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Supporting information

for

Seven-coordinate Mn(II) complex with a pyridine-based 15-membered macrocyclic ligand containing one acetate pendant arm: structure, stability and relaxation properties

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Figure S1. ESI mass spectrum of studied ligand L1 (positive mode).



Figure S2. ¹H NMR spectrum (400 MHz, dmso- d_6 ; 60 °C) of studied ligand L1.



Figure S3. ¹³C NMR spectrum (400 MHz, dmso-*d*₆) of studied ligand L1.



Figure S4. ¹H–¹³C HMQC NMR spectrum (400 MHz, dmso-*d*₆) of studied ligand L1



Figure S5. ¹H–¹³C HMBC NMR spectrum (400 MHz, dmso-*d*₆) of studied ligand L1



Figure S6. The crystal packing of complex **1**. View of the asymmetric unit (*top*) and view of the supramolecular 1D chains along a axis (*bottom*). Blue dashed lines represent hydrogen bonds. Hydrogen atoms were omitted for clarity.



Figure S7. Potentiometric titration curves, as a dependence of pH on base equivalents, obtained for ligand L1 in the absence and in the presence of selected divalent metal ions (I = 0.15M; 25 °C).



Figure S8. Dependence of the observed dissociation rate constants for **MnL1** on Zn^{2+} concentration. pH readings from the top are 4.9 (*black*), 5.0 (*red*), 5.2 (*green*), 5.5 (*dark blue*), 5.7 (*light blue*) and 6.1 (*pink*).

Table S1 Selected non-covalent contact parameters (A, °) for complexes 1.						
D	Η	Α	d (D – H)	d(H···A)	d(D···A)	<(D-H-A)
05	H5A	03	0.87	1.91	2.771(4)	169.9
05	H5B	$Cl1^1$	0.87	2.33	3.194(3)	171.9
06	H6C	05	0.87	2.21	2.918(3)	138.9
06	H6D	$O5^2$	0.87	2.22	2.918(3)	136.8
N3	H3	O5 ³	1.00	2.06	2.991(4)	153.7

Table S1 Selected non-covalent contact parameters (Å, °) for complexes 1.

Symmetry transformations used to generate equivalent atoms: ¹1/2-X,-1/2+Y,-1/2+Z; ²+X,1/2-Y,+Z; ³1/2+X,+Y,3/2-Z

Analysis of the ¹⁷O NMR and ¹H NMRD data¹

The reduced transverse ¹⁷O relaxation rates, $1/T_{2r}$, and reduced ¹⁷O chemical shifts, ω_r , are calculated from the measured relaxation rates $1/T_2$ and angular frequencies ω of the paramagnetic solutions and from the relaxation rates $1/T_{2A}$ and angular frequencies ω_A of the diamagnetic reference according to the Equations (S1) and (S2)²:

$$\frac{1}{T_{2r}} = \frac{1}{P_{m}} \left[\frac{1}{T_{2}} - \frac{1}{T_{2A}} \right] = \frac{1}{\tau_{m}} \frac{T_{2m}^{-2} + \tau_{m}^{-1} T_{2m}^{-1} + \Delta \omega_{m}^{2}}{(\tau_{m}^{-1} + T_{2m}^{-1})^{2} + \Delta \omega_{m}^{2}}$$
(S1)

$$\Delta \omega_{\rm r} = \frac{1}{P_{\rm m}} (\omega - \omega_{\rm A}) = \frac{\Delta \omega_{\rm m}}{(1 + \tau_{\rm m} T_{\rm 2m}^{-1})^2 + \tau_{\rm m}^2 \Delta \omega_{\rm m}^2}$$
(S2)

 $\Delta \omega_{\rm m}$ is determined by the hyperfine or scalar coupling constant, $A_{\rm O}/\hbar$, where *B* represents the magnetic field, *S* is the electron spin and $g_{\rm L}$ is the isotropic Landé *g* factor (Equation (S3)).

$$\Delta \omega_{\rm m} = \frac{g_L \mu_{\rm B} S(S+1) B}{3k_{\rm B} T} \frac{A_o}{\hbar}$$
(S3)

The chemical shifts are measured with high errors (given the large linewidths) and the reduced chemical shifts were not further used in the fit.

The ¹⁷O transverse relaxation rate is mainly determined by the scalar contribution, $1/T_{2sc}$, and it is given by Equation (S4).

$$\frac{1}{T_{2m}} \approx \frac{1}{T_{2sc}} = \frac{S(S+1)}{3} \left(\frac{A_o}{\hbar}\right)^2 \left(\tau_{s1} + \frac{\tau_{s2}}{1 + \tau_{s2}^2 \omega_s^2}\right) \qquad \qquad \frac{1}{\tau_{si}} = \frac{1}{\tau_m} + \frac{1}{T_{ie}}$$
(S4)

The exchange rate, k_{ex} , (or inverse binding time, τ_{m}) of the inner sphere water molecule is assumed to obey the Eyring equation (Equation (S5)) where ΔS^{\ddagger} and ΔH^{\ddagger} are the entropy and enthalpy of activation for the exchange, and k_{ex}^{298} is the exchange rate at 298.15 K.

$$\frac{1}{\tau_{\rm m}} = k_{\rm ex} = \frac{k_{\rm B}T}{h} \exp\left\{\frac{\Delta S^{\ddagger}}{R} - \frac{\Delta H^{\ddagger}}{RT}\right\} = \frac{k_{\rm ex}^{298}T}{298.15} \exp\left\{\frac{\Delta H^{\ddagger}}{R} \left(\frac{1}{298.15} - \frac{1}{T}\right)\right\}$$
(S5)

The electronic relaxation is mainly governed by modulation of the transient zero-field splitting, and for the electron spin relaxation rates, $1/T_{1e}$ and $1/T_{2e}$, McMachlan has developed Equations (S6)–(S8):

$$\left(\frac{1}{T_{1e}}\right) = \frac{32}{25} \Delta^2 \left(\frac{\tau_v}{1 + \omega_s^2 \tau_v^2} + \frac{4\tau_v}{1 + 4\omega_s^2 \tau_v^2}\right)$$
(S6)
$$\left(\frac{1}{T_{2e}}\right) = \frac{32}{50} \Delta^2 \left[3\tau_v + \frac{5\tau_v}{1 + \omega_s^2 \tau_v^2} + \frac{2\tau_v}{1 + 4\omega_s^2 \tau_v^2}\right]$$
(S7)

$$\tau_{v} = \tau_{v}^{298} \exp\left\{\frac{E_{v}}{R} \left(\frac{1}{T} - \frac{1}{298.15}\right)\right\}$$
(S8)

where Δ^2 is the trace of the square of the transient zero-field-splitting (ZFS) tensor, τ_v is the correlation time for the modulation of the ZFS with the activation energy E_v , and ω_s is the Larmor frequency of the electron spin.

The proton relaxivities (normalized to 1 mM Mn^{2+} concentration) originate from inner- and outer-sphere contributions (Equation (S9)):

$$r_1 = r_{1is} + r_{1os} (S9)$$

The inner-sphere term is given by Equation (S10), where q is the number of inner-sphere water molecules.

$$r_{\rm lis} = \frac{1}{1000} \times \frac{q}{55.55} \times \frac{1}{T_{\rm lm}^{\rm H} + \tau_{\rm m}}$$
(S10)

In the longitudinal relaxation rate of inner sphere water protons, $1/T_{1m}^{H}$, the dipolar contribution dominates (Equation (S11)):

$$\frac{1}{T_{1m}^{H}} \approx \frac{1}{T_{1}^{DD}} = \frac{2}{15} \left(\frac{\mu_{0}}{4\pi}\right)^{2} \frac{\hbar^{2} \gamma_{S}^{2} \gamma_{I}^{2}}{r_{MnH}^{6}} S(S+1) \left[\frac{3\tau_{d1H}}{1+\omega_{I}^{2} \tau_{d1H}^{2}} + \frac{7\tau_{d2H}}{1+\omega_{S}^{2} \tau_{d2H}^{2}}\right]$$
(S11)

Here r_{MnH} is the effective distance between the Mn^{2+} electron spin and the water protons, ω_{I} is the proton resonance frequency, τ_{diH} is given by Eq. S12, where τ_{RH} is the rotational correlation time of the Mn^{2+} –H_{water} vector:

$$\frac{1}{\tau_{\rm diH}} = \frac{1}{\tau_{\rm m}} + \frac{1}{\tau_{\rm RH}} + \frac{1}{T_{i\rm e}} \qquad \qquad i = 1, 2; \tag{S12}$$

No scalar relaxation term was included in the analysis, as the NMRD profiles did not present a second, low-field dispersion, typical of an important scalar contribution to relaxivity.¹

The outer-sphere contribution to the overall relaxivity is described by Equation (S13), where N_A is the Avogadro constant, and J_{os} is a spectral density function (Equation (S14)).

$$r_{1\text{os}} = \frac{32N_{\text{A}}\pi}{405} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\hbar^2 \gamma_{\text{S}}^2 \gamma_{1}^2}{a_{\text{MnH}} D_{\text{MnH}}} S(S+1) [3J_{\text{os}}(\omega_1, T_{1\text{e}}) + 7J_{\text{os}}(\omega_{\text{S}}, T_{2\text{e}})]$$
(S13)

$$J_{os}(\omega, T_{je}) = \operatorname{Re}\left[\frac{1 + \frac{1}{4}\left(i\omega\tau_{MnH} + \frac{\tau_{MnH}}{T_{je}}\right)^{1/2}}{1 + \left(i\omega\tau_{MnH} + \frac{\tau_{MnH}}{T_{je}}\right)^{1/2} + \frac{4}{9}\left(i\omega\tau_{MnH} + \frac{\tau_{MnH}}{T_{je}}\right) + \frac{1}{9}\left(i\omega\tau_{MnH} + \frac{\tau_{MnH}}{T_{je}}\right)^{3/2}}\right]$$

$$j = 1, 2$$
(S14)

The diffusion coefficient for the diffusion of a water proton away from a Mn^{2+} complex, D_{MnH} , obeys the exponential temperature dependence described by Equation (S15), with activation energy E_{MnH} :

$$D_{\rm MnH} = D_{\rm MnH}^{298} \exp\left\{\frac{E_{MnH}}{R} \left(\frac{1}{298.15} - \frac{1}{T}\right)\right\}$$
(S15)

In the fitting procedure, some parameters have been fixed to common values: the distance between the metal ion and the inner and outer sphere water protons ($r_{MnH} = 2.83$ Å and $a_{MnH} =$ 3.6 Å, respectively); the diffusion coefficient and its activation energy ($D_{MnH} = 26 \times 10^{-10} \text{ m}^2 \text{s}^{-1}$ and $E_{MnH} = 20 \text{ kJ mol}^{-1}$).

¹ E. Balogh, Z. He, W. Hsieh, S. Liu, É. Tóth; *Inorg. Chem.* 2007, **46**, 238-250. ² T. J. Swift, R. E. Connick, *J. Chem. Phys.* 1962, **37**, *2*, 307-320.