## Supplementary Material:

# High-resolution rovibrational and rotational spectroscopy of $\mathrm{HCCCO}^{+}$ 

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## 1 Frequency calibration

The high-resolution results shown in Fig. 1 of the main paper have been recorded in sections of about $1.5 \mathrm{~cm}^{-1}$, the scanning range of the OPO. Within such mode-hop-free sections, the precision of the measurement is a few $0.0001 \mathrm{~cm}^{-1}$, but the total accuracy is only on the order of $0.001 \mathrm{~cm}^{-1}$ (see also comments below). For illustration purposes in that Figure, such sections have been combined with the help of overlapping prominent spectral features (the applied shifts are typically much less than $0.001 \mathrm{~cm}^{-1}$ ), whereas for the analysis of the data the untouched measurements have been used.

Concerning the accuracy of our high-resolution data, it became clear during the analysis that the alignment of our wavemeter had not been optimal during the measurements and therefore showed slightly lower wavenumber readings. Therefore, the spectroscopic data were recalibrated using the $3204.730595 \mathrm{~cm}^{-1}$ line of neutral $\mathrm{C}_{2} \mathrm{H}_{2}$ contained within an absorption cell. Using this reference line (taken from the HITRAN database ${ }^{1}$ ), we shifted our original data up by $0.009 \mathrm{~cm}^{-1}$. With this, we assume the accuracy of the data given in the Tables to be a few $0.001 \mathrm{~cm}^{-1}$. The precision of the data, however, is much better than this, on the order of a few $0.0001 \mathrm{~cm}^{-1}$.

## 2 Early ro-vibrational measurements applying LIICG

In summer 2020 we did first rovibrational measurements on $\mathrm{HCCCO}^{+}$using the LIICG method. ${ }^{2}$ With the introduction of LOS, we continued the measurements using this novel method. The signal-to-noise ratio of LOS is superior as seen in the Figure below.


Figure S1: Comparison of measurements of the $\nu_{1} P(6)$ line of $\mathrm{HCCCO}^{+}$using LOS and LIICG. For LIICG, $\mathrm{HCCCO}^{+}$-He complexes are generated at 4 K in the ion trap, which are finally counted (on the order of 450 , see y2-axis). If the laser is in resonance with the $P(6)$ line of $\mathrm{HCCCO}^{+}$, the counts of $\mathrm{HCCCO}^{+}-\mathrm{He}$ decrease (dip at resonance). For LOS, on the other hand, one simply counts the $\mathrm{HCCCO}^{+}$ions which are kicked out of the ion trap at resonance (about 4000, see y1-axis). As visible, the $\mathrm{S} / \mathrm{N}$ ratio of LOS is superior.

## 3 Comparison with earlier detection of $\nu_{1}$

The Figure below shows a comparison of the low-resolution measurements of $\nu_{1}$ and $\nu_{2}+\nu_{4}$ of this work (Fig. 1a) with a former measurement of tagged $\mathrm{Ne}-\mathrm{HCCCO}^{+}$published by Thorwirth et al. ${ }^{3}$. Both measurements exhibit two band structures, but the Ne-tagged measurement (red trace in Fig. S2) is blue-shifted by about $11 \mathrm{~cm}^{-1}$. As the measurements in this work are corroborated by the well-calibrated high-resolution measurements (see Fig. 1 b and c), and a blueshift of $11 \mathrm{~cm}^{-1}$ caused by the Ne-tag seems highly unlikely, we suspect a calibration error in measurement of the red trace. Together with theoretical predictions, this formerly mislead us to assign the left red peak to $\nu_{1}$ and only speculate about the nature of the right peak, so that the $\nu_{2}+\nu_{4}$ band could not be assigned with confidence in that work. ${ }^{3}$


Figure S2: Comparison of the measurement of $\nu_{1}$ and $\nu_{2}+\nu_{4}$ of $\mathrm{HCCCO}^{+}$(this work, Fig. 1a) with infrared photodissociation (IRPD) of $\mathrm{Ne}-\mathrm{HCCCO}^{+}$performed by Thorwirth et al. ${ }^{3}$.

## 4 Resonance System

The rotational constants listed in Table 6 of the main manuscript yield experimental rotationvibration interaction constants $\left(\alpha_{1}=13.63(4) \mathrm{MHz}, \alpha_{2+4}=23.07(5) \mathrm{MHz}, \alpha_{2+5+7}=11.97(5)\right.$ MHz ) significantly deviating from those estimated from the calculation (see Table 1 main manuscript; $\alpha_{1}=7.31 \mathrm{MHz}, \alpha_{2+4} \approx \alpha_{2}+\alpha_{4}=31.78 \mathrm{MHz}, \alpha_{2+5+7} \approx \alpha_{2}+\alpha_{5}+\alpha_{7}=3.42 \mathrm{MHz}$ ). Although the presence of the $\nu_{2}+\nu_{4}$ combination band might have been speculated (cf., Ref. ${ }^{3}$ ), the appearance of a third band in close proximity to $\nu_{1}$ came somewhat as a surprise. A possible explanation for this finding is the existence of a (three-level) resonance system in which an initially dark state gains intensity from its resonance partners. In our case most likely a Fermi-resonance system is observed as the $\nu_{1}$ fundamental as well as the $\nu_{2}+\nu_{4}$ combination band are both of $\Sigma^{+}$-symmetry. Hence, the initially dark state of the Fermi-resonance system has to posses at least one component of $\Sigma^{+}$-symmetry in order to borrow intensity from the bright states. This criterion is in general fulfilled by bands such as the $\nu_{2}+\nu_{5}+\nu_{7}$ combination band ( $\Sigma^{+}, \Sigma^{-}$and $\Delta$-symmetry) due to the two involved bending motions.

Within a simple second order de-perturbation approach all appearing coupling terms between the individual resonating bands should be equal in magnitude but opposite in sign $\left(\frac{\left.\left|\left\langle\psi_{i}\right| W_{i j}\right| \psi_{j}\right\rangle\left.\right|^{2}}{E_{i}-E_{j}}=-\frac{\left.\left|\left\langle\psi_{j}\right| W_{j i}\right| \psi_{i}\right\rangle\left.\right|^{2}}{E_{j}-E_{i}}\right)$. Hence, the sum of their perturbed term values (Table S3 in the Supplementary Material), initially calculated from combination differences of the observed transitions, should be equal to the sum of their unperturbed term energies. An effective fit of the combined term values yields a rotational constant of the mixed states $B_{\text {eff }}=\frac{B_{a}+B_{b}+B_{c}}{3}=4444.54(19) \mathrm{MHz}$ which corresponds to an effective rotation-vibration interaction parameter $\alpha_{\text {eff }}=\frac{\alpha_{a}+\alpha_{b}+\alpha_{c}}{3}=16.05(19) \mathrm{MHz}$. Taking the calculated rotationvibration interaction constants of $\nu_{1}$ and the $\nu_{2}+\nu_{4}$ combination band as given above, the initially dark state of the resonance system should exhibit a rotation-vibration interaction constant $\alpha_{d a r k}=9.1(6) \mathrm{MHz}$. A reasonable estimate for the rotation-vibration coupling
constant $\alpha_{\nu_{2}+\nu_{5}+\nu_{7}, l=0}$ of its two $\ell=0$ ( $\Sigma^{+}$- and $\Sigma^{-}$-symmetry) components can be estimated from the values given in Table 1 in the main manuscript. The sum of the rotationvibration coupling constants of the involved fundamentals has to be raised and lowered by the difference between the $\ell$-type-doubling constants of the two bending motions yielding $\alpha_{\nu_{2}+\nu_{5}+\nu_{7}, \ell=0}=3.42 \pm 5.26=8.68 /-1.84 \mathrm{MHz}$. The first value lies fairly good within the uncertainties of the derived dark state rotation-vibration coupling constant $\alpha_{\text {dark }}=9.1$ (6) MHz supporting our tentative assignment.

## 5 Line lists and tentative assignment of hot bands of $\mathrm{HCCCO}^{+}$

Two hot bands are observed in panel b of Fig. 1 in the main paper, an intense one with band center at about $3234.60 \mathrm{~cm}^{-1}$ (marked green in the Figure and called "hot1" in the following), and one less intense, with a lower band center at about $3234.06 \mathrm{~cm}^{-1}$ (marked blue in the Figure and called "hot2" in the following). We suspect them to originate from the lowest bending modes ${ }^{3}$ of $\mathrm{HCCCO}^{+}$, the C-C-C bending vibration $\nu_{7}$ at $169^{1} / 174^{2} \mathrm{~cm}^{-1}$, its overtone $2 \nu_{7}$, or the C-C-O bending vibration $\nu_{6}$ at $558^{1} / 553^{2} \mathrm{~cm}^{-1}$ (c.f. Table S1). We base our tentative assignments on characteristic observables such as band intensities, rotational constants $B_{\nu}$ (or rotation-vibration constants $\alpha_{\nu}$ ), doubling constants $q_{\nu}$, and band centers (or anharmonicity constants $\chi$ ).

The line lists of these two hot bands are given below. Due to the congested character of the Q-bands, only well-separated lines are listed here. The given rotational assignment is utilized to derive the term energies of the involved states by calculating combination differences. The hereby obtained ground state term values are listed in Tables S4 and S5, respectively, while those in the upper state of the hot bands are given in Table S6. The rotational constants listed in the mentioned tables refer to a least squares fit of the term values to a standard rigid rotor Hamiltonian neglecting centrifugal distortion.

[^0]A comparison of the obtained values with the results of the quantum chemical calculations given in Table S1 leaves room for two different assignments of the involved vibrational motions. One possible assignment would be "hot1" to originate from the C-C-O bending motion $\nu_{6}\left({ }^{1} \Pi\right)$, while "hot2" originates from the C-C-C bending vibration $\nu_{7}\left({ }^{1} \Pi\right)$. The major argument for this assignment is that their rotational-vibrational coupling constants $\alpha_{\nu}$ agree in magnitude and sign with their predictions. Furthermore, the assigned $\nu_{1}+\nu_{6} \leftarrow \nu_{6}$ ( ${ }^{1} \Pi \leftarrow{ }^{1} \Pi$ ) band would originate above the $\nu_{1}+\nu_{7} \leftarrow \nu_{7}\left({ }^{1} \Pi \leftarrow{ }^{1} \Pi\right)$ band with a relative displacement $\left(0.545 \mathrm{~cm}^{-1}\right)$ which is comparably close to the difference between the calculated anharmonicity constants $\chi$ of these two hot-bands ( $0.76 \mathrm{~cm}^{-1}$, see Table S1). Also, a close inspection of "hot1" reveals resolved $\ell$-type splittings for higher rotational states (see inset in Fig. 1b) which correspond to an $\ell$-type doubling constant $q_{\text {hot } 1}=3.26(5) \mathrm{MHz}$. This value is again reasonably close to the value $q_{6}=3.06 \mathrm{MHz}$ predicted for $\nu_{6}$. On the other side, $\ell$-type doubling is not observed for "hot2" which is somewhat surprising, as the l-type coupling constant of the C-C-C bending vibration, $q_{7}$, is calculated to be more than twice than the one of the C-C-O bending $q_{6}$.

However, taking into account our cryogenic experimental conditions, the assignment given above is difficult to accept, as we do not expect the $\nu_{6}$ bending mode (at $558 \mathrm{~cm}^{-1}$ or about $803 \mathrm{~K})$ to be significantly populated at $T=12 \mathrm{~K}$. Therefore, we alternatively assign the more intense band "hot1" to $\nu_{1}+\nu_{7} \leftarrow \nu_{7}$ and the weaker band "hot2" to $\nu_{1}+2 \nu_{7} \leftarrow 2 \nu_{7}$. Apart from this intensity argument, this assignment is corroborated by quantum levels in "hot2" apparently only existing for $J \geq 2$, which suggest an assignment to a ${ }^{1} \Delta \leftarrow^{1} \Delta$ transition (but in this case we wonder why $2 \nu_{7}$ appears as $\Delta$-state $(\ell=2)$ with a missing $\Sigma$-state $(\ell=0))$. Finally, we note that $\alpha_{h o t 2} \approx-15.96$ has double the value of $\alpha_{h o t 1} \approx-7.88$, as expected for an overtone state.

In summary, we can only tentatively assign the two hot bands to $\nu_{1}+\nu_{7} \leftarrow \nu_{7}, \nu_{1}+2 \nu_{7} \leftarrow$ $2 \nu_{7}$, or $\nu_{1}+\nu_{6} \leftarrow \nu_{6}$. As an additional puzzle, both suggested assignment sets given above lack an explanation of the missing $\ell$-type doubling in "hot2". This requires additional future
investigations of the bending vibrations to resolve the open questions.

## line list of hot-band "hot1" of $\mathrm{HCCCO}^{+}$

Truncated Log file from PGOPHER output. Transitions are given in units of $\mathrm{cm}^{-1}$.

| J'S' | \#' | J"S" | \#" | Observed | Calculated | Obs-Calc |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 13 e | 1 | 14 e | 1 | 3230.3785 | 3230.3780 | 0.0005 |
| 13 f | 1 | 14 f | 1 | 3230.3812 | 3230.3810 | 0.0002 |
| 12 f | 1 | 13 f | 1 | 3230.6862 | 3230.6863 | -0.0001 |
| 11 e | 1 | 12 e | 1 | 3230.9886 | 3230.9885 | 0.0001 |
| 11 f | 1 | 12 f | 1 | 3230.9908 | 3230.9911 | -0.0003 |
| 10 e | 1 | 11 e | 1 | 3231.2928 | 3231.2929 | -0.0001 |
| 10 f | 1 | 11 f | 1 | 3231.2946 | 3231.2953 | -0.0007 |
| 9 e | 1 | 10 e | 1 | 3231.5967 | 3231.5967 | -0.0000 |
| 9 f | 1 | 10 f | 1 | 3231.5984 | 3231.5989 | -0.0005 |
| 8 e | 1 | 9 e | 1 | 3231.8998 | 3231.9000 | -0.0002 |
| 8 f | 1 | 9 f | 1 | 3231.9016 | 3231.9019 | -0.0003 |
| 7 e | 1 | 8 e | 1 | 3232.2027 | 3232.2027 | 0.0000 |
| 7 f | 1 | 8 f | 1 | 3232.2041 | 3232.2044 | -0.0003 |
| 6 e | 1 | 7 e | 1 | 3232.5050 | 3232.5048 | 0.0002 |
| 6 f | 1 | 7 f | 1 | 3232.5063 | 3232.5064 | -0.0001 |
| 5 e | 1 | 6 e | 1 | 3232.8072 | 3232.8064 | 0.0008 |
| 5 f | 1 | 6 f | 1 | 3232.8072 | 3232.8077 | -0.0005 |
|  |  |  |  | Blend | 3232.8071 | 0.0001 |
| 4 f | 1 | 5 f | 1 | 3233.1081 | 3233.1085 | -0.0004 |
| 4 e | 1 | 5 e | 1 | 3233.1081 | 3233.1074 | 0.0007 |
|  |  |  |  | Blend | 3233.1080 | 0.0001 |
| 3 f | 1 | 4 f | 1 | 3233.4086 | 3233.4087 | -0.0001 |
| 3 e | 1 | 4 e | 1 | 3233.4086 | 3233.4079 | 0.0007 |
|  |  |  |  | Blend | 3233.4083 | 0.0003 |
| 2 e | 1 | 3 e | 1 | 3233.7085 | 3233.7078 | 0.0007 |
| 2 f | 1 | 3 f | 1 | 3233.7085 | 3233.7084 | 0.0001 |
|  |  |  |  | Blend | 3233.7081 | 0.0004 |
| 1 f | 1 | 2 f | 1 | 3234.0077 | 3234.0075 | 0.0002 |
| 1 e | 1 | 2 e | 1 | 3234.0077 | 3234.0071 | 0.0006 |
|  |  |  |  | Blend | 3234.0073 | 0.0004 |
| 10 f | 1 | 10 e | 1 | 3234.5608 | 3234.5610 | -0.0002 |
| 9 f | 1 | 9 e | 1 | 3234.5688 | 3234.5688 | 0.0000 |
| 8 f | 1 | 8 e | 1 | 3234.5760 | 3234.5758 | 0.0002 |


| 7 f | 1 | 7 e | 1 | 3234.5822 | 3234.5821 | 0.0001 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 6 f | 1 | 6 e | 1 | 3234.5877 | 3234.5876 | 0.0001 |
| 4 f | 1 | 4 e | 1 | 3234.5962 | 3234.5962 | 0.0000 |
| 3 f | 1 | 2 f | 1 | 3235.4953 | 3235.4945 | 0.0008 |
| 3 e | 1 | 2 e | 1 | 3235.4953 | 3235.4952 | 0.0001 |
|  |  |  |  |  |  |  |

62 observations in total reduced after combining blends to 52

## line list of hot-band "hot2" of $\mathrm{HCCCO}^{+}$

Truncated Log file from PGOPHER output. Transitions are given in units of $\mathrm{cm}^{-1}$. All $\Lambda$-doubled lines are blended (there are $2 \times 22$ lines in the list).

| J'S' | \#' | J"S" | \#" | Observed | Calculated | Obs-Calc |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 10 f | 1 | 11 f | 1 | 3230.7501 | 3230.7495 | 0.0006 |
| 10 e | 1 | 11 e | 1 | 3230.7501 | 3230.7495 | 0.0006 |
| 9 f | 1 | 10 f | 1 | 3231.0529 | 3231.0529 | -0.0001 |
| 9 e | 1 | 10 e | 1 | 3231.0529 | 3231.0529 | -0.0001 |
| 8 f | 1 | 9 f | 1 | 3231.3556 | 3231.3559 | -0.0003 |
| 8 e | 1 | 9 e | 1 | 3231.3556 | 3231.3559 | -0.0003 |
| 7 f | 1 | 8 f | 1 | 3231.6585 | 3231.6584 | 0.0001 |
| 7 e | 1 | 8 e | 1 | 3231.6585 | 3231.6584 | 0.0001 |
| 6 e | 1 | 7 e | 1 | 3231.9602 | 3231.9604 | -0.0002 |
| 6 f | 1 | 7 f | 1 | 3231.9602 | 3231.9604 | -0.0002 |
| 4 f | 1 | 5 f | 1 | 3232.5630 | 3232.5630 | 0.0000 |
| 4 e | 1 | 5 e | 1 | 3232.5630 | 3232.5630 | 0.0000 |
| 3 f | 1 | 4 f | 1 | 3232.8636 | 3232.8635 | 0.0001 |
| 3 e | 1 | 4 e | 1 | 3232.8636 | 3232.8635 | 0.0001 |
| 2 f | 1 | 3 f | 1 | 3233.1634 | 3233.1636 | -0.0002 |
| 2 e | 1 | 3 e | 1 | 3233.1634 | 3233.1636 | -0.0002 |
| 8 f | 1 | 8 e | 1 | 3234.0430 | 3234.0433 | -0.0003 |
| 8 e | 1 | 8 f | 1 | 3234.0430 | 3234.0433 | -0.0003 |
| 7 e | 1 | 7 f | 1 | 3234.0472 | 3234.0472 | -0.0000 |
| 7 f | 1 | 7 e | 1 | 3234.0472 | 3234.0472 | -0.0000 |
| 6 f | 1 | 6 e | 1 | 3234.0508 | 3234.0506 | 0.0002 |
| 6 e | 1 | 6 f | 1 | 3234.0508 | 3234.0506 | 0.0002 |
| 5 f | 1 | 5 e | 1 | 3234.0531 | 3234.0535 | -0.0004 |
| 5 e | 1 | 5 f | 1 | 3234.0531 | 3234.0535 | -0.0004 |
| 3 f | 1 | 2 f | 1 | 3234.9536 | 3234.9537 | -0.0001 |
| 3 e | 1 | 2 e | 1 | 3234.9536 | 3234.9537 | -0.0001 |
| 4 e | 1 | 3 e | 1 | 3235.2504 | 3235.2503 | 0.0001 |
| 4 f | 1 | 3 f | 1 | 3235.2504 | 3235.2503 | 0.0001 |
| 5 e | 1 | 4 e | 1 | 3235.5467 | 3235.5465 | 0.0002 |
| 5 f | 1 | 4 f | 1 | 3235.5467 | 3235.5465 | 0.0002 |
| 6 f | 1 | 5 f | 1 | 3235.8419 | 3235.8422 | -0.0003 |
| 6 e | 1 | 5 e | 1 | 3235.8419 | 3235.8422 | -0.0003 |
| 7 f | 1 | 6 f | 1 | 3236.1375 | 3236.1374 | 0.0001 |
| 7 e | 1 | 6 e | 1 | 3236.1375 | 3236.1374 | 0.0001 |
| 8 e | 1 | 7 e | 1 | 3236.4325 | 3236.4321 | 0.0005 |


| 8 f | 1 | 7 f | 1 | 3236.4325 | 3236.4321 | 0.0005 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 9 e | 1 | 8 e | 1 | 3236.7268 | 3236.7263 | 0.0005 |
| 9 f | 1 | 8 f | 1 | 3236.7268 | 3236.7263 | 0.0005 |
| 10 e | 1 | 9 e | 1 | 3237.0207 | 3237.0200 | 0.0006 |
| 10 f | 1 | 9 f | 1 | 3237.0207 | 3237.0200 | 0.0006 |
| 11 e | 1 | 10 e | 1 | 3237.3123 | 3237.3132 | -0.0009 |
| 11 f | 1 | 10 f | 1 | 3237.3123 | 3237.3132 | -0.0009 |
| 12 e | 1 | 11 e | 1 | 3237.6058 | 3237.6060 | -0.0002 |
| 12 f | 1 | 11 f | 1 | 3237.6058 | 3237.6060 | -0.0002 |

## 6 Quantum chemical calculations

Table S1: Fundamental vibrational wavenumbers of $\mathrm{HC}_{3} \mathrm{O}^{+}$, their rotationvibration coupling constants $\alpha_{\nu}$, and doubling constants $\mathrm{q}_{\nu}$, calculated on the $\mathrm{fc}-\mathrm{CCSD}(\mathrm{T}) / \mathrm{pVTZ}$ level of theory (this work). The given anharmonicity constants $\chi_{1 \nu}$ refer to the hot-bands of the C-H stretching fundamental $\nu_{1}$.

| Mode | Harm <br> $\left[\mathbf{c m}^{-1}\right]$ | Anharm <br> $\left[\mathbf{c m}^{-1}\right]$ | $\mathbf{v}_{i, B E}^{a}$ <br> $\left[\mathbf{c m}^{-1}\right]$ | $\alpha_{i}$ <br> $[\mathbf{M H z ]}$ | $q_{i}$ <br> $[\mathbf{M H z ]}$ | $\chi_{1 i}{ }^{b}$ <br> $\left[\mathbf{c m}^{-1}\right]$ |
| :---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\nu_{1}(\Sigma)$ C-H stretch | 3362.1 | 3228.6 | 3231 | 7.31 | - | -54.52 |
| $\nu_{2}(\Sigma)$ C-O stretch | 2353.1 | 2312.7 | 2316 | 22.84 | - | -1.37 |
| $\nu_{3}(\Sigma)$ C-C stretch | 2097.4 | 2063.6 | 2074 | 14.56 | - | -7.86 |
| $\nu_{4}(\Sigma)$ C-C stretch | 917.8 | 901.3 | 911 | 8.94 | - | 1.14 |
| $\nu_{5}(\Pi)$ C-C-H bending | 773.2 | 757.5 | 773 | -0.49 | 2.13 | -18.56 |
| $\nu_{6}(\Pi)$ C-C-O bending | 552.3 | 547.3 | 558 | -8.00 | 3.06 | -0.39 |
| $\nu_{7}(\Pi)$ C-C-C bending | 181.1 | 166.1 | 169 | -18.93 | 7.39 | -1.49 |

${ }^{b}$ anharmonicity constants of the hot-bands with respect to the $\nu_{1} \mathrm{C}-\mathrm{H}$ stretching fundamental

## 7 Ground state term values

Table S2: Term values for the vibrational ground state $(\nu=0)$ of $\mathrm{HCCCO}^{+}$as derived from all observed transitions of the $\nu_{1} \mathrm{C}-\mathrm{H}$ stretching fundamental and resonating combination bands $\left(\nu_{2}+\nu_{4}, \nu_{2}+\nu_{5}+\nu_{7}\right)$. Two separate networks of states are derived from the experimental data, denoted as a and $b$. This leads to two reference points $\mathrm{E}-\mathrm{D}_{0}=0$ that are not directly linked to each other. These sets are combined based on their y-axis intercept assuming the validity of the rigid rotor approximation.

| $\mathbf{J}$ | Set | $\mathbf{E - D}_{0}^{i}\left[\mathbf{c m}^{-1}\right]$ | $\mathbf{E - D}_{0}^{i i}\left[\mathbf{c m}^{-1}\right]$ | $\mathbf{E}_{\text {combined }}\left[\mathbf{c m}^{-1}\right]$ |
| :---: | :---: | :---: | :---: | :---: |
| 0 | a | 0 | - | 0.00000 |
| 1 | b | - | 0 | $0.29760(1)$ |
| 2 | a | $0.89282(16)$ | - | $0.89282(16)$ |
| 3 | b | - | $1.48835(3)$ | $1.78595(2)$ |
| 4 | a | $2.97627(9)$ | - | $2.97627(9)$ |
| 5 | b | - | $4.16701(9)$ | $4.46461(10)$ |
| 6 | a | $6.25001(12)$ | - | $6.25001(12)$ |
| 7 | b | - | $8.03588(8)$ | $8.33349(9)$ |
| 8 | a | $10.71430(21)$ | - | $10.71430(21)$ |
| 9 | b | - | $13.09521(18)$ | $13.39281(20)$ |
| 10 | a | $16.36868(48)$ | - | $16.36868(48)$ |
| 11 | b | - | $19.34491(67)$ | $19.64251(68)$ |
| 12 | a | $23.21350(52)$ | - | $23.21350(52)$ |
| 13 | b | - | $26.78503(96)$ | $27.08263(97)$ |
| 14 | a | $31.24868(94$ | - | $31.24868(94)$ |
| 15 | b | - | $35.41525(105)$ | $35.71286(107)$ |
| 16 | a | $40.47311(260)$ | - | $40.47311(260)$ |
| $\mathbf{B}_{0}[\mathbf{M H z}$ | $4460.90(5)$ | $4461.00(3)$ | $4460.94(3)$ |  |

## 8 Term values of C-H stretching fundamental and resonat-

## ing combination bands

Table S3: Term energies for the observed vibrationally excited states lying in the $\mathrm{C}-\mathrm{H}$-stretching range of $\mathrm{HCCCO}^{+}$. The given energies are calculated from the observed transitions utilizing the combined ground state term values given in Table S2.

| J | $\nu_{1}=\mathbf{1}\left[\mathrm{cm}^{-1}\right]$ | $\nu_{2}=1, \nu_{4}=\mathbf{1}\left[\mathbf{c m}^{-1}\right]$ | $\nu_{2}=1, \nu_{5}=1, \nu_{7}=1\left[\mathrm{~cm}^{-1}\right]$ |
| :---: | :---: | :---: | :---: |
| 0 | 3237.13231(1) | 3220.46705(1) | $\frac{-}{\text { - }}$ |
| 1 | 3237.42911(16) | 3220.76302(16) | 3237.59800(16) |
| 2 | 3238.02246(2) | 3221.35538(2) | 3238.19314(2) |
| 3 | 3238.91240(9) | 3222.24355(9) | 3239.08143(9) |
| 4 | 3240.09956(10) | 3223.42815(10) | 3240.26938(10) |
| 5 | 3241.58269(12) | 3224.90820(12) | 3241.75301(12) |
| 6 | 3243.36323(9) | 3226.68482(9) | 3243.53421(9) |
| 7 | 3245.43974(21) | 3228.75722(21) | 3245.61172(21) |
| 8 | 3247.81359(20) | $3231.12580(20)$ | 3247.98621(20) |
| 9 | 3250.48364(48) | $3233.79005(48)$ | 3250.65727(48) |
| 10 | 3253.45050(68) | 3236.75096(68) | 3253.62520(68) |
| 11 | 3256.71401(52) | 3240.00706(52) | 3256.88964(52) |
| 12 | $3260.27453(97)$ | 3243.56005(97) | 3260.45164(97) |
| 13 | 3264.13109(94) | 3247.40882(94) | 3264.30933(94) |
| 14 | 3268.28504(107) | 3251.56214(107) | 3268.46464(107) |
| 15 | 3272.73379 (260) | 3255.99306(260) | 3272.91519 (260) |
| 16 | 3277.48248(107) | 3260.72864(107) | 3277.66526(107) |
| 17 | 3282.52468(260) | $3265.76015(260)$ | 3282.70927(260) |
| $\mathrm{B}_{\nu}$ [MHz] | 4447.20(3) | 4437.72(18) | 4448.71(4) |
| $\alpha_{\nu}^{a}$ [MHz] | 13.74(5) | 23.22(19) | 12.23(6) |
| $\alpha_{\nu, \text { calc }}^{b}[\mathrm{MHz}]$ | 7.31 | 31.78 | 3.42 |

## 9 Term values derived from hot-bands

Table S4: Term values for the ground state of the first hot-band (Hot 1) of $\mathrm{HCCCO}^{+}$derived from ground state combination differences. Four separate networks of states are derived from the experimental data, denoted as a,b,c and d, due to the two l-type components of the underlying motion. This leads to four reference points $E-D_{0}=0$ that are not directly linked to each other. These sets are combined based on their y-axis intercept assuming the validity of the rigid rotor approximation.

| $\mathbf{J}_{e, f}$ | Set | E- $\mathrm{D}_{0}^{i}\left[\mathrm{~cm}^{-1}\right]$ | E-D ${ }_{0}^{i i}\left[\mathrm{~cm}^{-1}\right]$ | $\mathbf{E}_{\text {combined }}\left[\mathbf{c m}^{-1}\right]$ |
| :---: | :---: | :---: | :---: | :---: |
| $1_{f}$ | a | 0.0000 | - | 0.0000 |
| $1_{e}$ | - | - | 0.0000 | 0.0000 |
| $2_{f}$ | b | 0.0000 |  | 0.5961(8) |
| $2_{e}$ | d | - | 0.0000 | 0.5961(4) |
| $3_{f}$ | a | 1.4908(6) | - | 1.4908(6) |
| $3_{e}$ | c | - | 1.4908 (6) | 1.4908(5) |
| $4_{f}$ | b | 2.0870(6) | - | $2.6830(10)$ |
| 4 e | d | - | 2.0870(6) | 2.6831(7) |
| $5{ }_{f}$ | a | 4.1742(8) |  | 4.1742 (8) |
| 5 | c | - | 4.1742(8) | $4.1742(10)$ |
| $6_{f}$ | b | 5.3658(8) | - | $5.9619(11)$ |
| 6 e | d | - | 5.3670 (8) | 5.9631(9) |
| $7_{f}$ | a | 8.0492(10) | - | 8.0492(10) |
| $7{ }_{e}$ | c | - | 8.0504(10) | 8.0504(11) |
| $8_{f}$ | b | $9.8373(10)$ | - | 10.4335(12) |
| 8 e | d | - | $9.8402(10)$ | 10.4363(10) |
| $9_{f}$ | a | 13.1162(11) | - | 13.1162(11) |
| $9_{e}$ | c | - | 13.1210(11) | 13.1210(12) |
| $10_{f}$ | b | 15.5000(11) | - | 16.0980(14) |
| $10_{e}$ | d | - | 15.5066(11) | 16.1027(12) |
| $11_{f}$ | a | 19.3745(13) | - | 19.3745(13) |
| $11_{e}$ | c | - | 19.3836(13) | 19.3836(14) |
| $12_{f}$ | b | 22.3537(13) | - | 22.9498(15) |
| $12{ }_{e}$ | d | - | 22.3653(13) | 22.9614(13) |
| $13_{f}$ | a | 26.8246(14) | - | 26.8246(14) |
| $13_{e}$ | c | - | 26.8391(14) | 26.8391(15) |
| $14_{f}$ | b | 30.4000(14) | - | 30.9960(16) |
| $14_{e}$ | d | - | 30.4173(14) | 31.0134(14) |
| $\mathbf{B}_{a}$ | MHz] | 4467.66(17) | $\mathrm{B}_{b}$ [MHz] | 4467.46(15) |
| $\mathbf{B}_{c}$ | MHz] | 4470.22(9) | $\mathrm{B}_{d}$ [MHz] | 4470.09(5) |
| $\mathbf{B}_{\text {combined }}$ |  | [ ${ }_{\text {[ }} \mathrm{MHz}$ ] | 4468.82(34) |  |
| $\alpha_{\text {combined }}^{a}$ |  | ${ }_{d}$ [MHz] | -7.88(34) |  |
| $\alpha_{\nu_{6}=1, \text { calc }}^{b}$ |  | [ MHz ] | -8.00 |  |
| $\alpha_{\nu_{7}=1, \text { calc }}^{b}$ [ MHz ] |  |  | -18.93 |  |

Table S5: Term values for the ground state of the second hot-band (Hot 2) of $\mathrm{HCCCO}^{+}$as derived from ground state combination differences. Four separate networks of states are derived from the experimental data, denoted as a,b,c and d, due to the two l-type components of the underlying bending motion. This leads to four reference points $E-D_{0}=0$ that are not directly linked to each other. These sets are combined based on their y-axis intercept assuming the validity of the rigid rotor approximation.

| $\mathbf{J}_{e, f}$ | Set | E-D ${ }_{0}^{a, c}\left[\mathrm{~cm}^{-1}\right]$ | E-D ${ }_{0}^{\text {b,d }}\left[\mathrm{cm}^{-1}\right]$ | $\mathbf{E}_{\text {combined }}\left[\mathrm{cm}^{-1}\right]$ |
| :---: | :---: | :---: | :---: | :---: |
| $1_{e}$ | - | $\cdots$ | - | 0.000000 |
| $1_{f}$ | - | - | - | 0.000000 |
| $2{ }_{e}$ | c | 0.000000 | - | 0.59733(1) |
| $2_{f}$ | d | - | 0.000000 | 0.59733(1) |
| $3_{e}$ | a | 0.000000 | - | 1.49336(3) |
| $3_{f}$ | b | - | 0.000000 | $1.49336(3)$ |
| $4_{e}$ | c | 2.09043(57) | - | 2.68776(57) |
| $4_{f}$ | d |  | 2.09043(57) | 2.68776(57) |
| $5{ }_{\text {e }}$ | a | 2.68793(57) | - | 4.18129(57) |
| $5{ }_{f}$ | b |  | 2.68793(57) | 4.18129 (57) |
| 6 e | c | 5.37611(80) | - | $5.97344(80)$ |
| $6_{f}$ | d |  | $5.37611(80)$ | $5.97344(80)$ |
| $7_{e}$ | a | 6.57075(80) | - | 8.06410(80) |
| $7_{f}$ | b | - | 6.57075(80) | 8.06410(80) |
| 8 e | c | 9.85601(98) | - | $10.45334(98)$ |
| $8_{f}$ | d | - | 9.85601(98) | $10.45334(98)$ |
| $9_{e}$ | a | 11.64856(98) | - | 13.14191(98) |
| $9_{f}$ | b | - | 11.64856(98) | 13.14191(98) |
| $10_{e}$ | c | 15.53081(113) | - | 16.12814(113) |
| $10_{f}$ | d | - | 15.53081(113) | 16.12814(113) |
| $11_{e}$ | a | 17.92007(113) | - | 19.41343(113) |
| $11_{f}$ | b | - | 17.92007(113) | 19.41343(113) |
| $12_{e}$ | c | 22.39959(127) | - | 22.99692(127) |
| $12_{f}$ | d | - | 22.39959(127) | 22.99692(127) |
| $\mathrm{B}_{a}$ [ | MHz] | 4476.97(8) | $\mathrm{B}_{b}$ [MHz] | 4476.97(8) |
| $\mathrm{B}_{c}$ | MHz] | 4476.86(5) | $\mathrm{B}_{d}$ [MHz] | 4476.86(5) |
| $\mathbf{B}_{\text {combined }}$ |  | ed [MHz] | 4476.90(3) |  |
| $\alpha_{\text {combined }}^{a}$ |  | [ ${ }^{\text {d }}$ [ ${ }^{\text {d }}$ | -15.96(5) |  |
| $\alpha_{\nu_{7}=1, \text { calc }}^{b}$ |  | ${ }_{l c}$ [MHz] | -18.93 |  |
| $\alpha_{\nu_{7}=2, \text { calc }}^{b}$ [ MHz$]$ |  |  | -37.86 |  |

Table S6: Term values of the excited states of the observed hot-bands of $\mathrm{HCCCO}^{+}$. The given energies are calculated from the observed transitions utilizing the combined ground state term values given in Tables S 4 and S 5 . All units are given in wavenumbers unless otherwise noted.


## References

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[^0]:    ${ }^{1} \mathrm{fc}$-CCSD $(\mathrm{T}) /$ aug-cc-pV $(\mathrm{T}+\mathrm{d}) \mathrm{Z}$ level of theory
    ${ }^{2} \mathrm{fc}$ - $\operatorname{CCSD}(\mathrm{T}) /$ ANO1 level of theory

