

The variation in the space-charge potential at grain boundaries in SrTiO₃: the effects of two restricted equilibria

— Supplemental Information —

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S1 Explanations of the temperature dependence of Φ_0

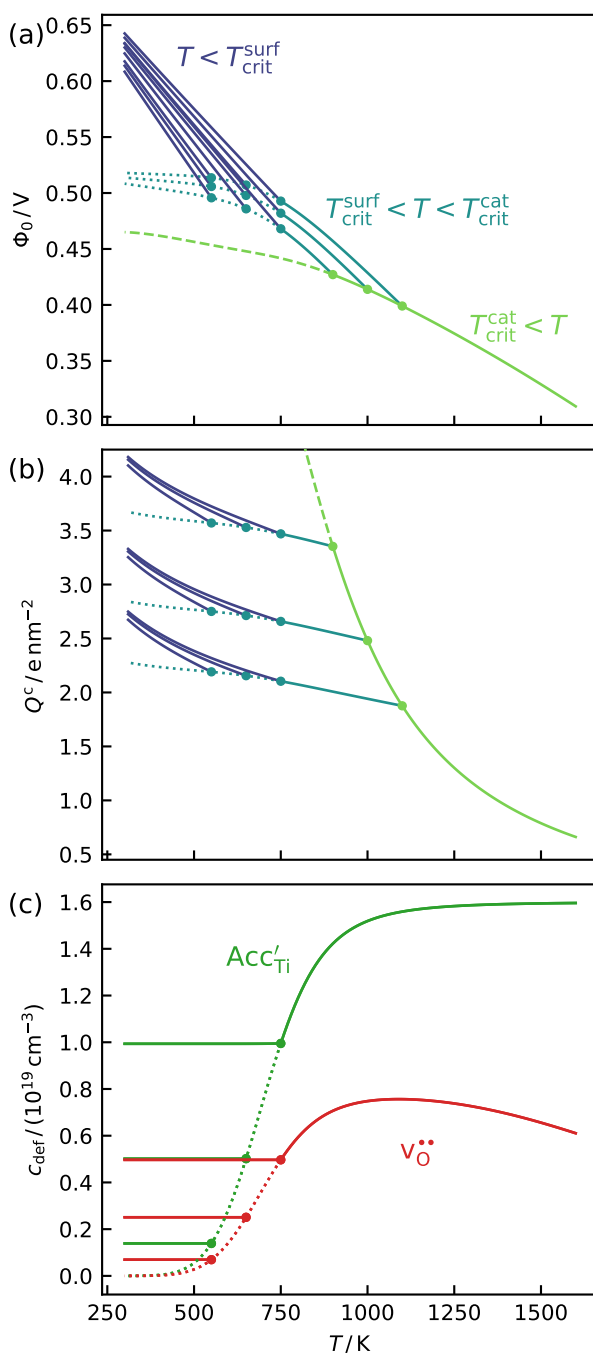


Figure S1: Temperature dependence of (a) the space-charge potential, Φ_0 ; (b) the grain-boundary core charge, Q^c ; and (c) the bulk defect concentrations of Acc'_{Ti} and $v_{\text{O}}^{\bullet\bullet}$, with the temperature range encompassing the full-equilibrium (FE) as well as one restricted-equilibrium (1RE) and two restricted-equilibria (2RE) scenarios.

$\Phi_0(T)$ steepens both at $T_{\text{crit}}^{\text{cat}}$ and $T_{\text{crit}}^{\text{surf}}$, but the reasons in both cases are quite different. The steepening below $T_{\text{crit}}^{\text{surf}}$ may be attributed to a steeper increase in the grain-boundary core charge, $Q^c(T)$. The abrupt change in the core charge's behaviour may be explained with $c^b(v_{\text{O}}^{\bullet\bullet})$ abruptly staying constant below $T_{\text{crit}}^{\text{cat}}$, at a value that is much higher than the chemical equilibrium concentrations for $T < T_{\text{crit}}^{\text{cat}}$. Larger Q^c generally have an enhancing effect on Φ_0 . Consistently, the Φ_0 values at a specific temperature are higher in the 2RE regime than in the 1RE regime. The abrupt steepening in $\Phi_0(T)$ at $T_{\text{crit}}^{\text{cat}}$, in contrast, stands opposed to an abrupt flattening in $Q^c(T)$. Increases in Q^c can be accommodated much more effectively in the FE regime than in the RE regimes, where an enrichment in Acc'_{Ti} is possible only by a valence change. Owing to this weaker shielding capability, the space-charge layer reacts, in the 1RE regime, much more sensitively to increases in Q^c than it does in the FE regime, and hence, Φ_0 increases more steeply despite the flatter increase in Q^c .

S2 Different grain-boundary parameters

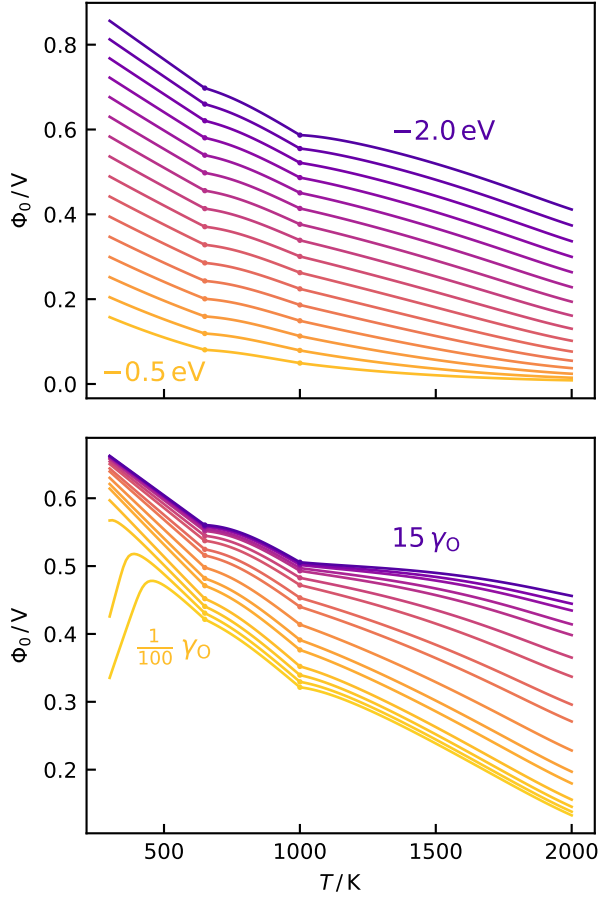


Figure S2: Temperature dependence of the space-charge potential, Φ_0 , for different grain-boundary parameters: (a) different values of the v_{O}^{\bullet} segregation energy, $\Delta\mu_v^{\ominus}$ (with $\gamma_{\text{O}}^c = \gamma_{\text{O}}/20$); (b) different values of the core oxygen-site density, γ_{O}^c (with $\Delta\mu_v^{\ominus} = -1.5$ eV). In all cases, we set $T_{\text{crit}}^{\text{cat}} = 1000$ K and $T_{\text{crit}}^{\text{surf}} = 650$ K.

Clearly, the non-trivial temperature dependence of Φ_0 with kinks at the critical temperatures applies over a wide range of grain-boundary core properties, i.e., the differences that are due to different temperature ranges must be expected to be encountered across different grain-boundary types.