

***Supporting Information for***

**Two dimensional Janus Ti-trihalides monolayers with half-metallic characteristic, Mott insulator and tunable magnetic anisotropy**

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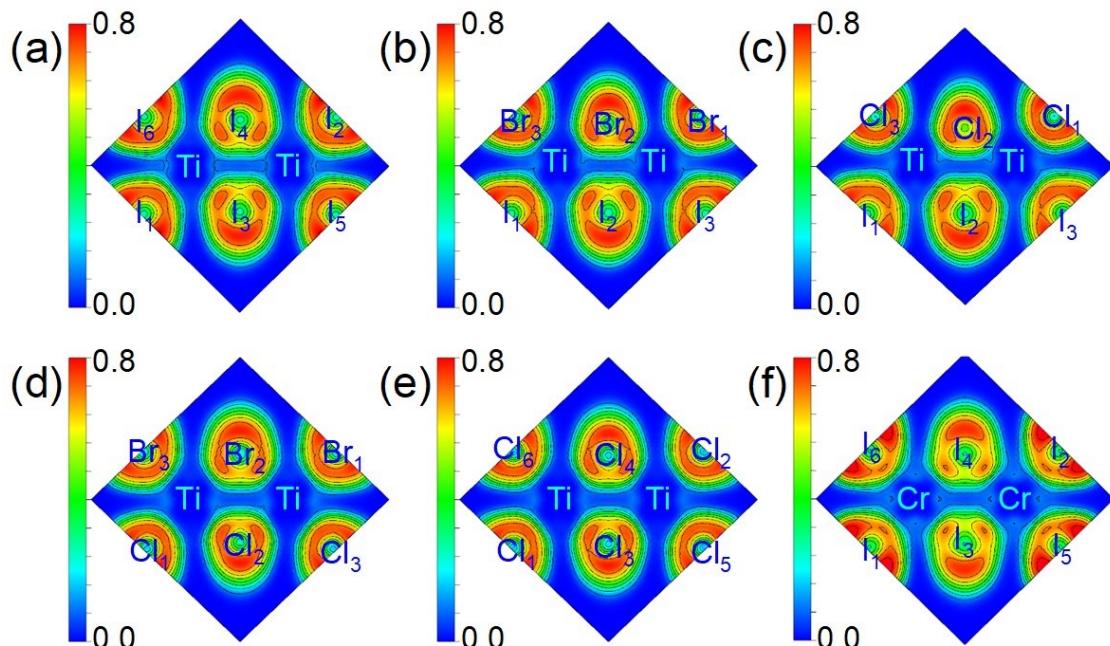
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**Table S1.** The bond lengths ( $d_1$ ,  $d_2$ ) and bond angles ( $\theta_1$ ,  $\theta_2$ ) of each monolayer. Local magnetic moments of Ti/Cr ( $M_{\text{Ti}}/M_{\text{Cr}}$ ) and  $X/Y$  ( $M_X/M_Y$ ) atoms.

	$d_1$ (Å)	$d_2$ (Å)	$\theta_1$ (°)	$\theta_2$ (°)	$M_{\text{Ti}}$ ( $\mu_{\text{B}}$ )	$M_X$ ( $\mu_{\text{B}}$ )	$M_Y$ ( $\mu_{\text{B}}$ )
TiI <sub>3</sub>	2.833	2.833	92.224	92.224	1.067	-0.031	-0.031
Ti <sub>2</sub> Br <sub>3</sub> I <sub>3</sub>	2.627	2.824	98.251	89.375	1.054	-0.023	-0.032
Ti <sub>2</sub> Cl <sub>3</sub> I <sub>3</sub>	2.481	2.821	104.063	87.779	1.046	-0.018	-0.033
Ti <sub>2</sub> Cl <sub>3</sub> Br <sub>3</sub>	2.464	2.609	98.224	91.107	1.023	-0.018	-0.022
TiCl <sub>3</sub>	2.457	2.457	94.855	94.855	1.011	-0.017	-0.017

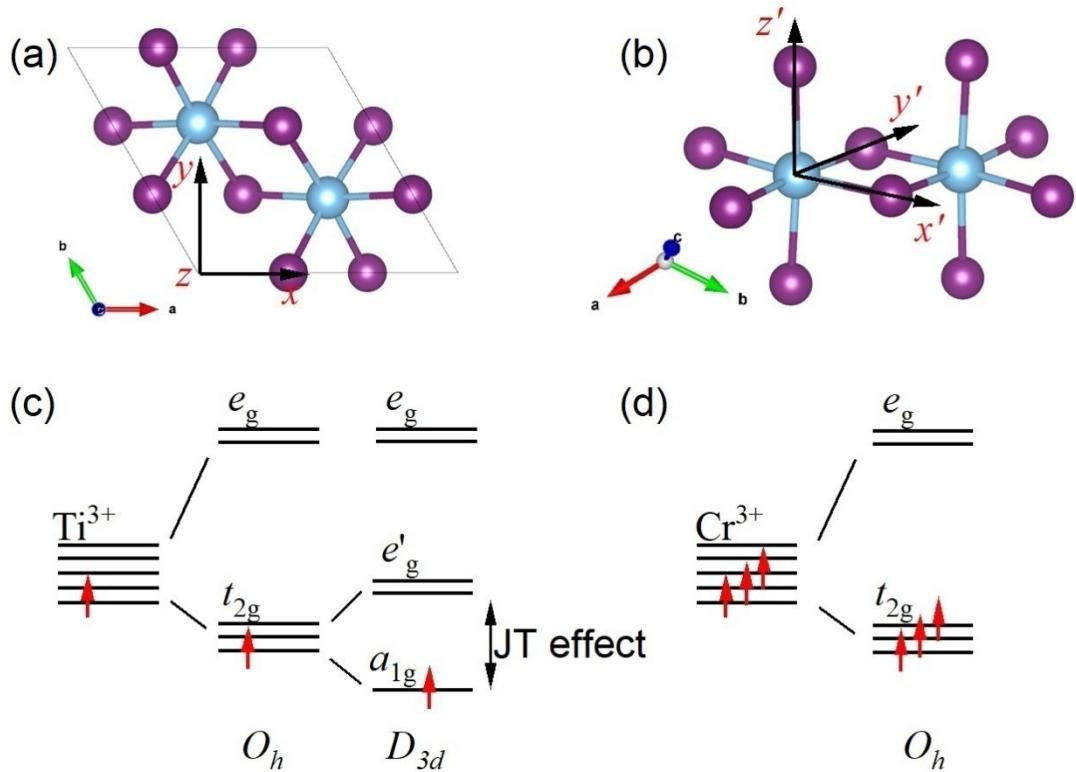
  

	$d_1$ (Å)	$d_2$ (Å)	$\theta_1$ (°)	$\theta_2$ (°)	$M_{\text{Cr}}$ ( $\mu_{\text{B}}$ )	$M_X$ ( $\mu_{\text{B}}$ )	$M_Y$ ( $\mu_{\text{B}}$ )
CrI <sub>3</sub>	2.747	2.747	92.772	92.772	3.400	-0.134	-0.134

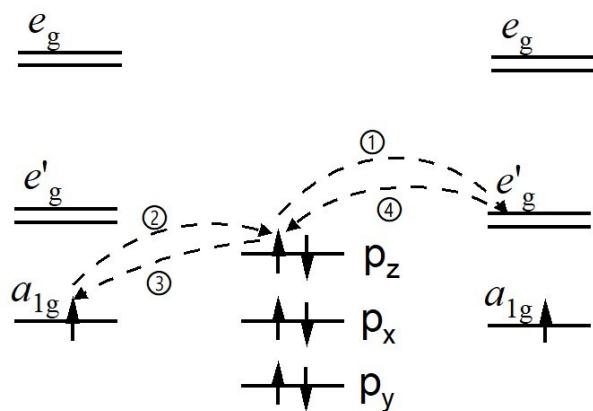


**Fig. S1.** ELF of the (a) TiI<sub>3</sub>, (b) Ti<sub>2</sub>Br<sub>3</sub>I<sub>3</sub>, (c) Ti<sub>2</sub>Cl<sub>3</sub>I<sub>3</sub>, (d) Ti<sub>2</sub>Cl<sub>3</sub>Br<sub>3</sub>, (e) TiCl<sub>3</sub>, (f) CrI<sub>3</sub> monolayer. ELF=1 (red)

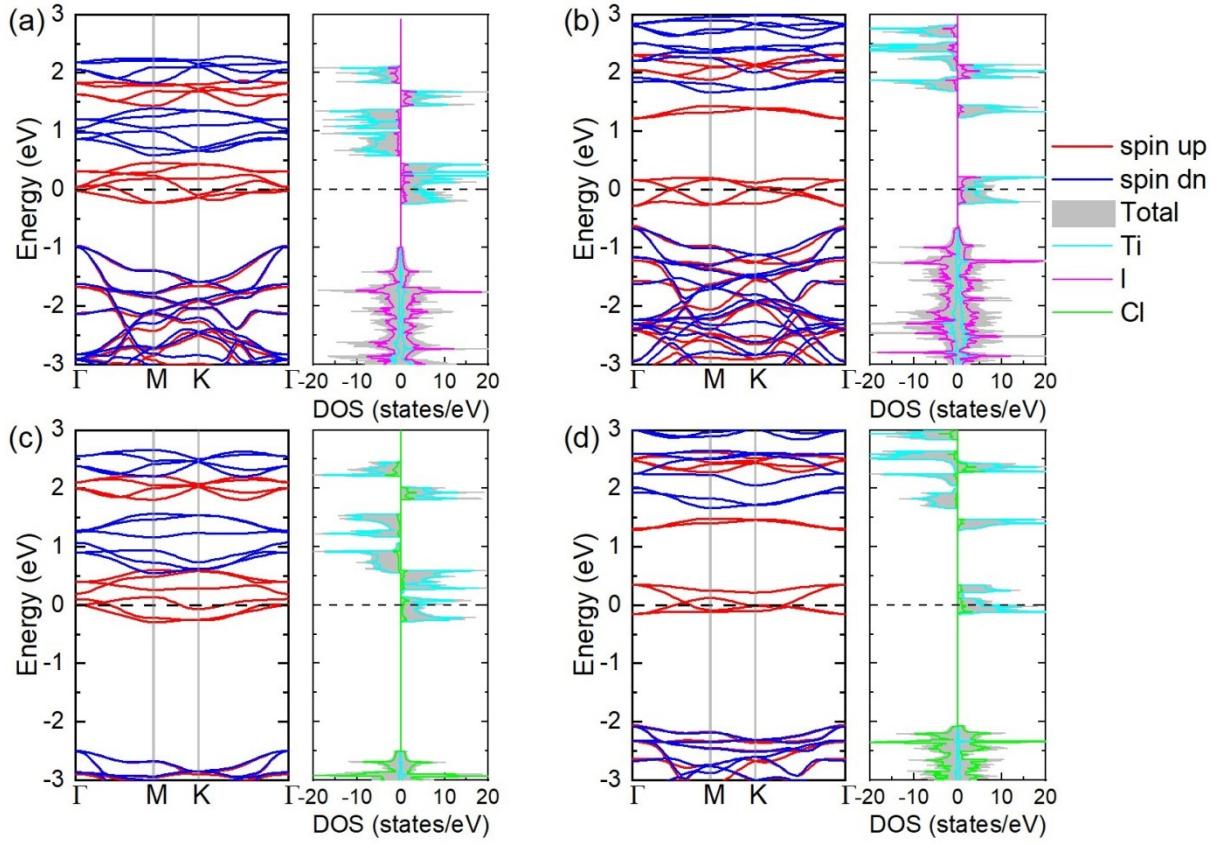
and 0 (blue) indicate the accumulation and depletion of electron density, respectively, and the range chosen here is 0 to 0.8.



**Fig. S2.** Edge-sharing  $\text{TiI}_6$  octahedra with (a) the global  $xyz$  coordinate system and (b) the local  $x'y'z'$  coordinate system. The crystal field splitting of  $d$  levels in (c)  $\text{Ti } 3d^1$  and (d)  $\text{Cr } 3d^3$ .



**Fig. S3.** Schematic plot of FM superexchange interactions in  $\text{TiI}_3$  monolayer.



**Fig. S4.** The band structure and PDOS of the half-metallic state with (a)  $t_{2g}^1$  and (b)  $e_g^1 a_{1g}^0$  configurations for  $\text{TiI}_3$  monolayer. The band structure and PDOS of the half-metallic state with (c)  $t_{2g}^1$  and (d)  $e_g^1 a_{1g}^0$  configurations for  $\text{TiCl}_3$  monolayer.

Two different half-metallic states were obtained by the PBE and the occupation-matrix-control PBE+ $U_{\text{eff}}$  method in Fig. S4. The steps of the occupation-matrix-control PBE+ $U_{\text{eff}}$  method are described. First of all, occupation-matrix-control PBE+ $U_{\text{eff}}$  starts structure optimization with input file OCCMATRIX initializing the occupancy density matrix. Then, the relaxed structure is used to perform a self-consistent-field (SCF) calculation with occupation control. The output CHGCAR and WAVECAR are read for a next relaxation without occupation control. Finally, the relaxed structure is used for the SCF calculation (this usually also requires a step of occupation control calculation first, e.g. reading the occupation density matrix output from the previous step, otherwise the expected electronic structure may not be obtained).

It should be pointed out that the complete SCF PBE+ $U_{\text{eff}}$  calculation without occupation control can also obtain twin orbital-order phases, but this is relatively difficult because the self-consistent solution often depends on the parameter settings.<sup>1-3</sup> Following is the CONTCAR file (used for VASP) for  $a_{1g}^1 e_g^0$  and  $e_g^1 a_{1g}^0$  configurations of the TiCl<sub>3</sub> monolayer in a complete SCF PBE+ $U_{\text{eff}}$  calculation.

$a_{1g}^1 e_g^0$  electronic configuration

1.0000000000000000		
6.2681495088052177	0.0000000000000000	0.0000000000000000
-3.1340747544026133	5.4283767093442652	0.0000000000000000
0.0000000000000000	0.0000000000000000	20.0000000000000000

C1    Ti

6        2

Direct

0.0000000000000000	0.6435794387893308	0.4301205835494915
0.0000000000000000	0.3564205612106697	0.5698794164505085
0.3564205612106697	0.3564205612106697	0.4301205835494915
0.6435794387893308	0.6435794387893308	0.5698794164505085
0.6435794387893308	0.0000000000000000	0.4301205835494915
0.3564205612106697	0.0000000000000000	0.5698794164505085
0.3333333333333333	0.6666666666666666	0.5000000000000000
0.6666666666666667	0.3333333333333334	0.5000000000000000

$e_g^1 a_{1g}^0$  electronic configuration

1.000000000000000

6.1086044213302273	0.000000000000000	0.000000000000000
-3.0543022106651172	5.2902066105419117	0.000000000000000
0.000000000000000	0.000000000000000	20.000000000000000

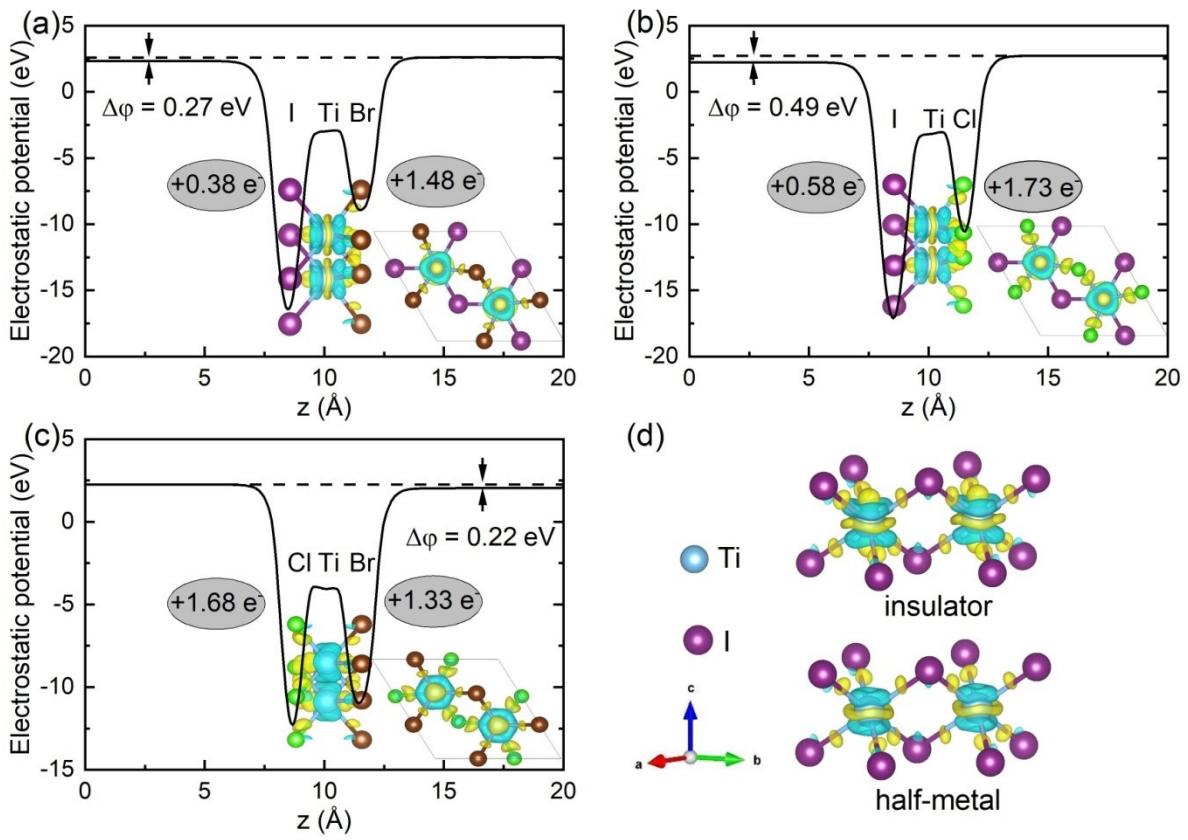
C1 Ti

6 2

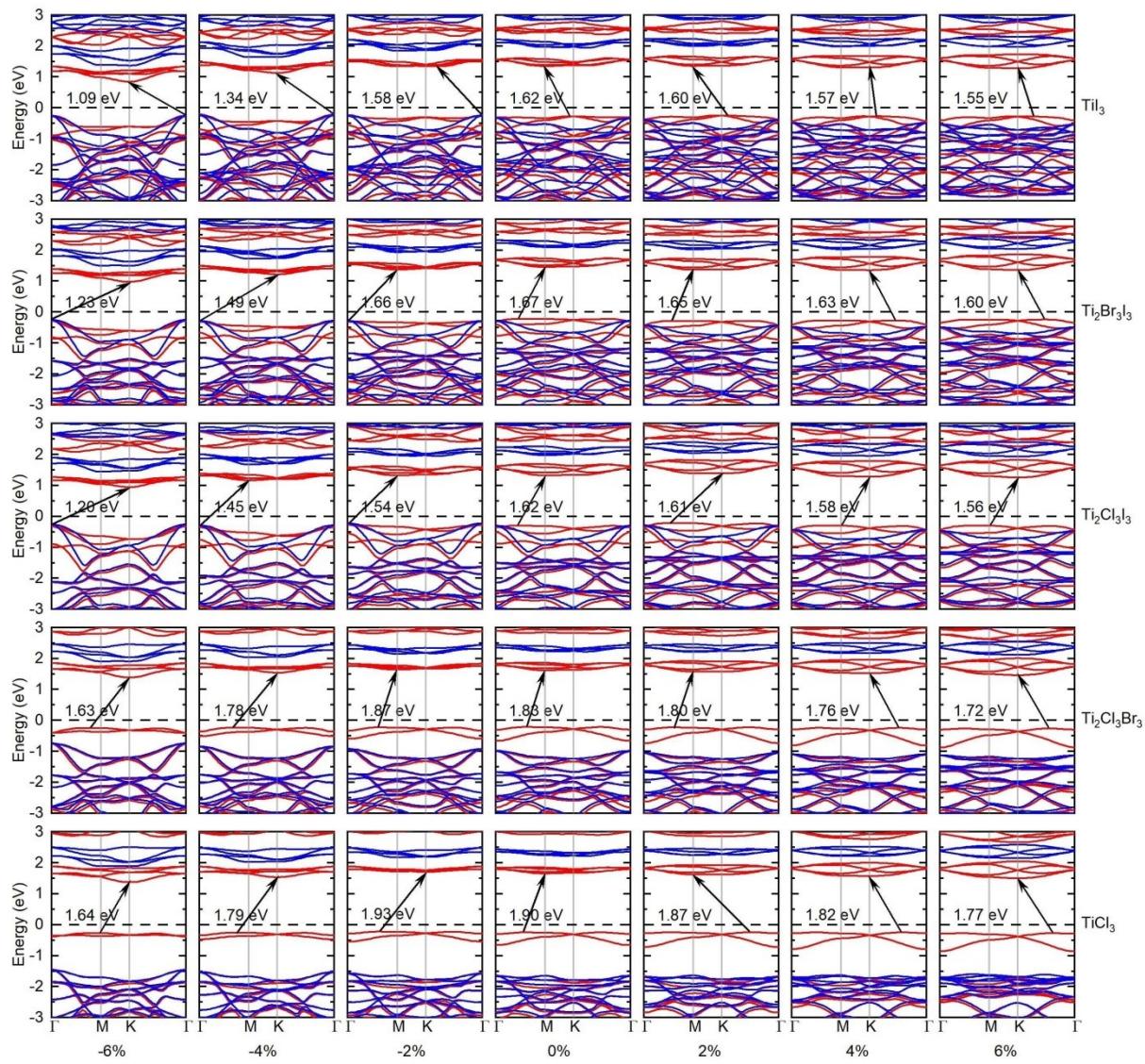
Direct

0.000000000000000	0.6589168856719224	0.4299829405071204
0.000000000000000	0.3410831143280780	0.5700170594928793
0.3410831143280780	0.3410831143280780	0.4299829405071204
0.6589168856719224	0.6589168856719224	0.5700170594928793
0.6589168856719224	0.000000000000000	0.4299829405071204
0.3410831143280780	0.000000000000000	0.5700170594928793
0.333333333333333	0.666666666666666	0.500000000000000
0.666666666666667	0.333333333333334	0.500000000000000

One can perform SCF calculations with these two relaxed structures, and the results are consistent with those of the occupation-matrix-control PBE+ $U_{\text{eff}}$  method.



**Fig. S5.** Electrostatic potential energy and differential charge density of the Janus (a)  $\text{Ti}_2\text{Br}_3\text{I}_3$ , (b)  $\text{Ti}_2\text{Cl}_3\text{I}_3$  and (c)  $\text{Ti}_2\text{Cl}_3\text{Br}_3$  monolayers.  $\Delta\varphi$  is the electrostatic potential difference and the value of isosurface is set to  $0.015 \text{ e/Bohr}^3$ . (d) Differential charge density of the  $\text{TiI}_3$  monolayer, where the value of isosurface is set to  $0.01 \text{ e/Bohr}^3$ . Yellow and blue regions represent charge gain and loss, respectively.



**Fig. S6.** The band structures of Ti<sub>2</sub>X<sub>3</sub>Y<sub>3</sub> monolayers at in-plane biaxial strain of -6% to 6%.

## References

1. B. Meredig, A. Thompson, H. A. Hansen, C. Wolverton and A. van de Walle, *Phys. Rev. B*, 2010, **82**, 195128.
2. X. Fang, B. Zhou, N. Sun, L. Fu and X. Wang, *Phys. Chem. Chem. Phys.*, 2022, **24**, 4374-4383.
3. Y. P. Wang and M. Q. Long, *Phys. Rev. B*, 2020, **101**, 024411.