a=b=c=4.608 $\alpha=\beta=\gamma=90.0$	Xe (0.000, 0.000, 0.000)
Xe	-22.480	75	<i>P6<sub>3</sub>/m</i>	a=b=3.135, c=5.136 $\alpha=\beta=90.0, \gamma=120.0$	Xe (0.333, 0.667, 0.250)
Xe	-19.273	100	<i>P6<sub>3</sub>/m</i>	a=b=3.047, c=4.999 $\alpha=\beta=90.0, \gamma=120.0$	Xe (0.333, 0.667, 0.250)
N	-9.187	0	<i>P<math>\bar{a}</math>-3<sup>a</sup></i>	a=5.856, b=5.856, c=5.856 $\alpha=90.0, \beta=90.0, \gamma=90.0$	N (0.055, 0.055, 0.055)
N	-8.685	5	<i>R-3c</i>	a=b=8.008, c=12.104 $\alpha=\beta=90.0, \gamma=120.0$	N (0.000, 0.000, 0.451), (0.061, 0.790, 0.721)
N	-7.311	25	<i>R-3c</i>	a=b=7.240, c=11.164 $\alpha=\beta=90.0, \gamma=120.0$	N (0.000, 0.000, 0.446), (0.063, 0.463, 0.882)
N	-5.935	50	<i>R-3c</i>	a=b=6.849, c=9.648 $\alpha=\beta=90.0, \gamma=120.0$	N (0.000, 0.000, 0.443), (0.064, 0.466, 0.878)
N	-4.740	75	<i>R-3c</i>	a=b=6.559, c=9.305 $\alpha=\beta=90.0, \gamma=120.0$	N (0.000, 0.000, 0.441), (0.065, 0.469, 0.876)
N	-3.658	100	<i>R-3c</i>	a=b=6.342, c=9.093 $\alpha=\beta=90.0, \gamma=120.0$	N (0.000, 0.000, 0.440), (0.068, 0.472, 0.874)
N	-7.323	25	<i>P4<sub>1</sub>2<sub>1</sub>2</i>	a=b=3.000, c=8.454 $\alpha=\beta=\gamma=90.0$	N (0.057, 0.814, 0.523)
N	-5.956	50	<i>P4<sub>1</sub>2<sub>1</sub>2</i>	a=b=2.834, c=7.949 $\alpha=\beta=\gamma=90.0$	N (0.064, 0.807, 0.524)
N	-4.940	75	<i>I2<sub>1</sub>3</i>	a=b=c=3.552 $\alpha=\beta=\gamma=90.0$	N (0.174, 0.174, 0.174)
N	-4.083	100	<i>I2<sub>1</sub>3<sup>b</sup></i>	a=b=c=3.501 $\alpha=\beta=\gamma=90.0$	N (0.176, 0.176, 0.176)

<sup>a</sup> Molecular crystal (N-N = 1.11 Å in triple bonded N<sub>2</sub>); <sup>b</sup>3D covalent net (N-N = 1.34 Å, 3-coordinated N, 10-rings).

**Table S6.** Selected parameters for bonding analysis in XeN<sub>x</sub> phases: interatomic separations, van der Waals radii, and electronegativities.

Pressure (GPa)	0	25	100
<b>Interatomic separation<sup>a</sup> (Å)</b>			
d(N≡N) triple bond <sup>b</sup>	1.11	-	-
d(N=N) double bond <sup>c</sup>	1.36	-	-
d(N-N) single bond <sup>d</sup>	1.41	1.38	1.34
d(Xe-Xe)	4.27 <sup>e</sup>	3.45 <sup>e</sup>	3.06 <sup>f</sup>
<b>van der Waals diameter<sup>g</sup> (Å)</b>			
r <sub>vdw</sub> (Xe-N)	4.11	3.02 <sup>h</sup>	2.70
r <sub>vdw</sub> (Xe-Xe)	4.64	3.40 <sup>h</sup>	3.04
<b>Electronegativity<sup>g</sup> (eV/e<sup>-1</sup>)</b>			
χ (N)	16.85	15.81 <sup>h</sup>	14.01
χ (Xe)	14.95	13.64 <sup>h</sup>	11.48
Δχ (Xe-N)	1.90	2.17 <sup>h</sup>	2.53

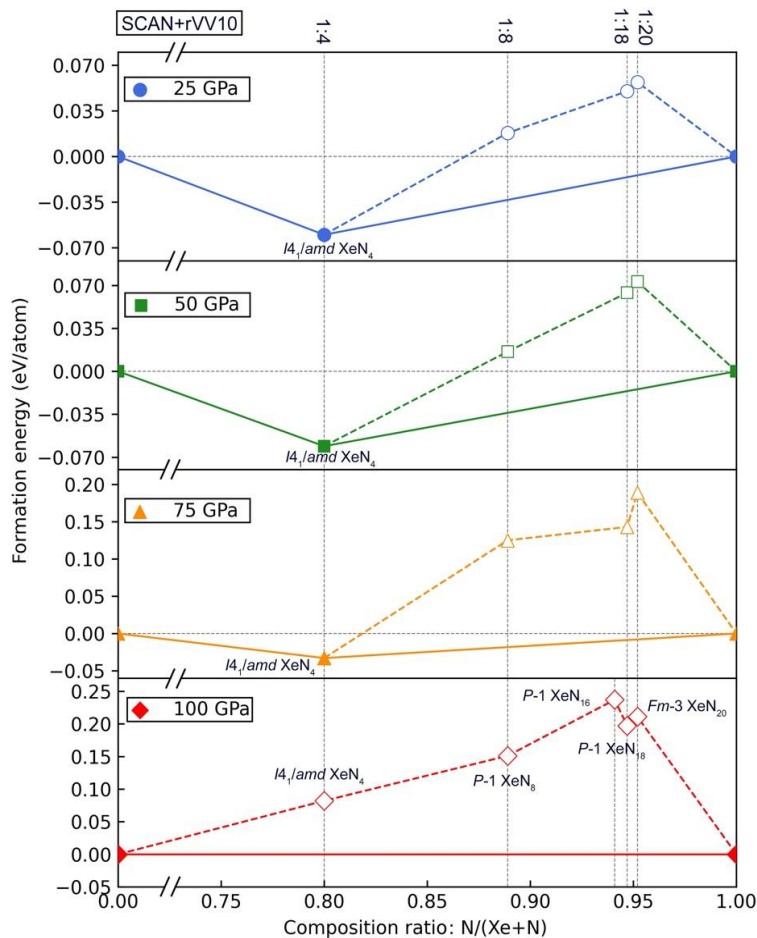
<sup>a</sup>SCAN+rVV10 level; <sup>b</sup>In molecular N<sub>2</sub>; <sup>c</sup>In molecular HN=NH; <sup>d</sup>In *cg*-N phase; <sup>e</sup>In *fcc* (*Fm-3m*) Xe phase; <sup>f</sup>In *hcp* (*P6<sub>3</sub>/m*) Xe phase; <sup>g</sup>The [Atoms Under Pressure \(AUP\) database](#)<sup>19–22</sup>; <sup>h</sup>at 24 GPa.



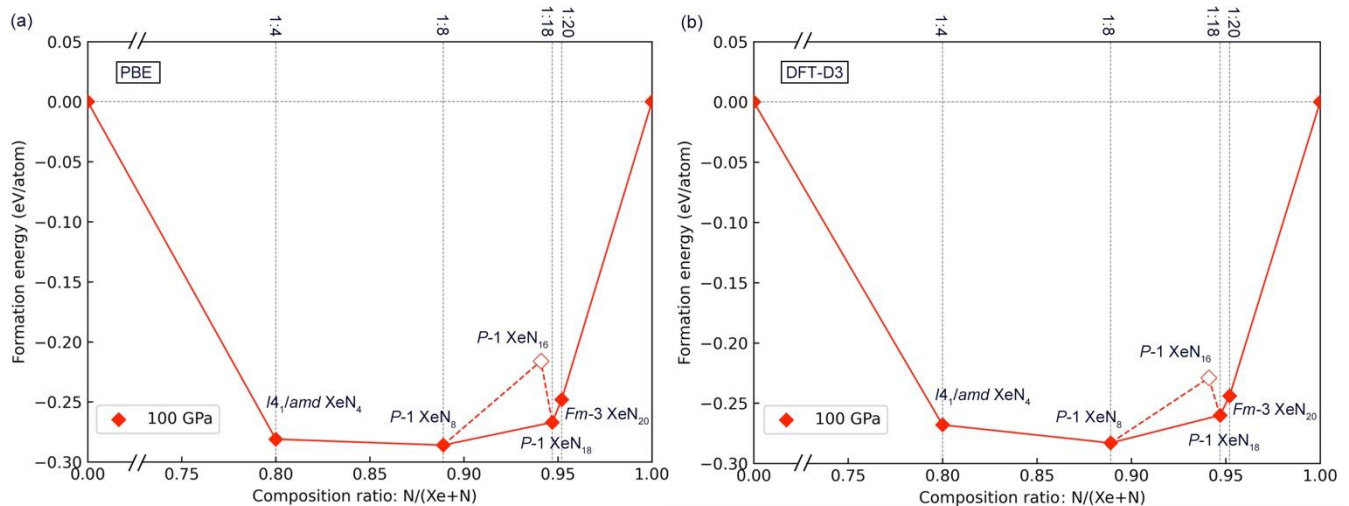
**Figure S1.** Relative enthalpies as a function of pressure for molecular nitrogen  $\alpha$ -N<sub>2</sub> (*Pa-3*),  $\epsilon$ -N<sub>2</sub> (*R-3c*), and *P4<sub>1</sub>2<sub>1</sub>2-N<sub>2</sub>*, as well as the polymeric *cg*-N.

Clearly,  $\alpha$ -N<sub>2</sub> is the ground stable phase at 0 GPa. When the pressure increases,  $\alpha$ -N<sub>2</sub> transfers to *P4<sub>1</sub>2<sub>1</sub>2-N<sub>2</sub>* above ~4 GPa. In the pressure range from 4 to 59 GPa, the  $\epsilon$ -N<sub>2</sub> is slightly above *P4<sub>1</sub>2<sub>1</sub>2-N<sub>2</sub>*. Finally, the polymeric *cg*-N turns out to be the most stable phase above 59 GPa.



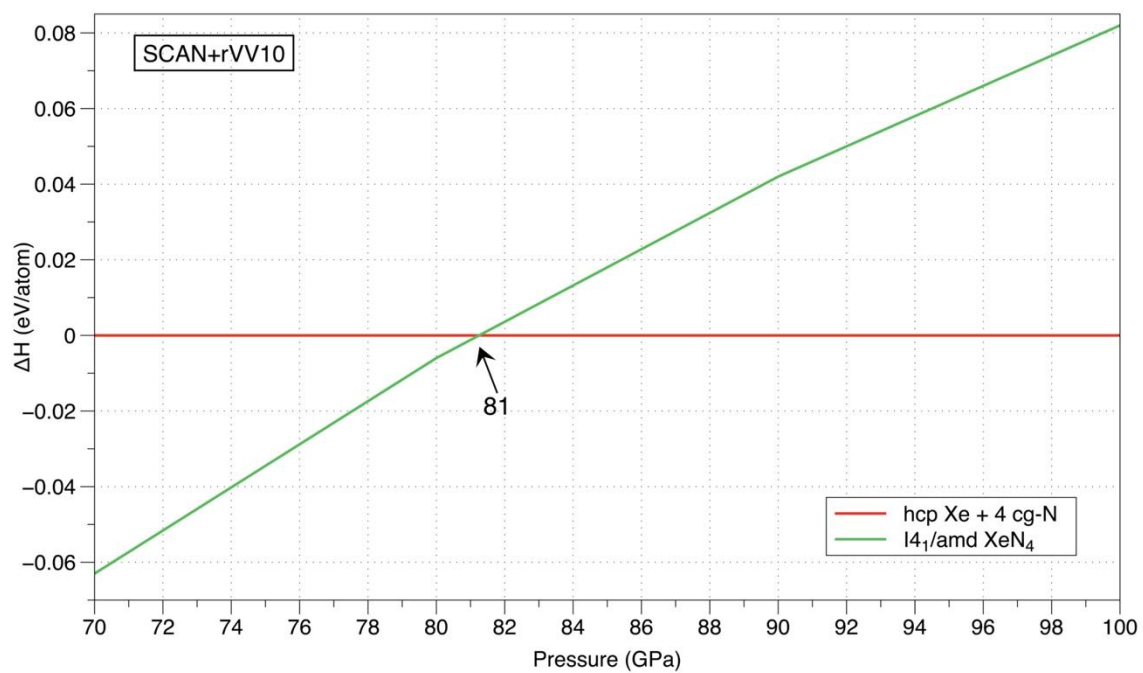


**Figure S2.** Convex-hull diagrams for the Xe-N system under selected pressure regarding the ground-state elements.  $P4_12_12-N_2$  and  $fcc$  Xe at 25 and 50 GPa, while  $cg-N$  and  $hcp$  Xe at 75 and 100 GPa were used to calculate the formation enthalpies. Solid symbols denote stable structures, and empty ones represent metastable structures. The space group is indicated for each novel stable structure.



**Figure S3.** Convex-hull diagrams for the Xe-N system under 100 GPa regarding the  $hcp$ -Xe ( $P6_3/m$ ) and  $\epsilon$ - $N_2$  ( $R-3c$ ) with (a) PBE and (b) DFT-D3 levels of theory. Solid symbols denote stable structures, and empty ones represent metastable structures. The space group is indicated for each novel stable structure.

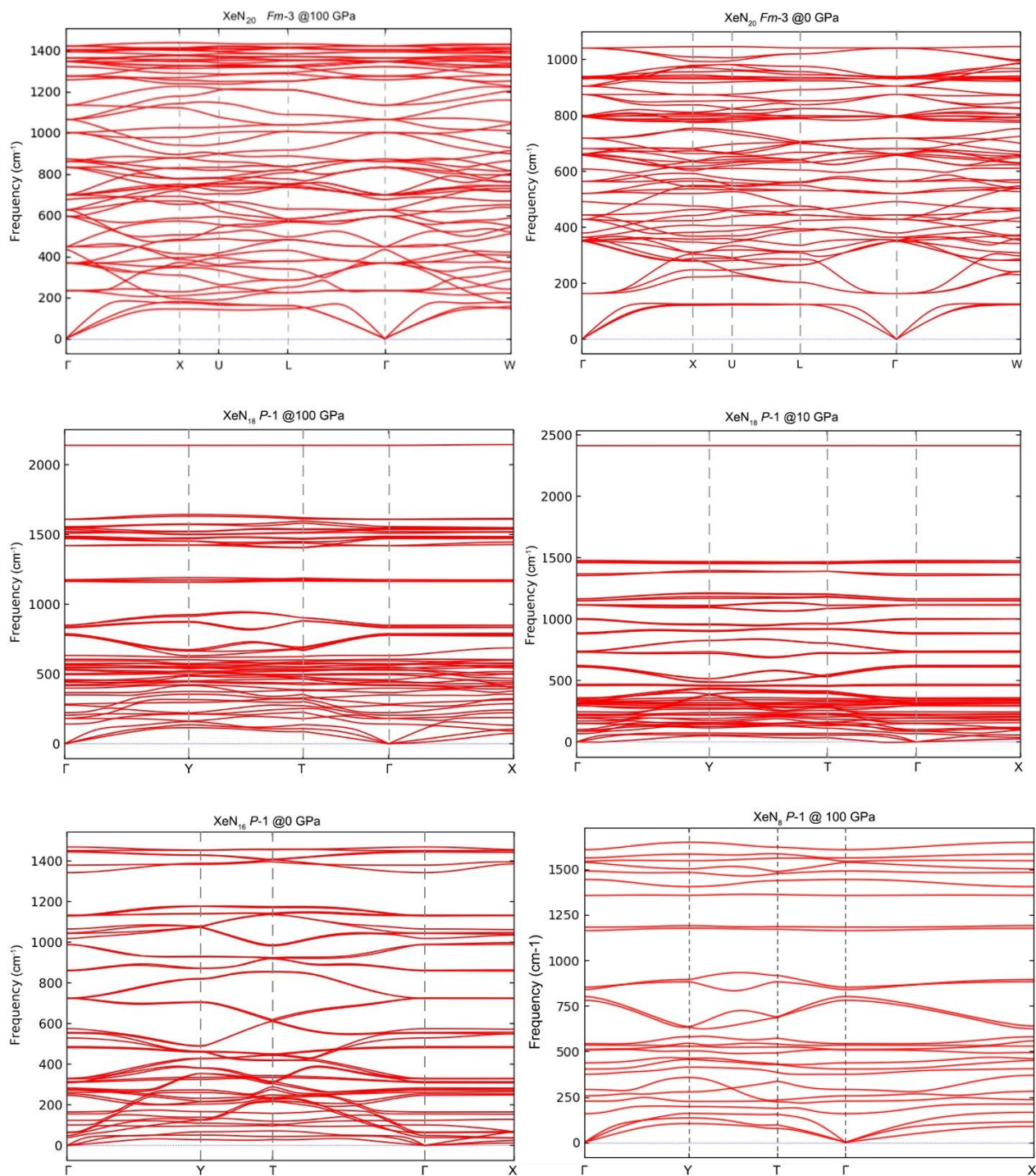
As shown in Figure S3, the PBE and DFT-D3 moderately shift the stability figures, but the phase stability order is not changed.

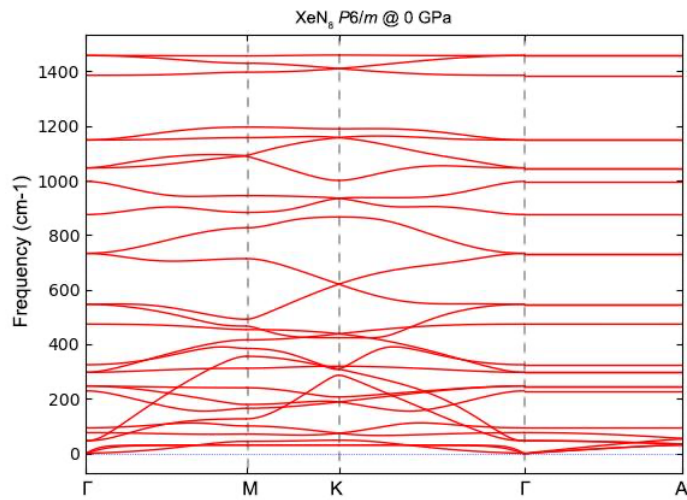


**Figure S4.** Calculated enthalpy curves of  $I_{41}/amd$  XeN<sub>4</sub> with respect to  $hcp$  Xe and  $cg$ -N under pressure (SCAN+rVV10 level).  $I_{41}/amd$  XeN<sub>4</sub> has a positive formation enthalpy above 81 GPa.

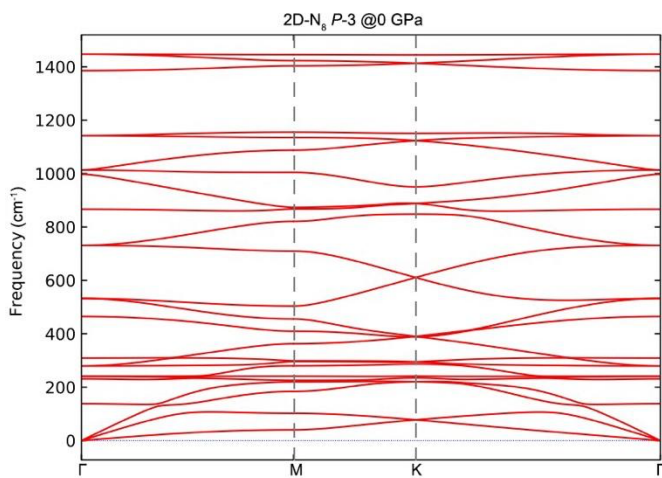
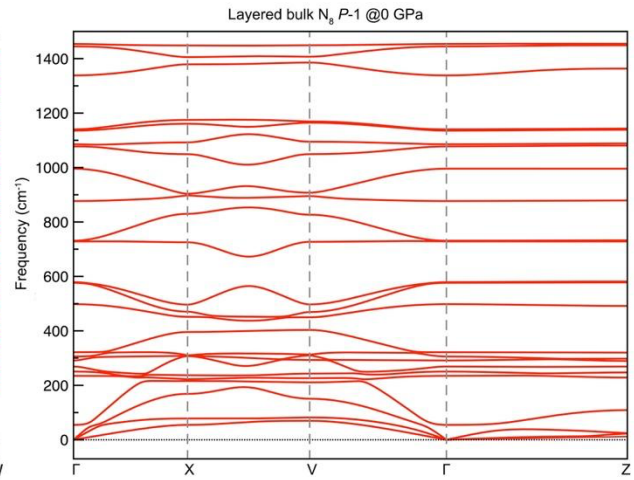
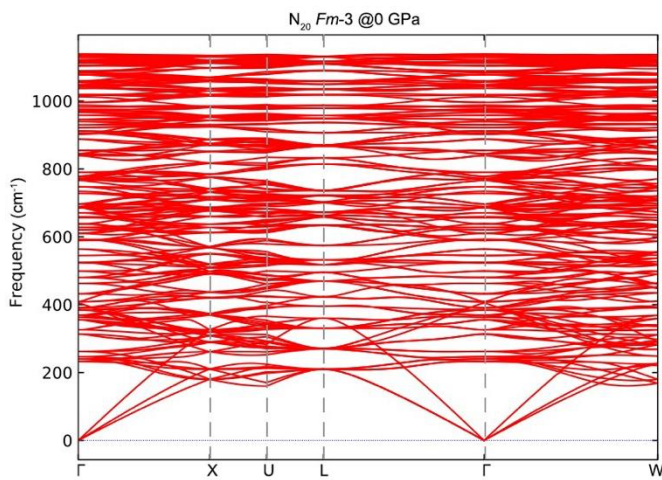
### S3. Dynamical stability

#### a. Xenon nitrides





b. Polymeric nitrogen compounds ( $N_{20}$  and layered 3D and 2D  $N_8$ )



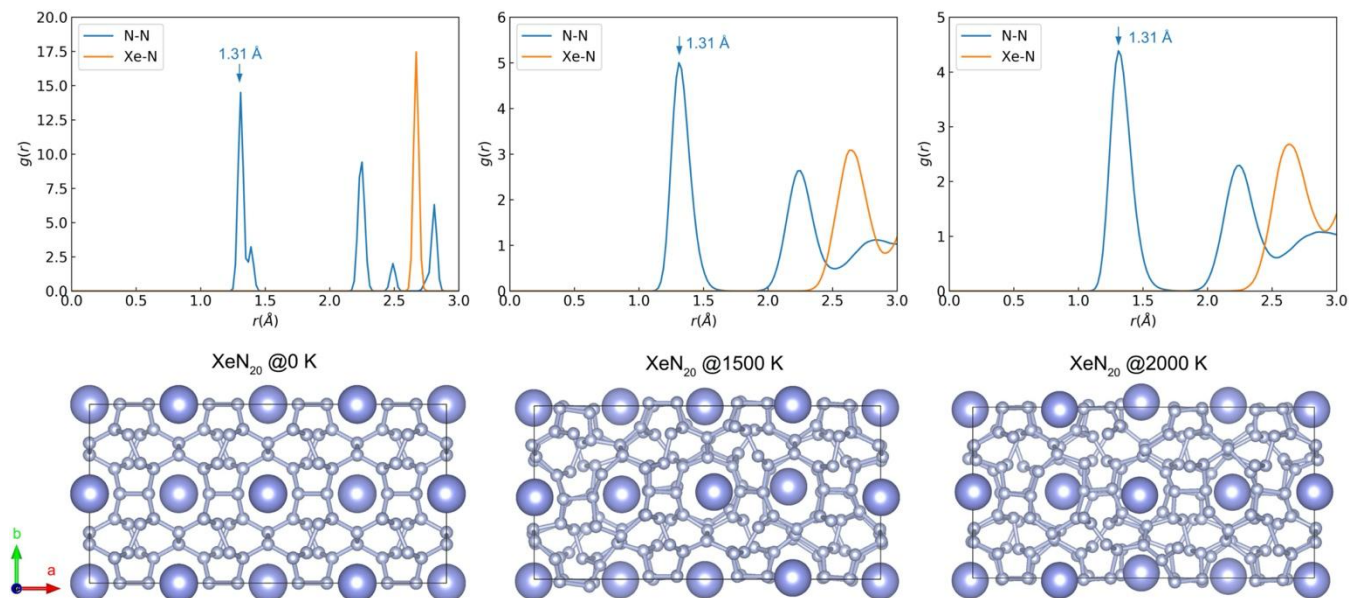
**Figure S5.** Phonon dispersion curves of  $XeN_{20}$ ,  $XeN_{18}$ ,  $XeN_8$  and polymeric nitrogen compounds ( $N_{20}$  and  $N_8$ ) at selected pressures (DFT-D3).

No imaginary phonon modes are found in the whole Brillouin zone. This confirms the dynamical stability of each phase.



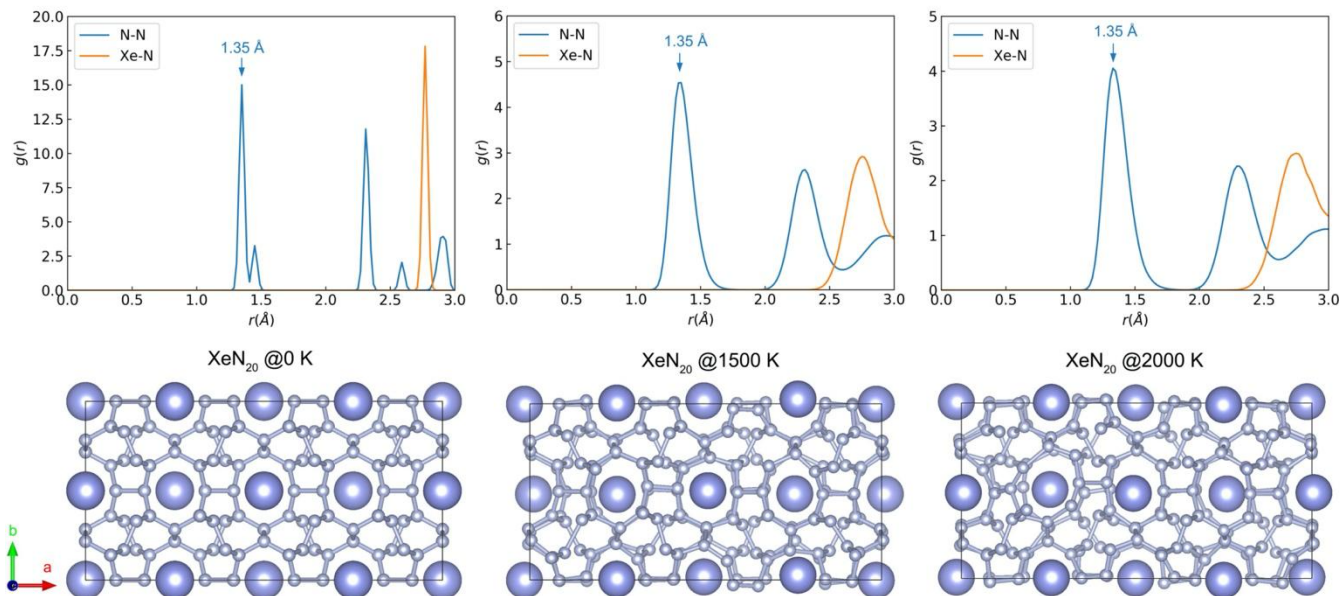
## S4. AIMD simulations

AIMD simulations with canonical ensemble using the Nosé heat bath scheme were performed to evaluate the thermal stability of specific phases up to 2000 K for 12 ps, see Figure S6-S12. Radial distribution functions (RDF) of specific phases are also presented, which reveal if the contacts between atoms (N-N, or Xe-N) are preserved at high temperatures.



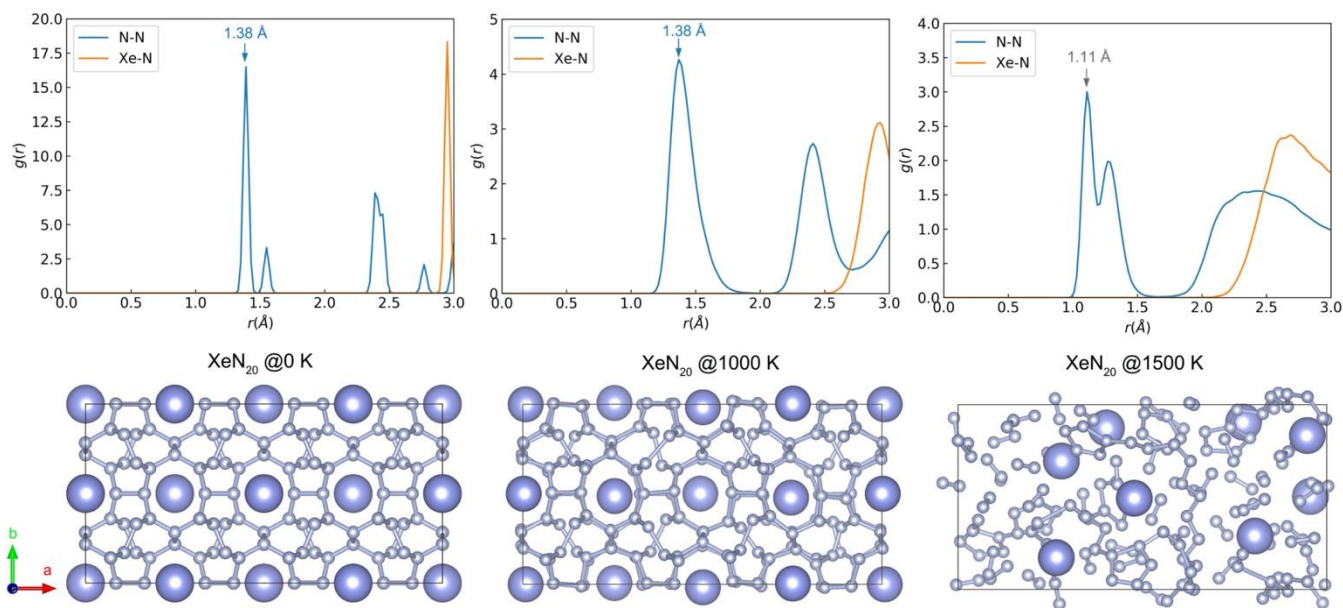
**Figure S6.** RDF (upper panel) and snapshots (lower panel) of the *Fm-3*  $\text{XeN}_{20}$  at 100 GPa with  $2 \times 1 \times 1$  supercell (168 atoms) after 12 ps AIMD simulations at 0 K, 1500 K, and 2000 K, respectively. Large portage spheres are xenon atoms, and small gray spheres are nitrogen atoms. The single bonds N-N under 100 GPa (1.31 Å) are indicated by blue arrows in RDF.

These AIMD simulations demonstrate the single bonds N-N framework within  $\text{XeN}_{20}$  could be maintained up to 2000 K at 100 GPa.



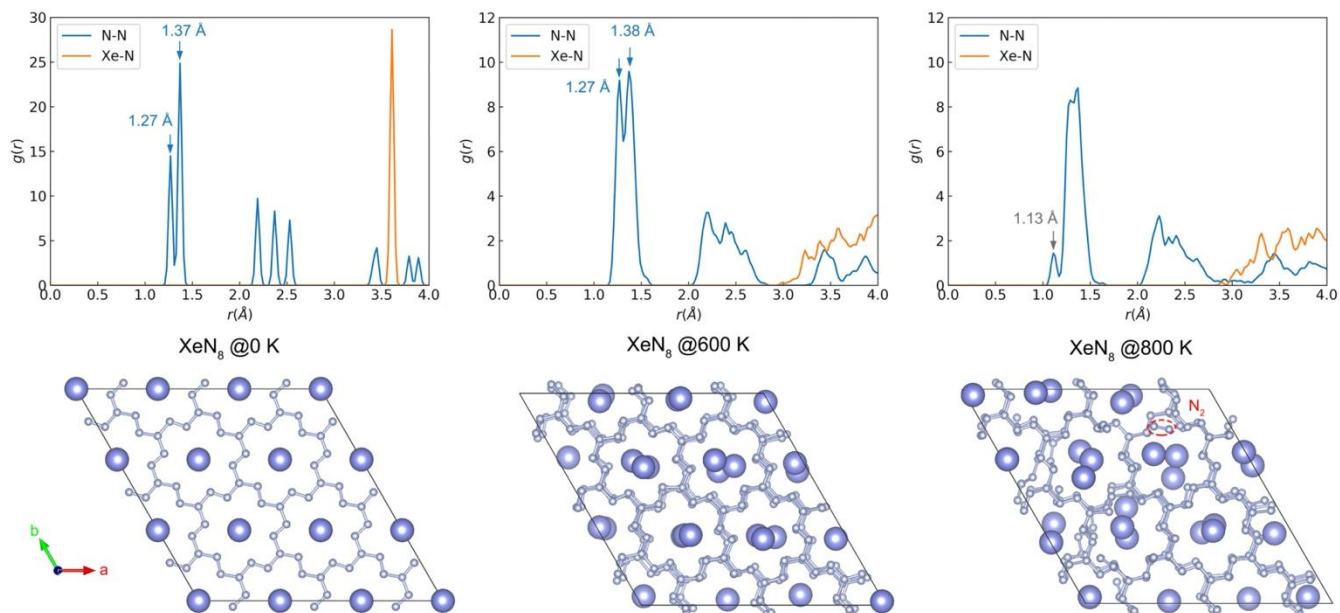
**Figure S7.** RDF (upper panel) and snapshots (lower panel) of the *Fm-3* XeN<sub>20</sub> at 50 GPa with 2 × 1 × 1 supercell (168 atoms) after 12 ps AIMD simulations at 0 K, 1500 K and 2000 K, respectively. The single bonds N-N under 50 GPa (1.35 Å) are shown by blue arrows in RDF.

These AIMD simulations analyses indicate that the single bonds N-N framework within XeN<sub>20</sub> could be maintained up to 2000 K at 50 GPa.



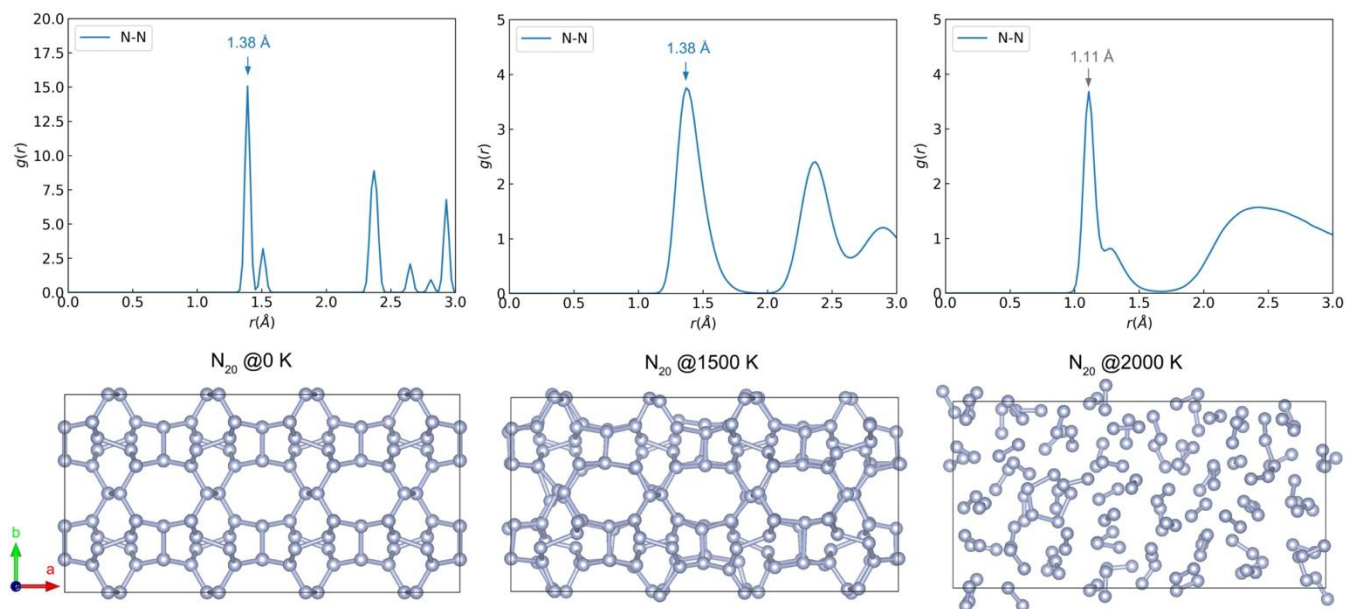
**Figure S8.** RDF (upper panel) and snapshots (lower panel) of the *Fm-3* XeN<sub>20</sub> at 0 GPa with 2 × 1 × 1 supercell (168 atoms) after 12 ps AIMD simulations at 0 K, 1000 K, and 1500 K, respectively. The single bonds N-N under 0 GPa (1.38 Å) are shown by blue arrows in RDF for 0 K and 1000 K.

The polymeric Nitrogen is crashed at 1500 K, and the signals for N<sub>2</sub> dimers (1.11 Å) are observed from both RDF and snapshots, which indicate the single bonds N-N framework within XeN<sub>20</sub> could be maintained up to 1000 K, but will be crashed under 1500 K at 0 GPa.



**Figure S9.** RDF (upper panel) and snapshots (lower panel) of the  $P6/m$   $\text{XeN}_8$  at 0 GPa with  $3 \times 3 \times 3$  supercells (243 atoms) after 12 ps AIMD simulations at 0 K, 600 K, and 800 K, respectively. The single bonds (1.27 Å) and double bonds (1.37 Å) N-N under 0 GPa are shown by blue arrows in RDF for 0 K and 600 K.

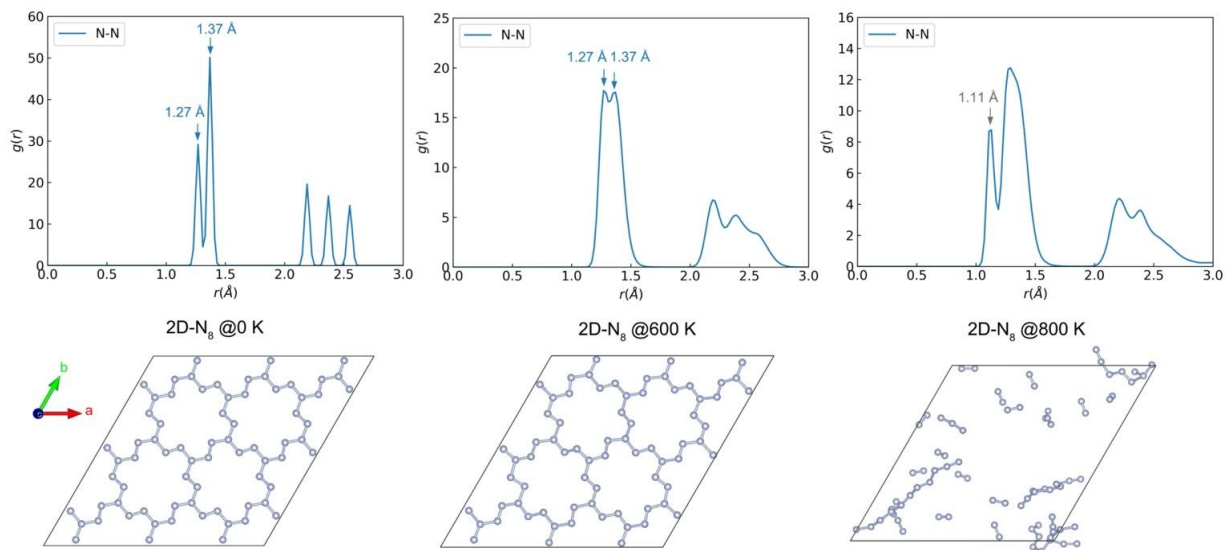
The polymeric Nitrogen is crashed at 800 K, and the signals for  $\text{N}_2$  dimers (1.13 Å) are observed from both RDF and snapshots, which indicate the single bonds N-N framework within  $P6/m$   $\text{XeN}_8$  could be maintained up to 600 K, but will be crashed under 800 K at 0 GPa.



**Figure S10.** RDF (upper panel) and snapshots (lower panel) of the  $Fm-3$   $\text{N}_{20}$  at 0 GPa with  $2 \times 1 \times 1$  supercell (160 atoms) after 12 ps AIMD simulations at 0 K, 1500 K, and 2000 K, respectively. The single bonds N-N under 0 GPa (1.38 Å) are shown by blue arrows in RDF for 0 K and 1500 K.

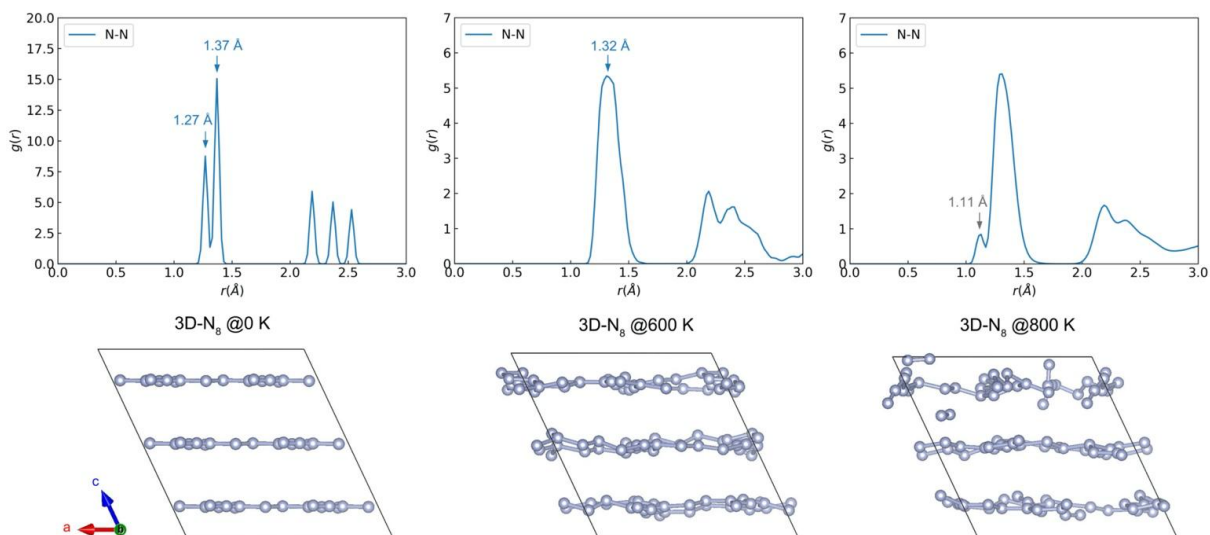


The polymeric Nitrogen is crashed at 1500 K, and the signals for N<sub>2</sub> dimers (1.11 Å) are observed from both RDF and snapshots, which indicate the single bonds N-N framework within N<sub>20</sub> could be maintained up to 1500 K, but will be crashed under 2000 K at 0GPa.



**Figure S11.** RDF (upper panel) and snapshots (lower panel) of the *P*-3 2D-N<sub>8</sub> at 0 GPa with 3 × 3 × 1 supercell (72 atoms) after 10 ps AIMD simulations at 0 K, 600 K, and 800 K, respectively. The single bonds (1.27 Å) and double bonds (1.37 Å) N-N under 0 GPa are shown by blue arrows in RDF for 0 K and 600 K.

The polymeric Nitrogen is crashed at 800 K, and the signals for N<sub>2</sub> dimers (1.11 Å) are observed from both RDF and snapshots, which indicate the single bonds N-N framework within 2D-N<sub>8</sub> could be maintained up to 600 K, but will be crashed under 800 K at 0 GPa.



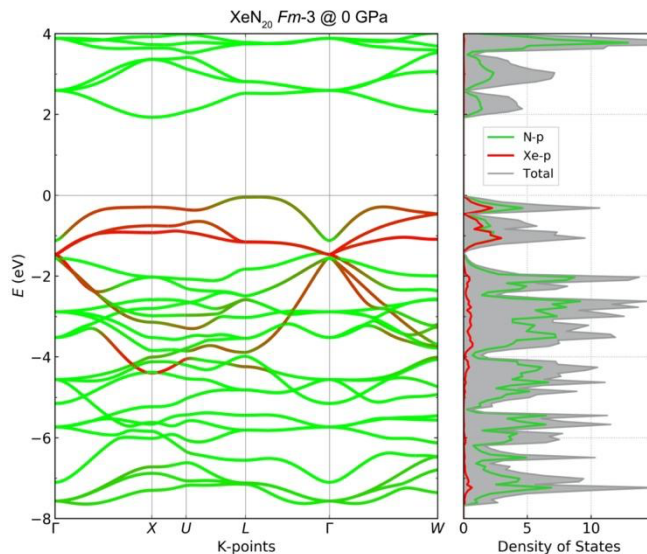
**Figure S12.** RDF (upper panel) and snapshots (lower panel) of the *P*-1 3D-N<sub>8</sub> at 0 GPa with 2 × 2 × 3 supercell (96 atoms) after 10 ps AIMD simulations at 0 K, 600 K, and 800 K, respectively. The single bonds (1.27 Å) and double bonds (1.37 Å) N-N under 0 GPa are shown by blue arrows in RDF for 0 K and 600 K.



The polymeric Nitrogen is crashed at 800 K, and the signals for N<sub>2</sub> dimers (1.11 Å) are observed from both RDF and snapshots, which indicate the single bonds N-N framework within 3D-N<sub>8</sub> could be maintained up to 600 K, but will be crashed under 800K at 0GPa.

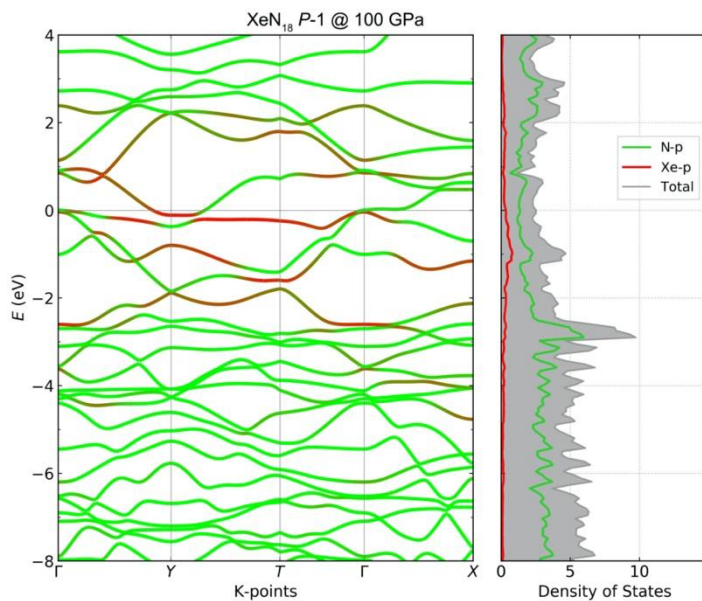
## S5. Electronic properties

- a. Projected band structures and density of states (DOS) at SCAN+rVV10 level. The colored bands of band structures indicate the atomic contributions to the band structure, from red (high Xe contributions) to green (high N contribution).



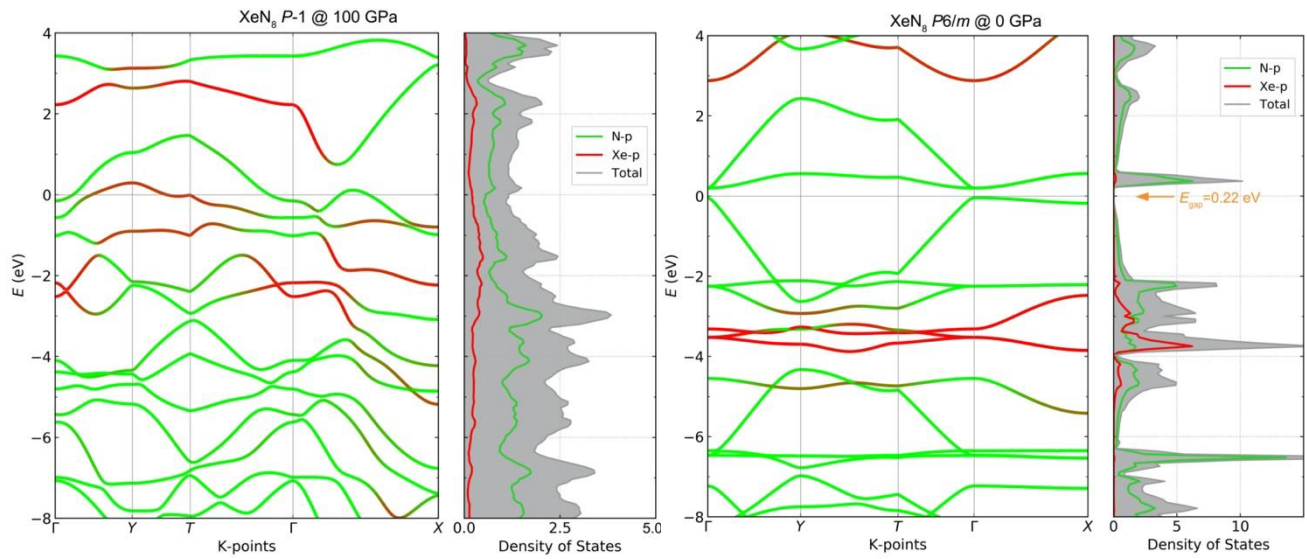
**Figure S13.** Band structure and DOS for  $Fm-3$   $XeN_{20}$  at 0 GPa.

A bandgap is observed for  $Fm-3$   $XeN_{20}$  at 100 and 0 GPa using SCAN+U+rVV10 functional.

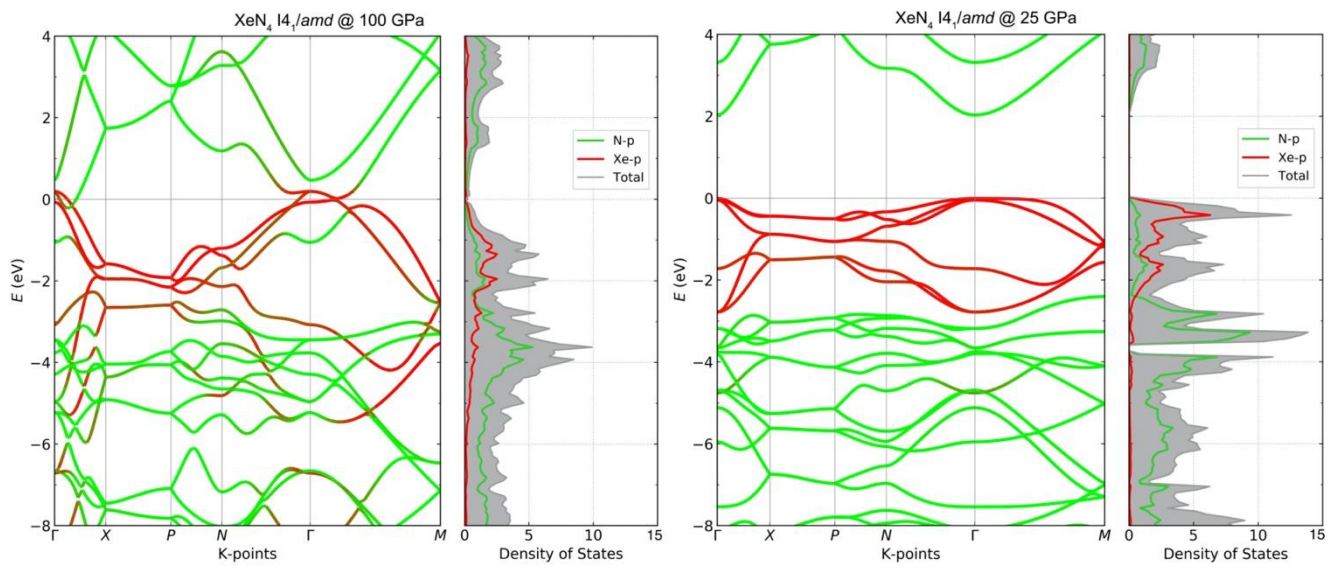


**Figure S14.** Band structure and DOS for  $P-1$   $XeN_{18}$  at 100 GPa.

Fermi level crosses the conductive bands, indicating that  $P-1$   $XeN_{18}$  is metallic at 100 GPa using SCAN+rVV10 functional.

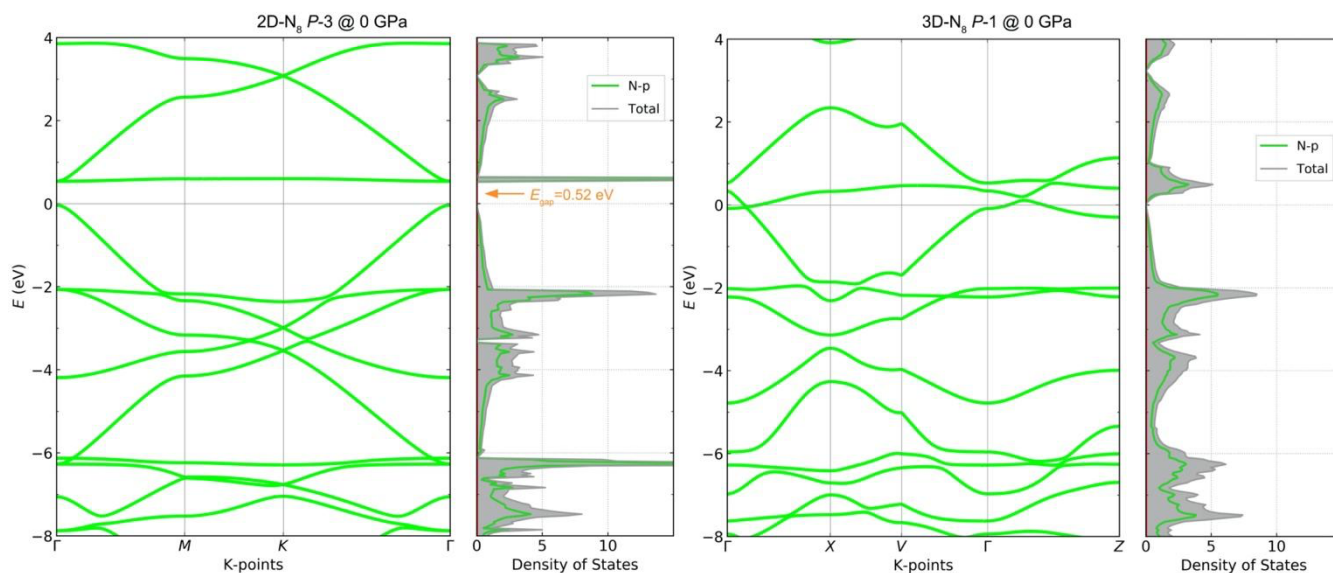


**Figure S15.** Band structure and DOS for  $P-1$   $\text{XeN}_8$  at 100 GPa, and  $P6/m$   $\text{XeN}_8$  0 GPa.



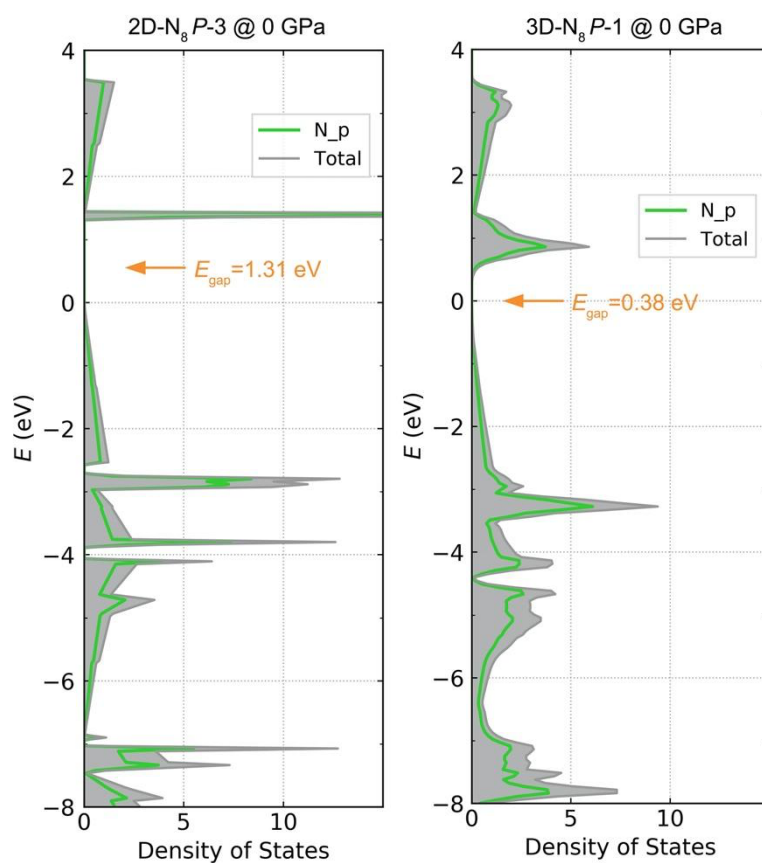
**Figure S16.** Band structure and DOS for  $I4_1/amd$   $\text{XeN}_4$  at 100 and 25 GPa.

The calculated metallic behavior of  $\text{XeN}_4$  at 100 GPa has good agreement with the previous experimental report.<sup>23,24</sup>



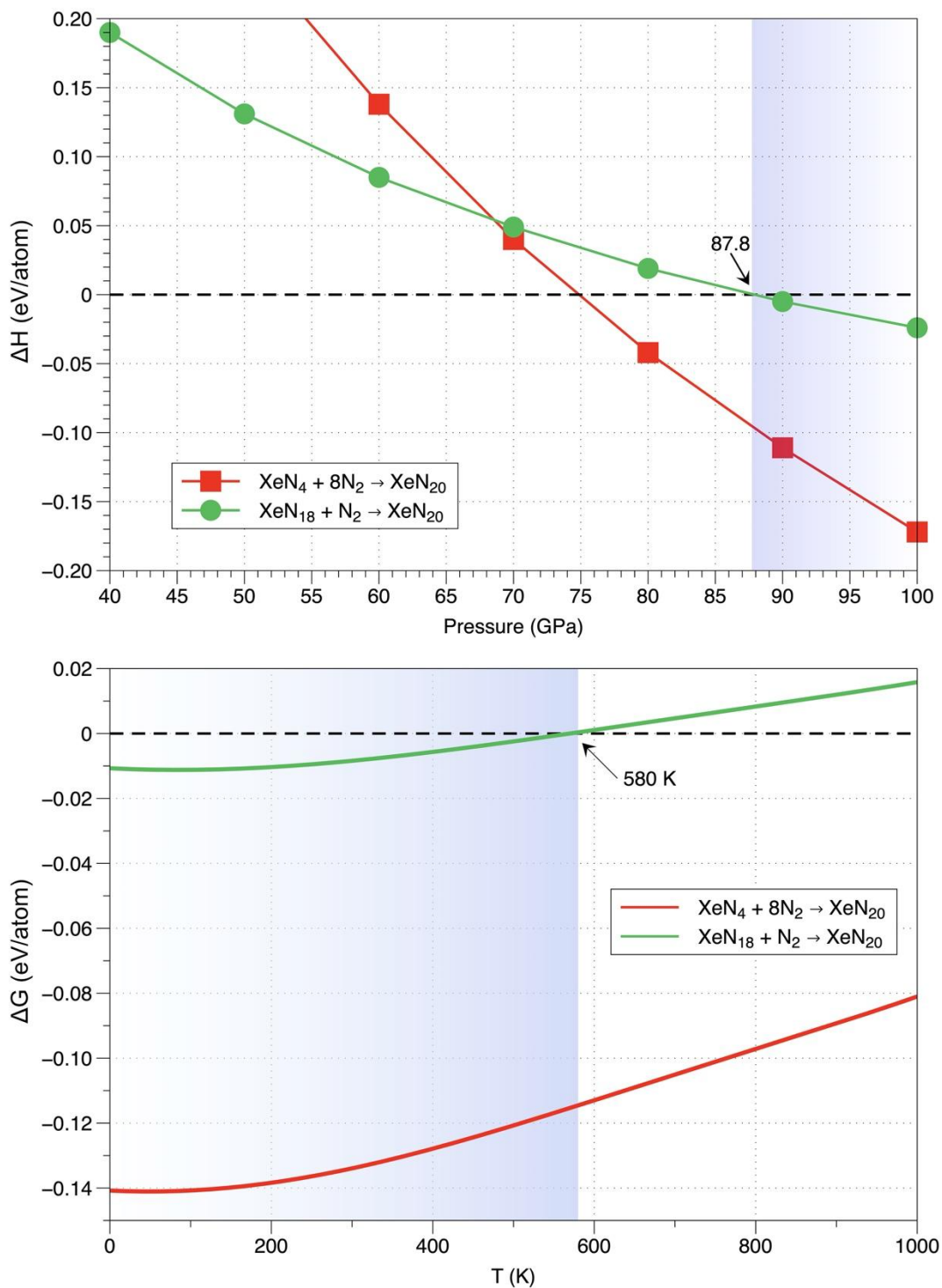
**Figure S17.** Band structure and DOS for  $P$ -3 2D- $N_8$  and  $P$ -1 3D- $N_8$  at 0 GPa.

b. Projected DOS-HSE06 at SCAN+rVV10 level.



**Figure S18.** DOS computed with HSE06//SCAN+rVV10 showing band gap for both  $P$ -3 2D- $N_8$  ( $E_{\text{gap}} = 1.31$  eV) and  $P$ -1 3D- $N_8$  ( $E_{\text{gap}} = 0.38$  eV) at 0 GPa.

## S6. Gibbs free energy of formation of XeN<sub>20</sub>



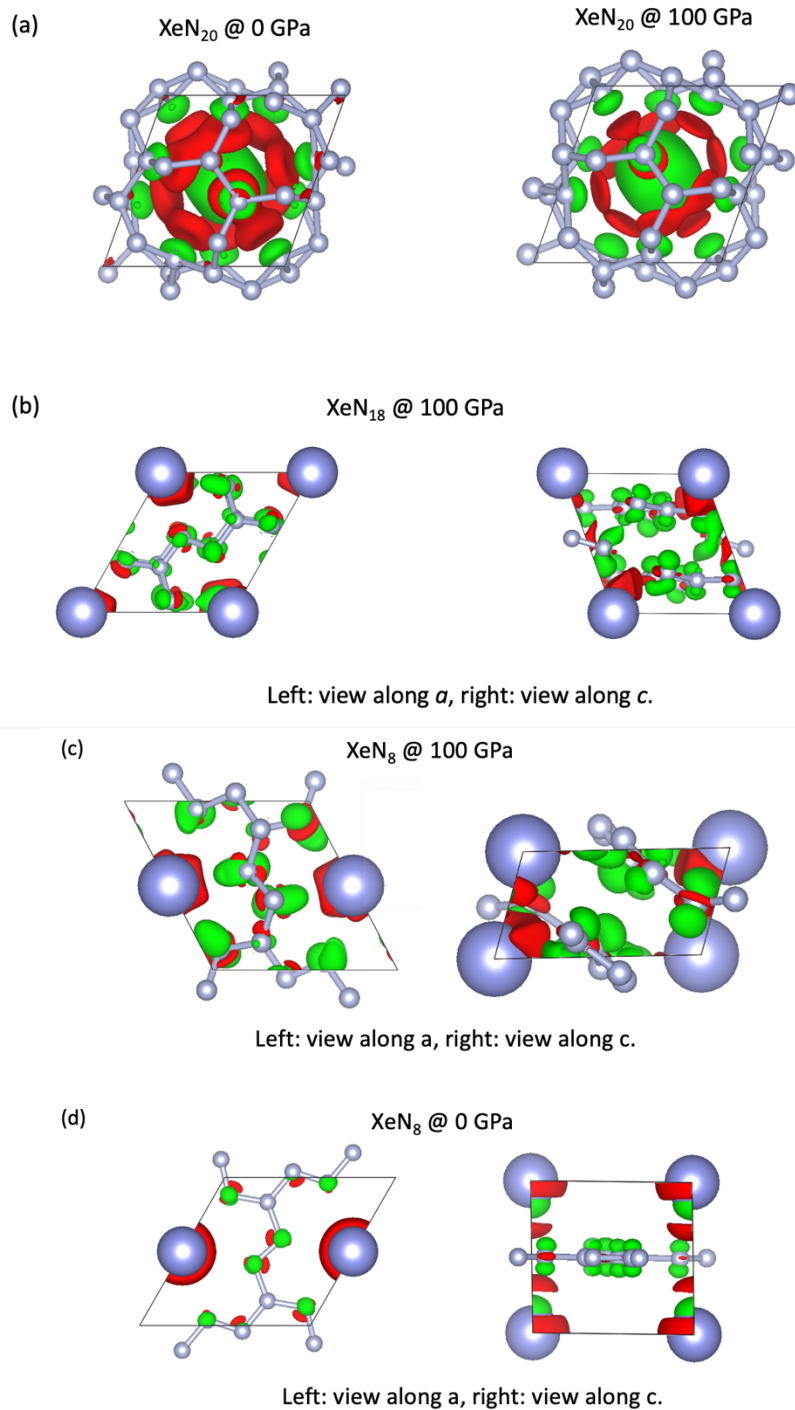
**Figure S19.** Relative enthalpies and Gibbs free energy of proposed reactions for  $\text{XeN}_4 + 8\text{N}_2 \rightarrow \text{XeN}_{20}$  (red line) and  $\text{XeN}_{18} + \text{N}_2 \rightarrow \text{XeN}_{20}$  (green line) with PBE functionals at high-pressure and high-temperature.

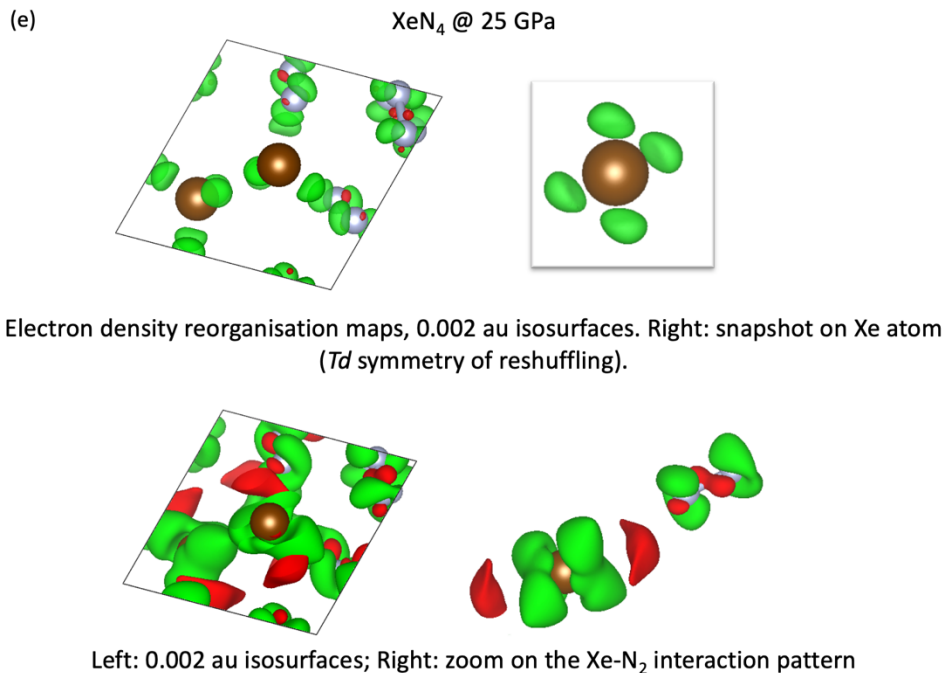
The calculations within DFT-D3 indicate that the XeN<sub>20</sub> is stable at high pressure and high temperatures.



## S7. Intermolecular interaction patterns in XeN<sub>x</sub> (x = 4, 8, 18, and 20)

### S7.1 Electron density reorganisation maps





**Figure S20.** Electron density reorganization map with isosurface for  $\text{XeN}_4$ ,  $\text{XeN}_8$ ,  $\text{XeN}_{18}$ , and  $\text{XeN}_{20}$  under specific pressure. Red and green regions indicate depletion and accumulation of electrons, respectively.

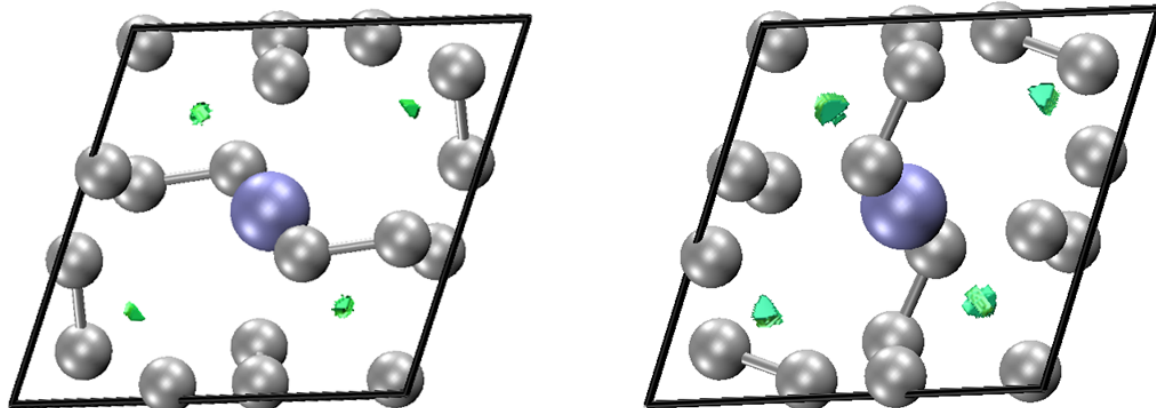
As shown in Figure S20, in most cases both Xe and polynitrogen subunits display both negative and positive contributions of  $\Delta\rho$ , which suggests each fragment is experiencing an electron density polarisation due to the interaction, rather than a genuine charge transfer. Nevertheless, the cases of  $\text{XeN}_{18}$  and  $\text{XeN}_8$  at 100 GPa could indicate a slight charge transfer from the Xe atom to the N matrix (only depletion areas being visible on Xe atoms in these cases). It may be noted in the case of  $\text{XeN}_8$  that this charge transfer is suppressed at 0 GPa, in line with the slight change in the 3D structure of the polynitrogen motif (fully planar at 0 GPa, corrugated at 100 GPa) and with the variations in the DOS (see S5). In both cases, the possibility to remove the Xe atoms without destruction of the polynitrogen moieties is a further indication of the limited extent of the charge transfer - which is also hinted by the NCI analysis (vide infra).

## S7.2 Non-Covalent Interaction Index analysis

(a)

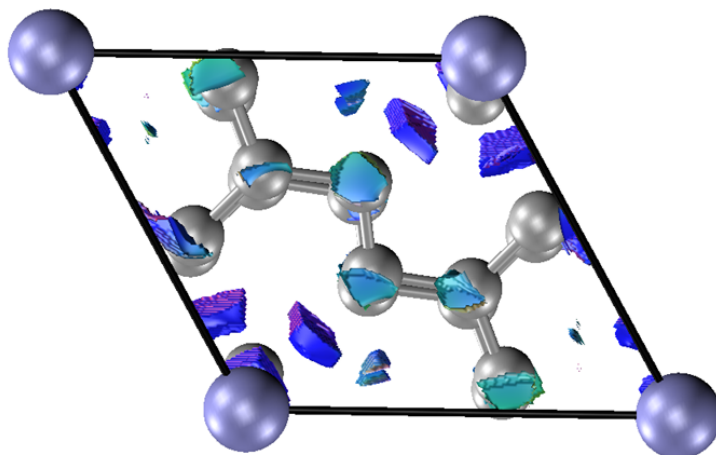
XeN<sub>20</sub> @ 0 GPa

XeN<sub>20</sub> @ 100 GPa

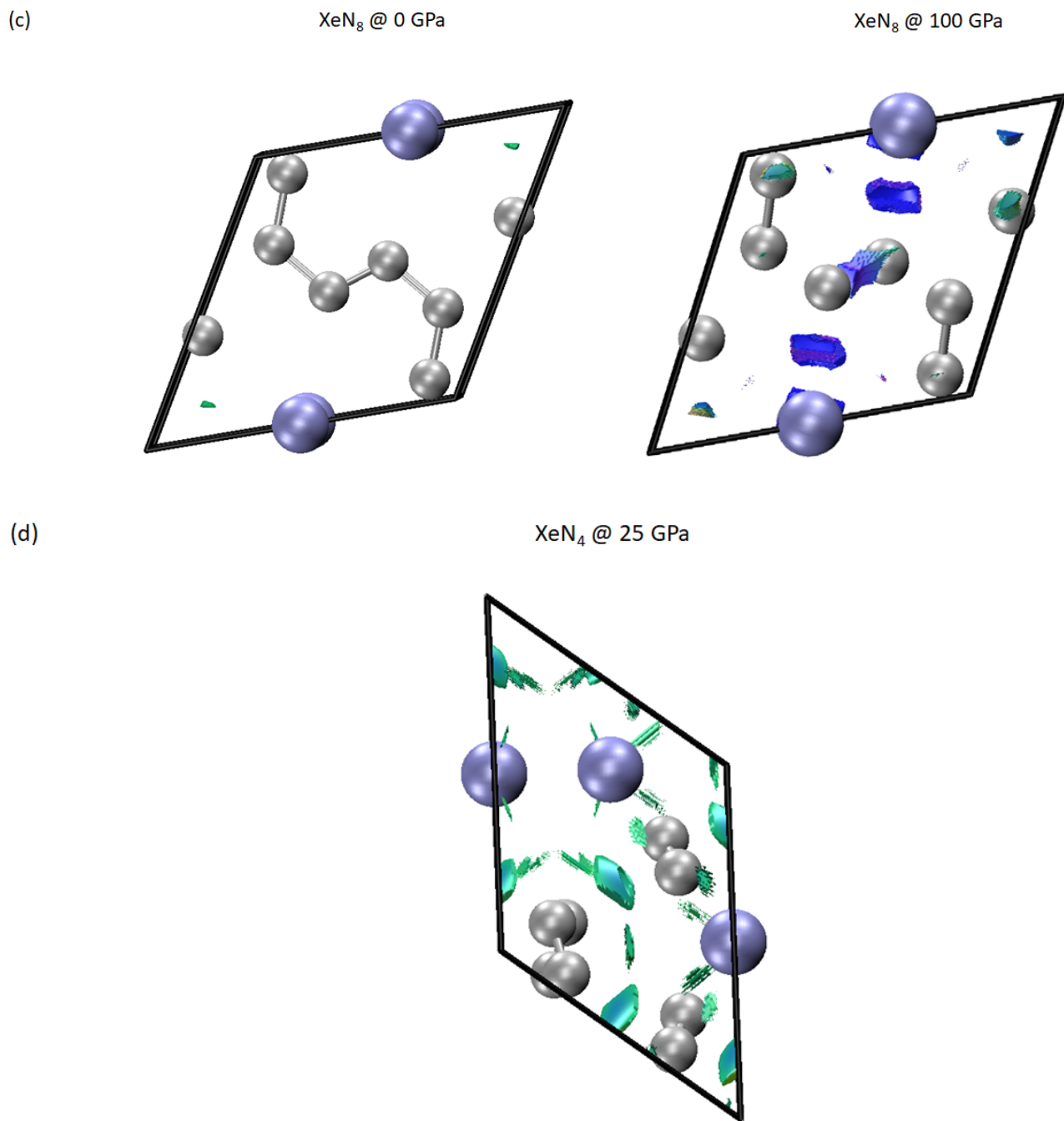


(b)

XeN<sub>18</sub> @ 100 GPa





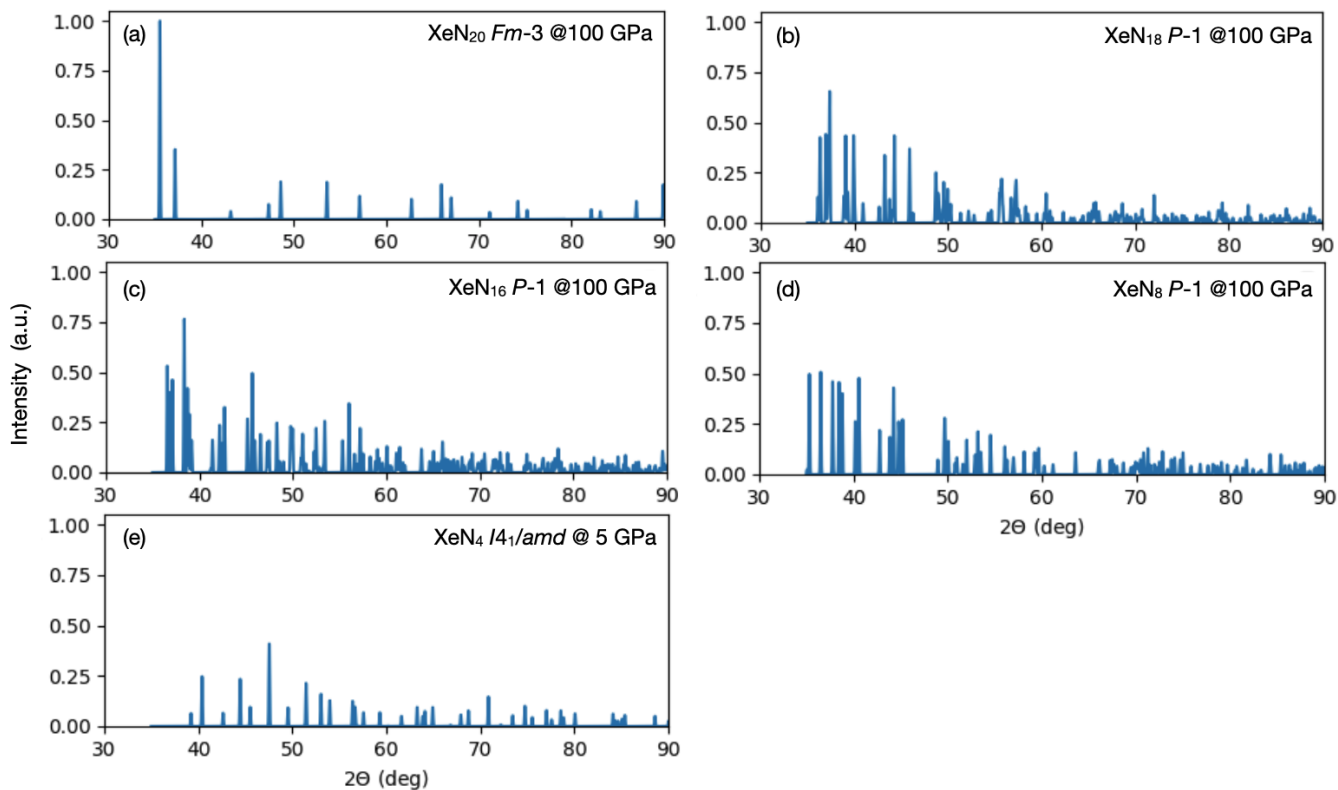


**Figure S21.** Reduced gradient isosurfaces ( $s=0.5$  a.u.) for  $\text{XeN}_4$ ,  $\text{XeN}_8$ ,  $\text{XeN}_{18}$ , and  $\text{XeN}_{20}$  under specific pressure. Xe atoms are depicted in blue and N atoms in gray. Surfaces are colored according to the usual blue-green-red color scheme associated with the sign of the second Hessian eigenvalue ( $<0$ , blue;  $=0$ , green;  $>0$  red). Only negative/nil values are depicted for the sake of clarity.

As shown in Figure S21, NCI analysis is in line with electron density difference maps: in most cases, only moderately stabilizing (dispersive) interactions are observed, associated to greenish volumes. Stronger interactions are observed only in the case of  $\text{XeN}_{18}$  and  $\text{XeN}_8$  at 100 GPa, in line with the observed slight charge transfer (*vide supra*).

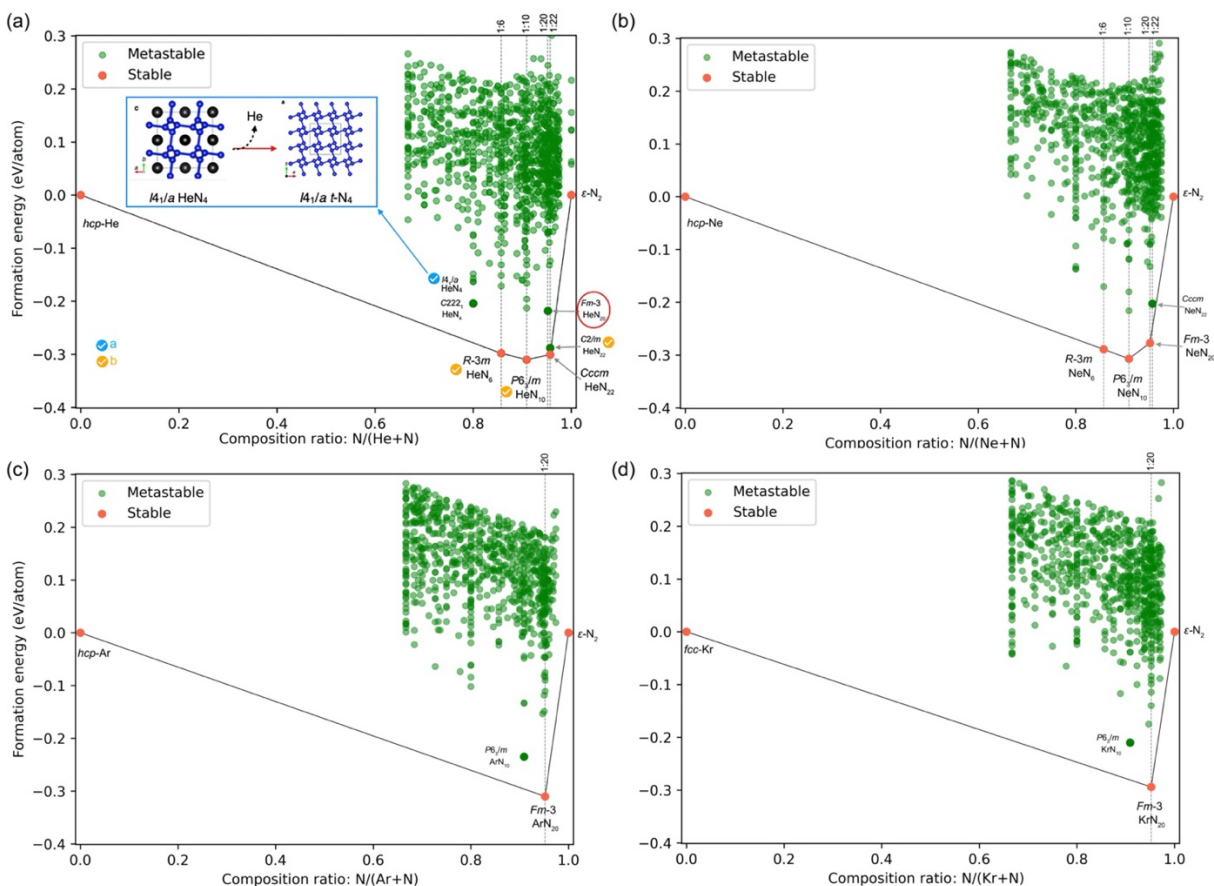
## S8. Simulated X-ray diffraction patterns

As X-Ray diffraction (XRD) is routinely used to identify crystalline phases and can provide unit cell information, XRD spectrums were produced using VESTA software. The wavelength  $\lambda$  of 1.48 Å was used for XRD simulations. The calculated XRD spectrum of the  $Fm-3$  XeN<sub>20</sub>,  $P-1$  XeN<sub>18</sub>,  $P-1$  XeN<sub>8</sub>, and  $I4_1/amd$  XeN<sub>4</sub> are presented in Figure S22.



**Figure S22.** Calculated X-ray diffraction patterns for  $Fm-3$  XeN<sub>20</sub>,  $P-1$  XeN<sub>18</sub>,  $P-1$  XeN<sub>8</sub>, and  $I4_1/amd$  XeN<sub>4</sub> under specific pressure.

## S9. $Fm-3$ Ng@N<sub>20</sub> compounds (Ng=He, Ne, Ar, and Kr)

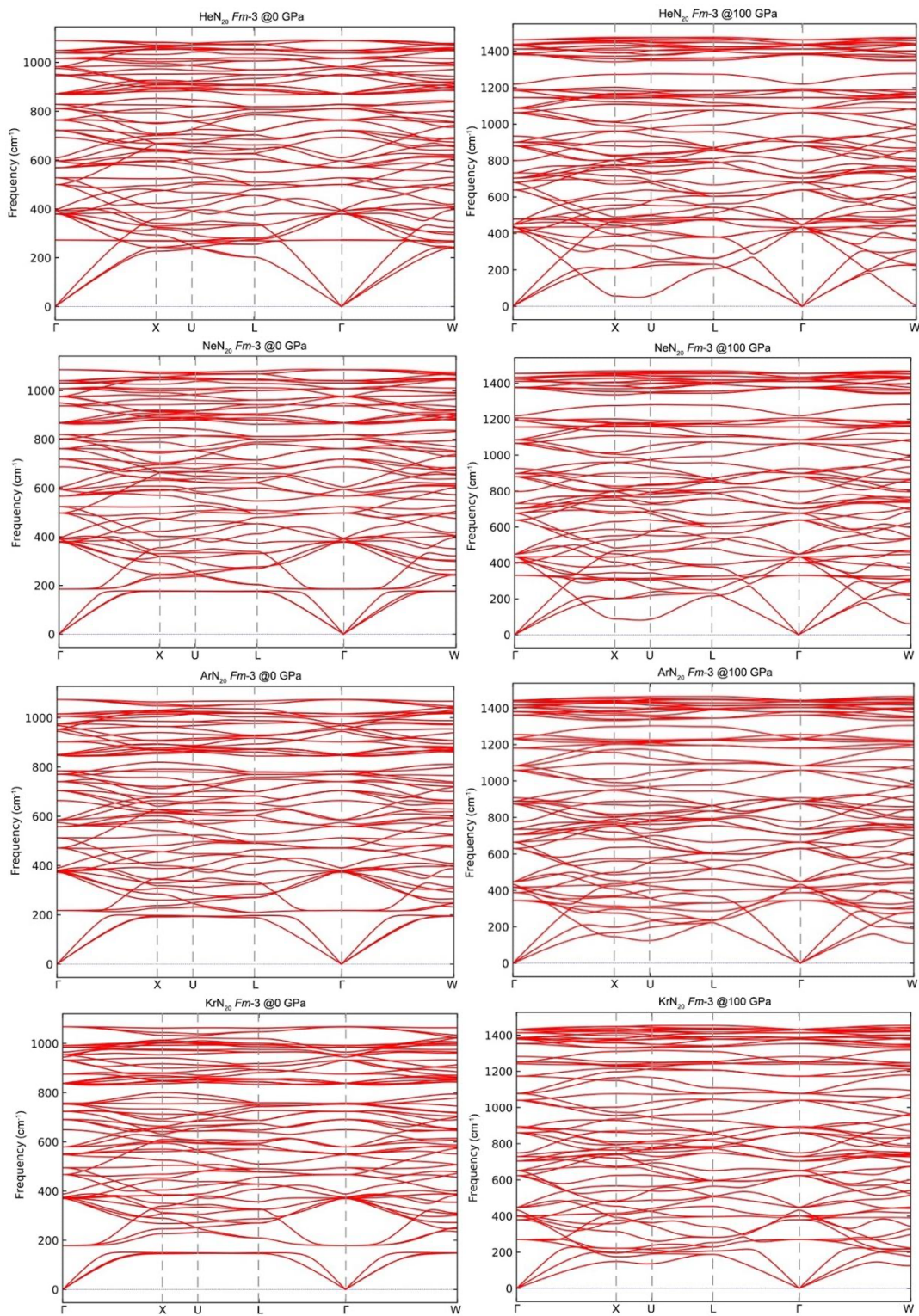


**Figure S23.** Convex-hull diagrams for the Ng–N system (Ng = He, Ne, Ar, and Kr) under 100 GPa with PBE level of theory. The red dots denote thermodynamic stable phases, and the green ones represent metastable structures. The space group is indicated for each novel dynamical stable structure. Here, the  $\epsilon$ -N<sub>2</sub> and accordingly ground state Ng phases (*hcp*-He, *hcp*-Ne, *hcp*-Ar, and *fcc*-Kr) were used to calculate the formation enthalpies at 100 GPa and 0 K. To be noted, only the data (composition ratio > 2/3) are depicted. In (a) He–N system, the blue and yellow circles indicate the previously reported phases by Li *et al.*<sup>29</sup> and Hou *et al.*<sup>31</sup>, respectively.

Remarkably, all the  $Fm-3$  Ng@N<sub>20</sub> are shown on their ( $\epsilon$ -N<sub>2</sub>+Ng) convex hulls at 100 GPa, except the HeN<sub>20</sub>. The  $Fm-3$  HeN<sub>20</sub> is located at 82 meV/atom above the convex hull. Moreover, the previously reported metastable  $I4_1/a$  HeN<sub>4</sub><sup>29</sup> are located above the convex hull, while the reported stable  $R-3m$  HeN<sub>6</sub>,  $P6_3/m$  HeN<sub>10</sub>, and  $C2/m$  He(N<sub>2</sub>)(N<sub>20</sub>) are on the convex hull at our level of theory (PBE).<sup>31</sup> Interestingly, a distorted  $C2/m$  He(N<sub>2</sub>)(N<sub>20</sub>) with rotated N<sub>2</sub> molecular, named  $Cccm$  He(N<sub>2</sub>)(N<sub>20</sub>) is predicted on the convex hull in our structure searches, whose formation enthalpy is slightly (13 meV/atom) below that of reported  $C2/m$  He(N<sub>2</sub>)(N<sub>20</sub>).<sup>31</sup>

In the Ne–N system, the  $R-3m$  NeN<sub>6</sub> and  $P6_3/m$  NeN<sub>10</sub> crystallize in the same structure as HeN<sub>6</sub> and HeN<sub>10</sub>. They are predicted to be the thermodynamic phases. Moreover, the  $Fm-3$  NeN<sub>20</sub> turns out to be on the convex hull.

For the Ar/Kr–N systems, only the  $Fm-3$  ArN<sub>20</sub> and KrN<sub>20</sub> are calculated on the convex hull, as shown in Figure S23.



**Figure S24.** Phonon dispersion curves of  $Fm-3$  Ng@N<sub>20</sub> compounds (Ng = He, Ne, Ar, and Kr) at 0 and 100 GPa.

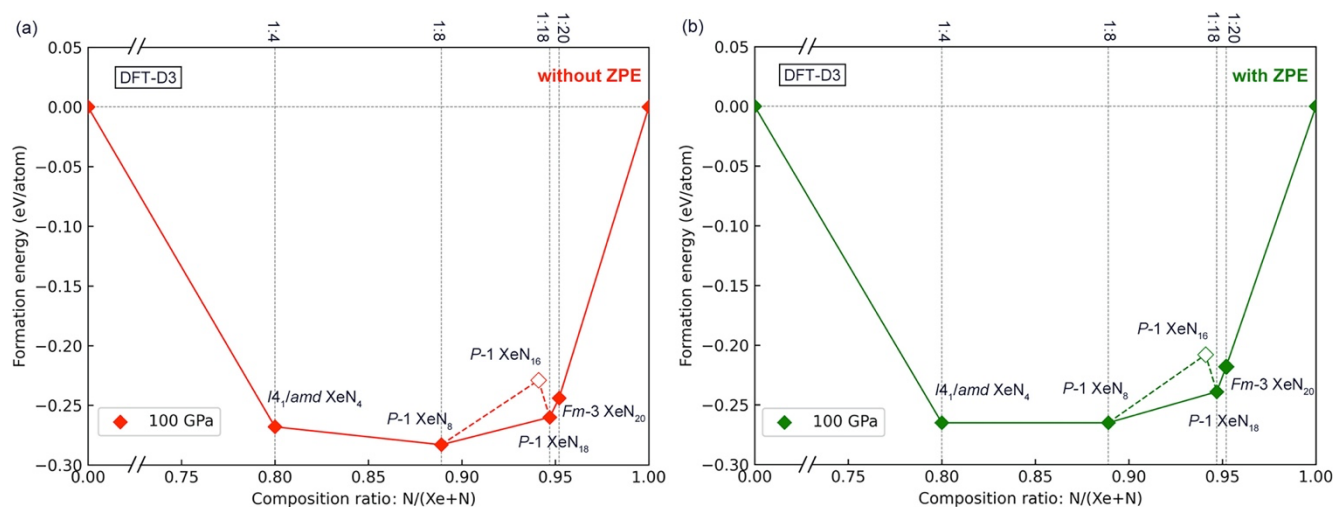
No imaginary phonon modes are found in the whole Brillouin zone, confirming their dynamical stability.

**Table S7.** Structural parameters of dynamically stable Noble-nitride phases (distances in Å, angles in °, energy in eV/atom) at the SCAN+rVV10 level of theory.

Phase	Pressure (GPa)	Space group	Z	Lattice parameters	Atomic coordinates (fractional)
KrN <sub>20</sub>	100	<i>Fm-3</i>	4	a=b=c=7.953 α=β=γ=90.0	Kr (0.000, 0.000, 0.500); N (0.000, 0.222, 0.167), (0.141, 0.141, 0.141)
KrN <sub>20</sub>	0	<i>Fm-3</i>	4	a=b=c=8.559 α=β=γ=90.0	Kr (0.000, 0.000, 0.500); N (0.000, 0.215, 0.169), (0.139, 0.139, 0.139)
ArN <sub>20</sub>	100	<i>Fm-3</i>	4	a=b=c=7.893 α=β=γ=90.0	Ar (0.000, 0.000, 0.500); N (0.000, 0.226, 0.165), (0.141, 0.141, 0.141)
ArN <sub>20</sub>	0	<i>Fm-3</i>	4	a=b=c=8.500 α=β=γ=90.0	Ar (0.000, 0.000, 0.500); N (0.000, 0.218, 0.168), (0.139, 0.139, 0.139)
NeN <sub>20</sub>	100	<i>Fm-3</i>	4	a=b=c=7.790 α=β=γ=90.0	Ne (0.000, 0.000, 0.500); N (0.000, 0.231, 0.163), (0.141, 0.141, 0.141)
NeN <sub>20</sub>	0	<i>Fm-3</i>	4	a=b=c=8.408 α=β=γ=90.0	Ne (0.000, 0.000, 0.500); N (0.000, 0.221, 0.167), (0.139, 0.139, 0.139)
HeN <sub>20</sub>	100	<i>Fm-3</i>	4	a=b=c=7.785 α=β=γ=90.0	He (0.000, 0.000, 0.500); N (0.000, 0.232, 0.161), (0.141, 0.141, 0.141)
HeN <sub>20</sub>	0	<i>Fm-3</i>	4	a=b=c=8.456 α=β=γ=90.0	He (0.000, 0.000, 0.500); N (0.000, 0.220, 0.166), (0.140, 0.140, 0.140)



## S10. The effect of zero-point energy



**Figure S25.** Convex-hull diagrams (a) without and (b) with ZPE correction (DFT-D3 level of theory) for the Xe-N system under 100 GPa regarding the *hcp*-Xe ( $P6_3/m$ ) and  $\epsilon$ -N<sub>2</sub> ( $R-3c$ ).

The effect of zero-point energy (ZPE) on the stability of Xe-N compounds is studied. We find that the inclusion of ZPE only moderately shifts the stability figures but doesn't change the phase stability order.

**Table S8.** Calculated enthalpies of the predicted XeN<sub>x</sub> compounds using DFT-D3.

Structure	Space Group	Pressure (GPa)	ZPE (eV/atom)	Enthalpy with ZPE (eV/atom)	Formation Enthalpy with ZPE (eV/atom)
Xe	$P6_3/m$	100	0.037	16.077	0.000
XeN <sub>4</sub>	$I4_1/amd$	100	0.119	0.871	-0.264
XeN <sub>8</sub>	$P-1$	100	0.143	-0.789	-0.265
XeN <sub>16</sub>	$P-1$	100	0.151	-1.708	-0.208
XeN <sub>18</sub>	$P-1$	100	0.152	-1.855	-0.239
XeN <sub>20</sub>	$Fm-3$	100	0.158	-1.928	-0.218
$\epsilon$ -N <sub>2</sub>	$R-3c$	100	0.136	-2.599	0.000

## S11. References

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