

Supporting Information

Adaptable sublattice stabilized high-entropy materials with superior thermoelectric performance

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A: Lattice thermal conductivity calculations

The lattice thermal conductivity is calculated by subtracting the carrier contribution κ_e and bipolar contribution κ_{bipolar} from the total κ , i.e. $\kappa_L = \kappa - \kappa_e - \kappa_{\text{bipolar}}$, where κ_L is lattice thermal conductivity, κ is measured total thermal conductivity, κ_e is electrical thermal conductivity, and κ_{bipolar} is bipolar thermal conductivity. The carrier thermal conductivity can be determined by the Wiedemann-Franz law: $\kappa_e = L\sigma T$, where T is absolute temperature, σ is experimental electrical conductivity and L is the Lorenz factor. The simple two-band model proposed by Simon¹, i.e. one parabolic conduction band and one parabolic valence band, is used to calculate the bipolar contribution κ_{bipolar} . The Lorenz factor L can be estimated by Equations S1–S2¹:

$$S = m \frac{k_B}{e} \left(\frac{(\lambda + 2)F_{\lambda+1}(\eta)}{(\lambda + 1)F_{\lambda}(\eta)} - \eta \right) \quad (\text{S1})$$

$$L = \left(\frac{k_B}{e} \right)^2 \left\{ \frac{(\lambda + 3)F_{\lambda+2}(\eta)}{(\lambda + 1)F_{\lambda}(\eta)} - \left[\frac{(\lambda + 2)F_{\lambda+1}(\eta)}{(\lambda + 1)F_{\lambda}(\eta)} \right]^2 \right\} \quad (\text{S2})$$

where e is elementary charge, k_B is Boltzmann constant, S is experimental Seebeck coefficient, η is reduced Fermi level, λ is scattering parameter, and $F_{\lambda}(\eta)$ is Fermi

integrate defined as $F_{\lambda}(\eta) = \int_0^{\infty} \frac{x^{\lambda}}{1 + e^{x-\eta}} dx$.

The room temperature reduced Fermi level is derived from both the carrier

concentration and Seebeck coefficient on the basis of single band approximation. Assuming that the intrinsic carrier n_0 is negligible and external doping is fully excitation, external dopant N_D is equal to be experimental carrier concentration at room temperature. Therefore, the electron concentration, hole concentration and reduced Fermi level can be solved by Equations S3–S5¹:

$$n = p_0 + N_D \quad (\text{S3})$$

$$n = 4\pi \left(\frac{2m_e^* k_B T}{h^2} \right)^{3/2} F_\lambda(\eta) \quad (\text{S4})$$

$$p = 4\pi \left(\frac{2m_h^* k_B T}{h^2} \right)^{3/2} F_\lambda(-\eta - E_g) \quad (\text{S5})$$

where m_e^* is electron effective mass, m_h^* is hole effective mass and E_g is band gap.

The electron Seebeck coefficient S_n and hole Seebeck coefficient S_p can be calculated by:

$$S_n = m \frac{k_B}{e} \left(\frac{(\lambda + 2)F_{\lambda+1}(\eta)}{(\lambda + 1)F_\lambda(\eta)} - \eta \right) \quad (\text{S6})$$

$$S_p = m \frac{k_B}{e} \left(\frac{(\lambda + 2)F_{\lambda+1}(-\eta - E_g)}{(\lambda + 1)F_\lambda(-\eta - E_g)} + \eta + E_g \right) \quad (\text{S7})$$

The electron conductivity σ_n and hole conductivity σ_p can be estimated by Equations S8–S10:

$$\sigma_n = n_0 \mu_{\text{aver}} e \quad (\text{S8})$$

$$\sigma_h = A p_0 \mu_{\text{aver}} e \quad (\text{S9})$$

$$\sigma_n S_n + \sigma_p S_p = \sigma S \quad (\text{S10})$$

where μ_{aver} is average mobility, and A is a parameter to fit calculated Seebeck coefficient and measured Seebeck coefficient. Finally, bipolar thermal conductivity can be calculated by¹:

$$\kappa_{\text{bipolar}} = \frac{\sigma_n \sigma_p}{\sigma_n + \sigma_p} (S_n - S_p)^2 T \quad (\text{S11})$$

B: Calculation of the Mg content

It is assumed that each Bi_{Sn} introduces one electron and each V_{Mg} introduces two holes. Then, the content of Mg vacancies can be estimated by²:

$$x_v = \frac{1}{2} \left(x - \Delta n \frac{M}{\rho N_A} \right) \quad (\text{S12})$$

$$\Delta n = n - n_0 \quad (\text{S13})$$

where x_v is the content of V_{Mg} , x is the nominal Bi content, n is the measured carrier concentration for our $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.75-x}\text{Bi}_x$ samples, n_0 is the carrier concentration of $\text{Mg}_2\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.75}$, M is the relative molecular mass, ρ is the sample's density, and N_A is the Avogadro constant. Therefore, the content of Mg (x_{Mg}) is given by

$$x_{\text{Mg}} = 2 - x_v \quad (\text{S14})$$

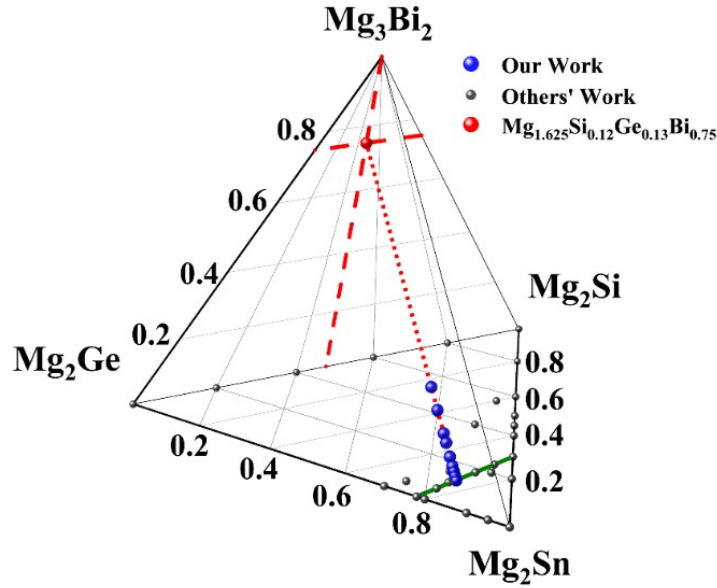


Fig. S1. Phase diagram of Mg_2Si - Mg_2Ge - Mg_2Sn - Mg_3Bi_2 quaternary system. The black sphere symbols represent the data from references³⁻¹⁰, while the blue sphere symbols represent the samples synthesized in this work.

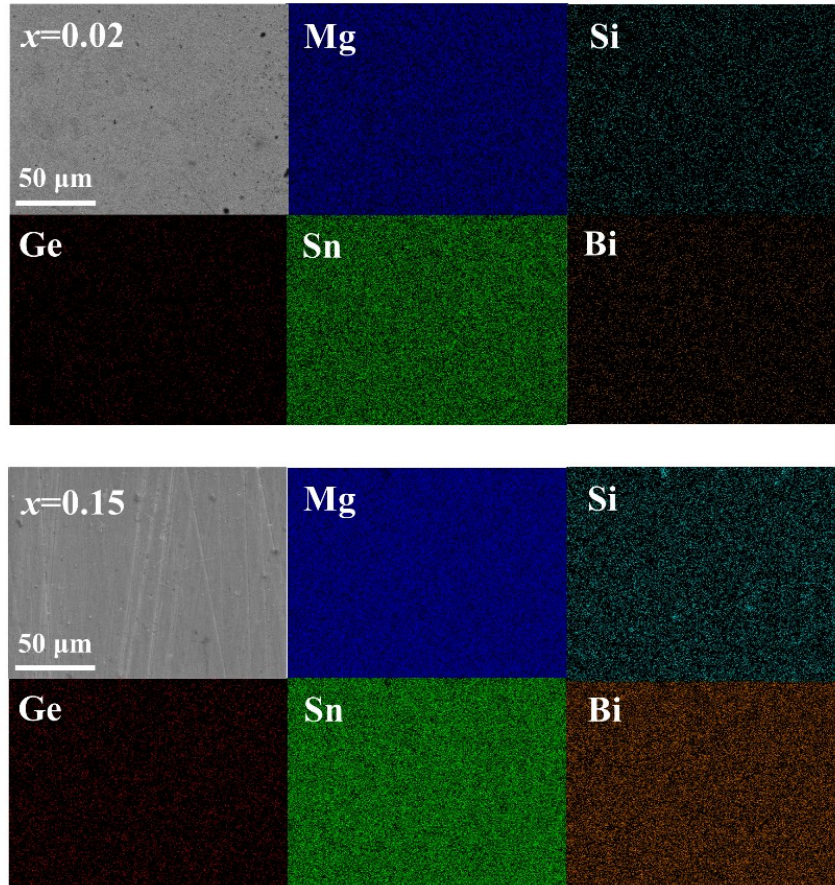


Fig. S2. Backscattered electron microscopy (BSE) images and corresponding elemental energy-dispersive X-ray spectroscopy (EDS) mapping of $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.75-x}\text{Bi}_x$ ($x = 0.02, 0.15$)

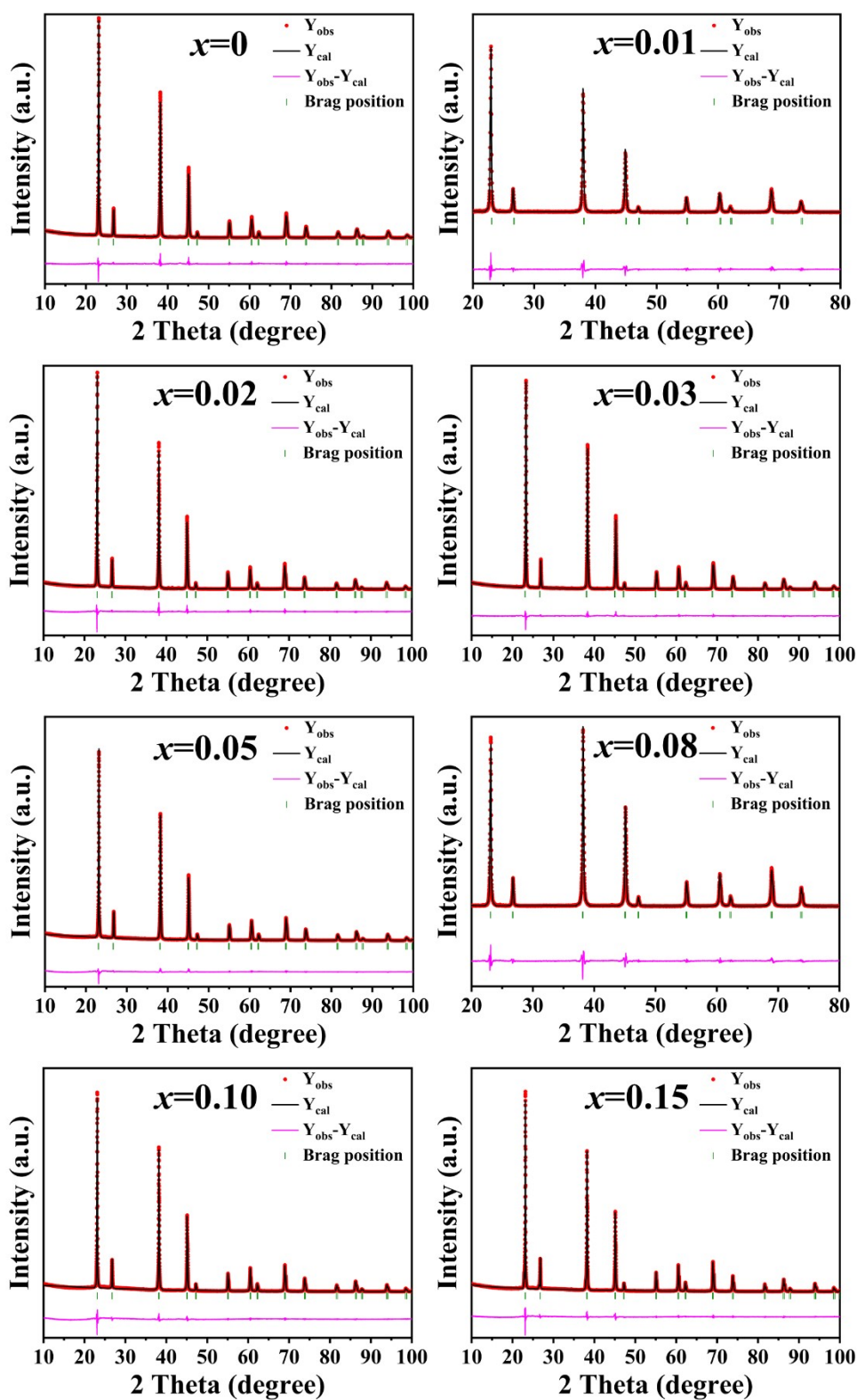


Fig. S3. Observed (red) and calculated (black) diffraction patterns and difference profiles (pink) for samples $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.75-x}\text{Bi}_x$ ($x=0, 0.01, 0.02, 0.03, 0.05, 0.08, 0.10, 0.15$) at room temperature.

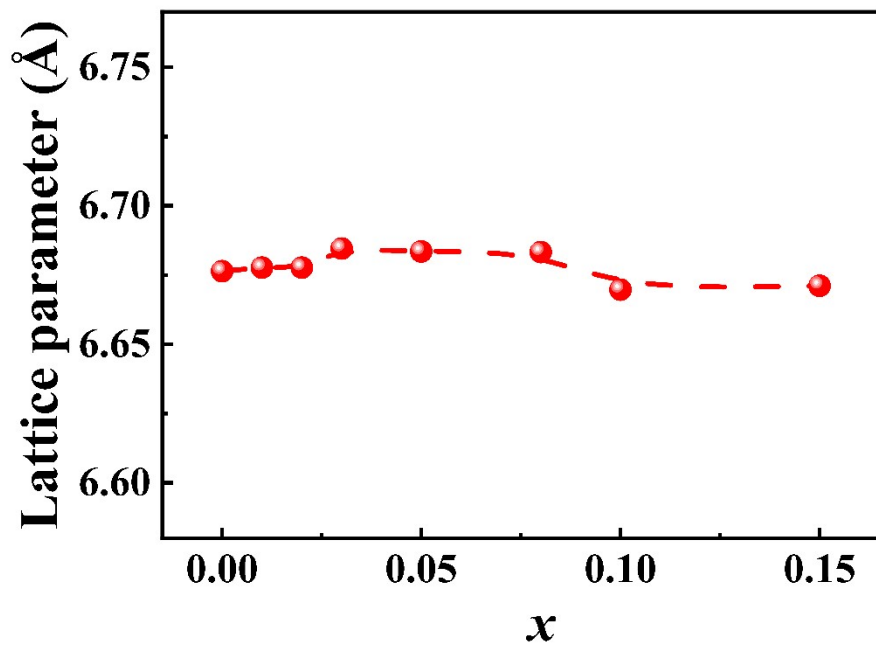


Fig. S4. Lattice parameter as a function of Bi content x .

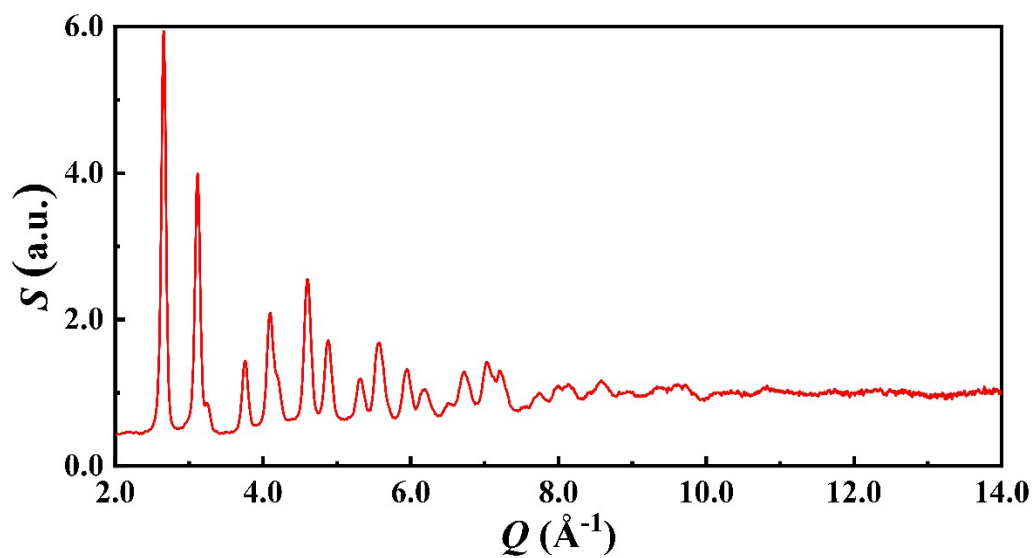


Fig. S5. Structure factor S as a function of Q .

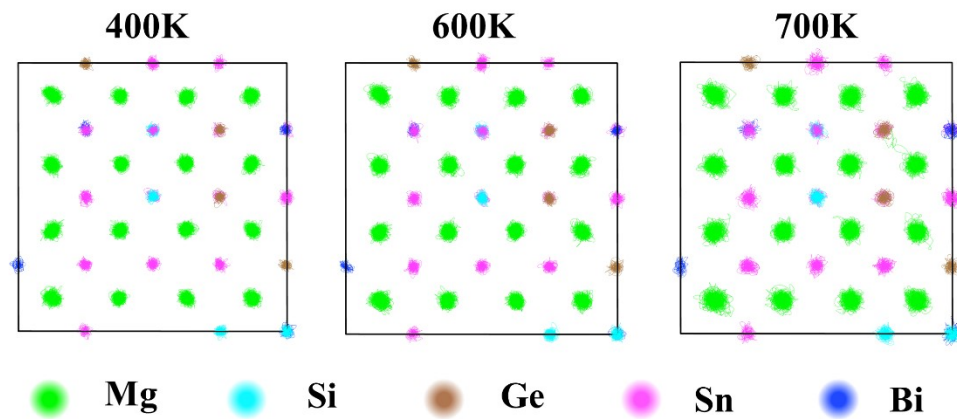


Fig. S6. Molecular dynamics simulations of the trajectory of Mg, Si, Ge, Sn, Bi at 400 K, 600 K, and 700 K

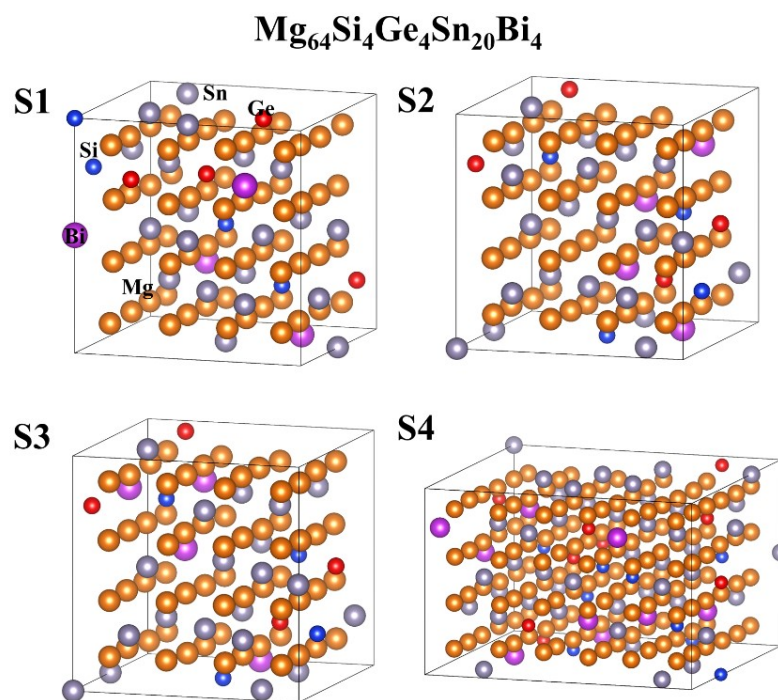


Fig. S7. Crystal structure of different supercells. (a) S1 ($2 \times 2 \times 2$ supercell), (b) S2 ($2 \times 2 \times 2$ supercell), (c) S3 ($2 \times 2 \times 2$ supercell), and (d) S4 ($3 \times 3 \times 2$ supercell).

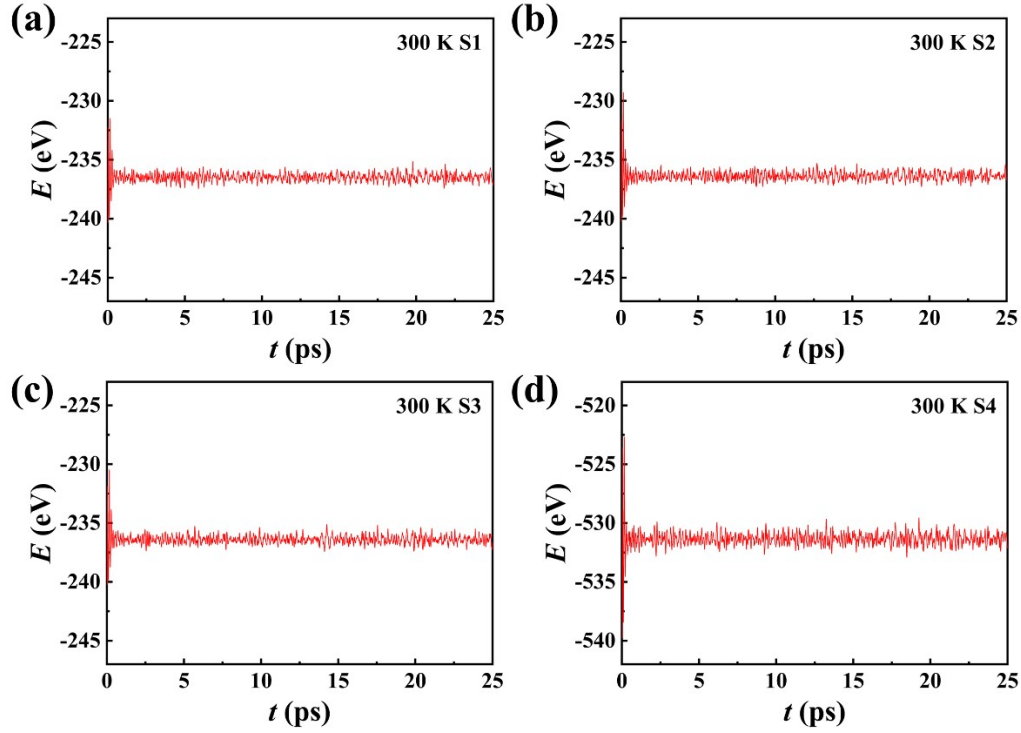


Fig. S8. Total energy E as a function of stimulation time t for different supercells. (a) S1 ($2 \times 2 \times 2$ supercell), (b) S2 ($2 \times 2 \times 2$ supercell), (c) S3 ($2 \times 2 \times 2$ supercell), and (d) S4 ($3 \times 3 \times 2$ supercell).

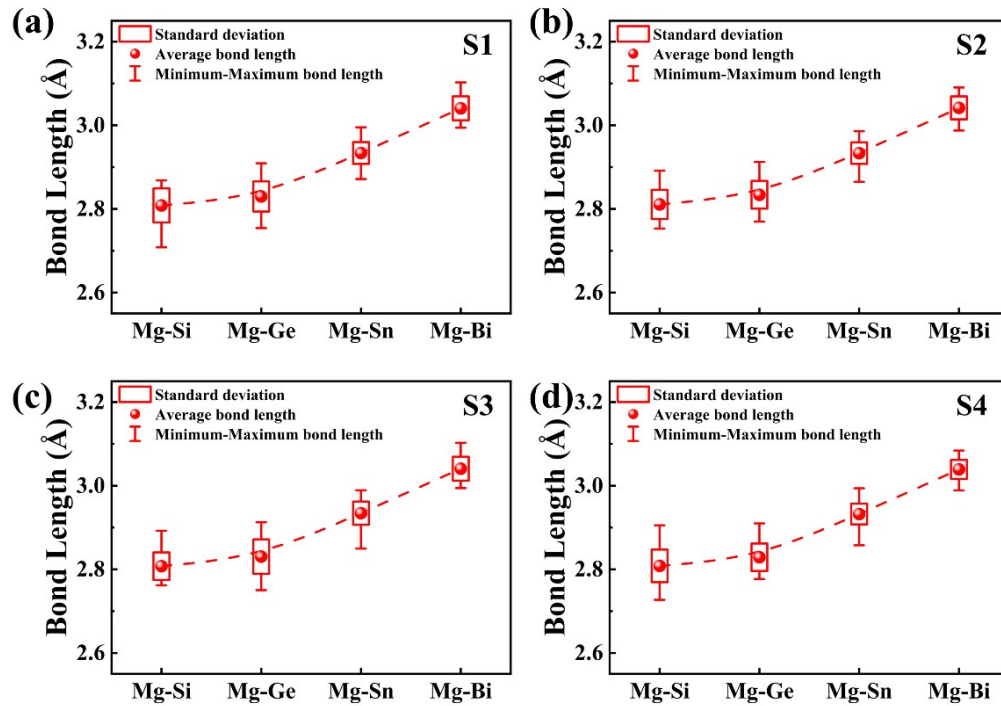


Fig. S9. Calculated bond length for different supercells. (a) S1 ($2 \times 2 \times 2$ supercell), (b) S2 ($2 \times 2 \times 2$ supercell), (c) S3 ($2 \times 2 \times 2$ supercell), and (d) S4 ($3 \times 3 \times 2$ supercell).

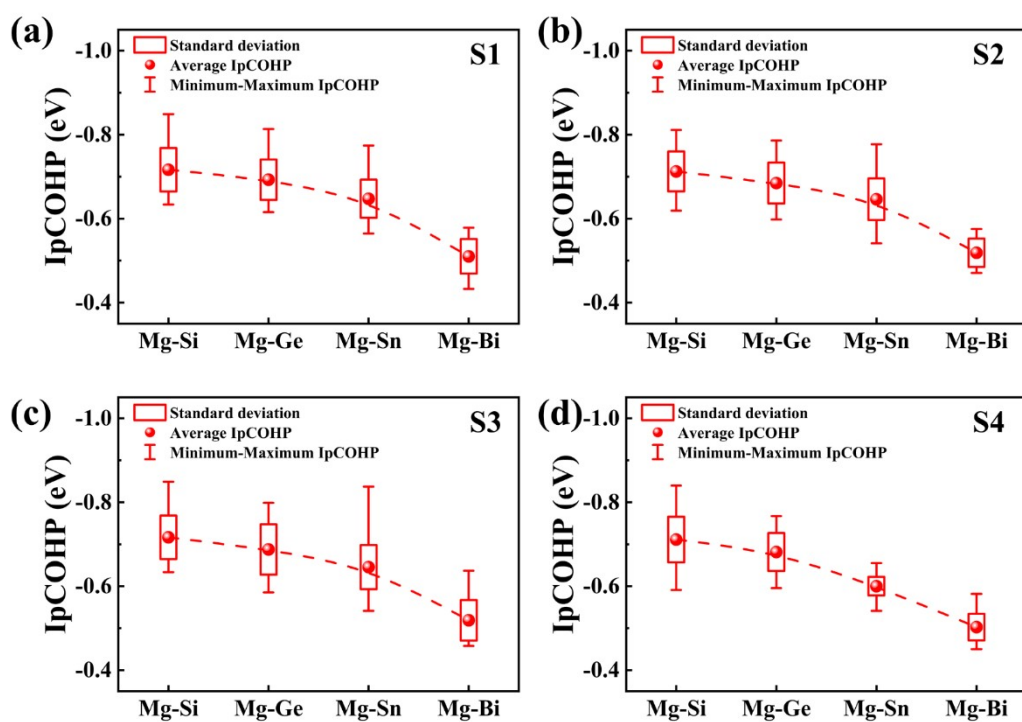


Fig. S10. Calculated $I_p\text{COHP}$ for different supercells. (a) S1 ($2\times 2\times 2$ supercell), (b) S2 ($2\times 2\times 2$ supercell), (c) S3 ($2\times 2\times 2$ supercell), and (d) S4 ($3\times 3\times 2$ supercell).

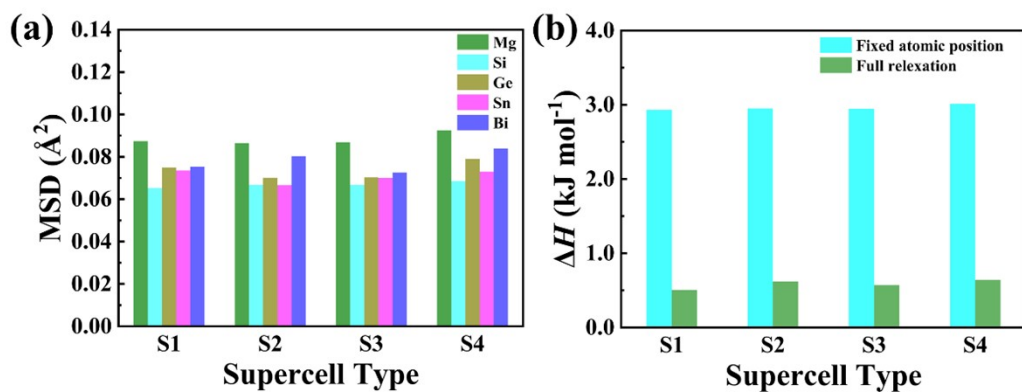


Fig. S11. Calculated (a) MSD at 300 K and (b) mixing enthalpy ΔH for different supercells.

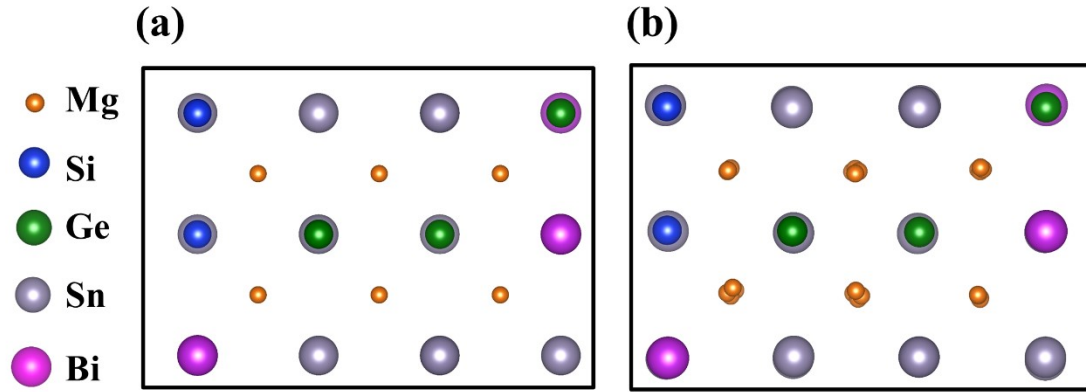


Fig. S12. The obtained crystal structures of $\text{Mg}_2\text{Si}_4\text{Ge}_4\text{Sn}_{20}\text{Bi}_4$: (a) volume relaxation with fixed atomic positions within the cubic lattice, (b) full relaxation with the cell volume, shape and atomic positions optimized.

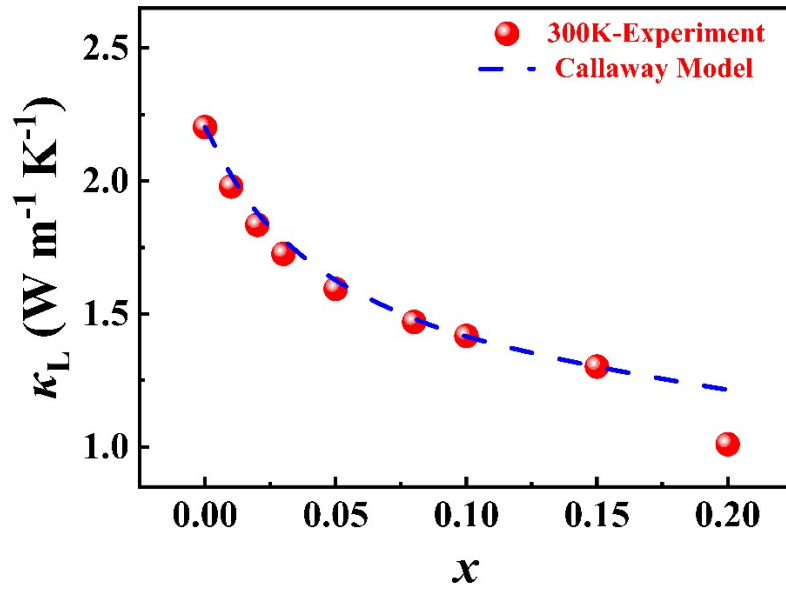


Fig. S13. Lattice thermal conductivity κ_L at 300 K as a function of Bi content x . The dashed line denotes the calculated κ_L by the Callaway model.

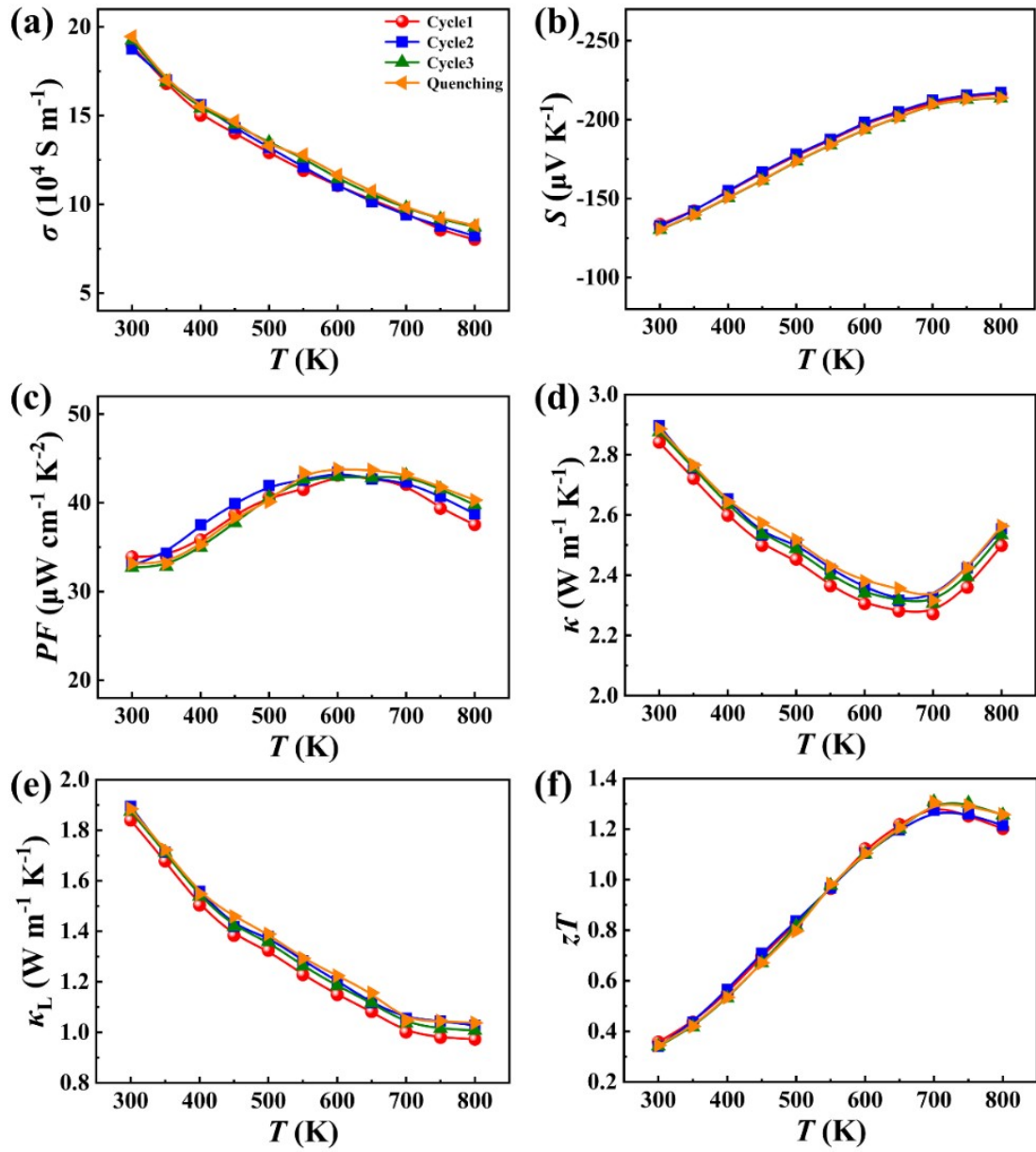


Fig. S14. Temperature dependence of (a) electrical conductivity σ , (b) Seebeck coefficient S , (c) power factor PF , (d) thermal conductivity κ , (e) lattice thermal conductivity κ_L and (d) zT of sample $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.73}\text{Bi}_{0.02}$ after three cycles and after quenching at 800 K.

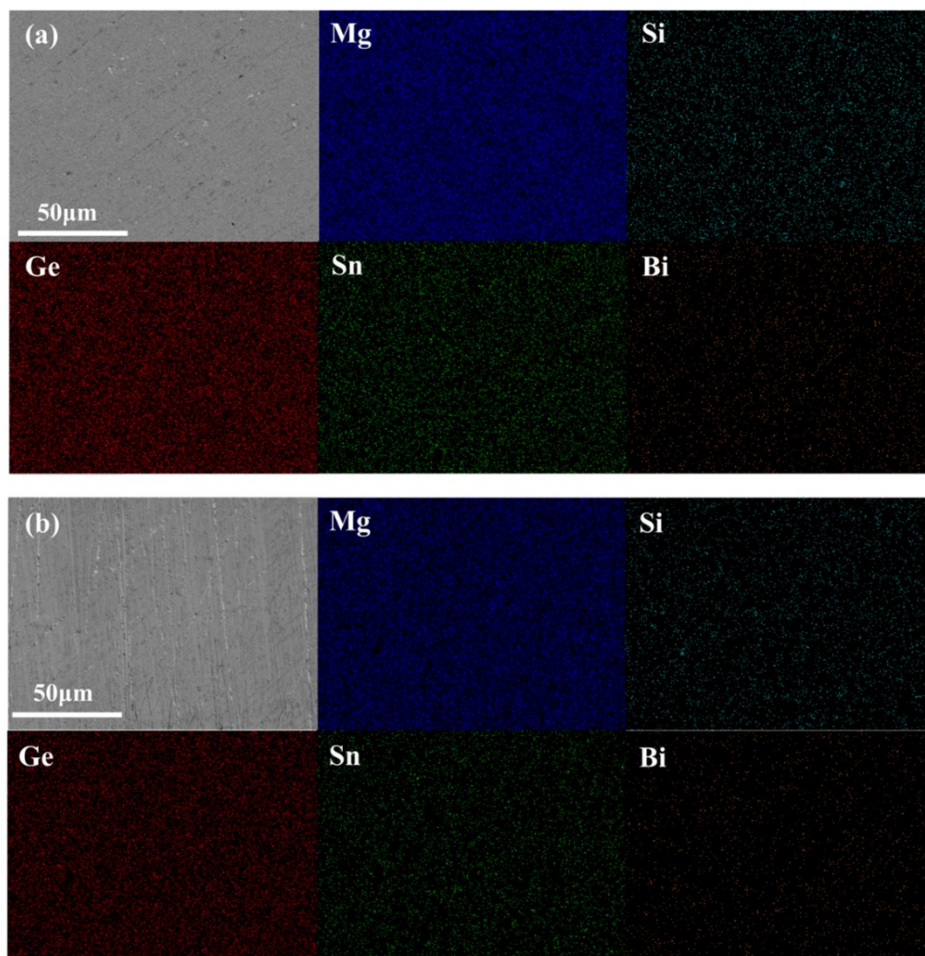


Fig. S15. Backscattered electron microscopy (BSE) images and corresponding elemental energy-dispersive X-ray spectroscopy (EDS) mapping of $\text{Mg}_2\text{-}\delta\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.73}\text{Bi}_{0.02}$. (a) after three cycles; (b) after quenching.

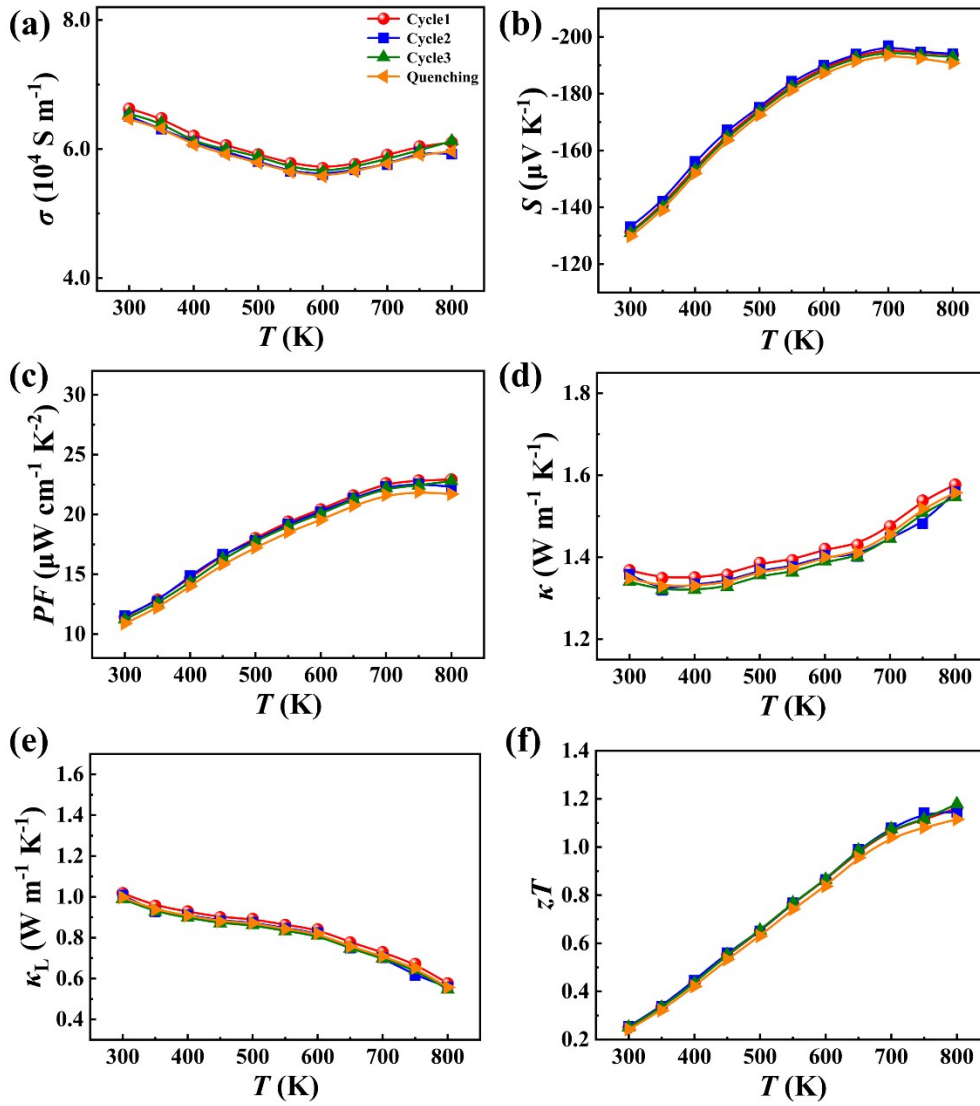


Fig. S16. Temperature dependence of (a) electrical conductivity σ , (b) Seebeck coefficient S , (c) power factor PF , (d) thermal conductivity κ , (e) lattice thermal conductivity κ_L and (d) zT of sample $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.55}\text{Bi}_{0.20}$ after three cycles and after quenching at 800 K.

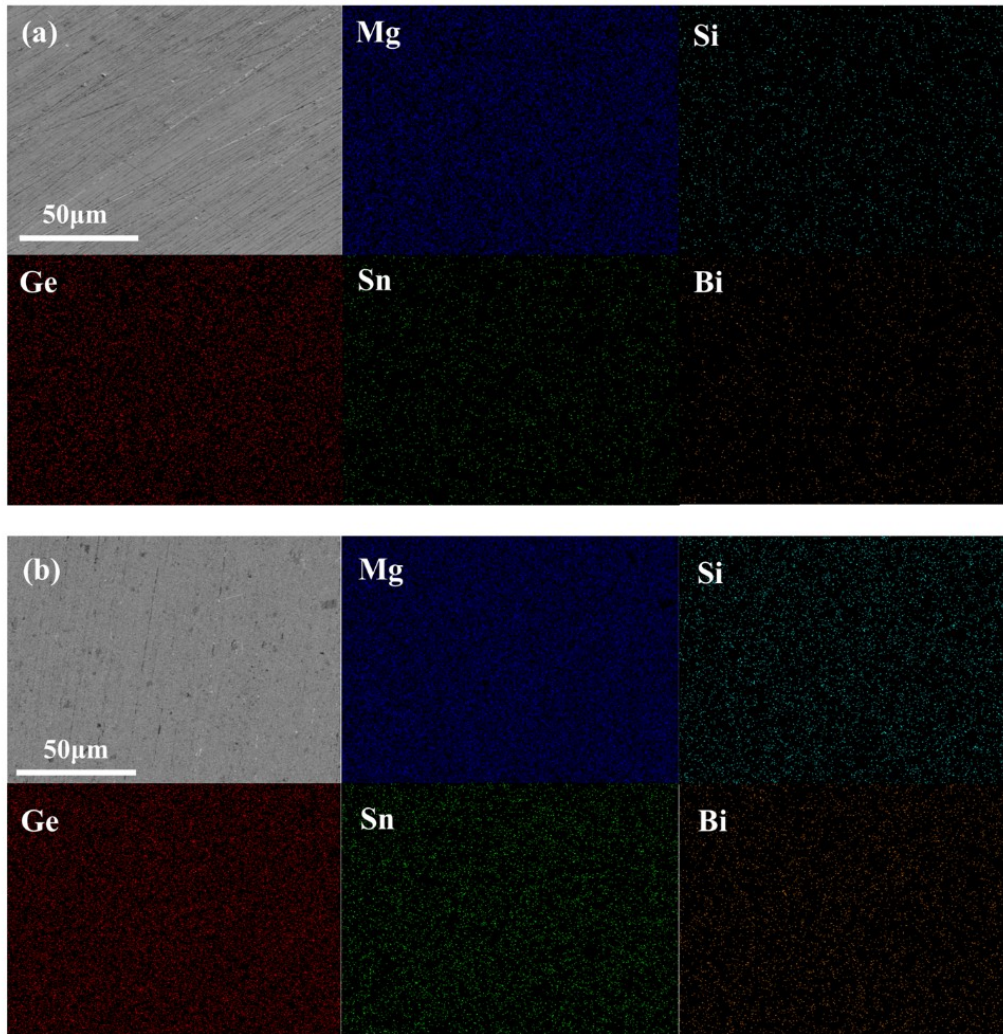


Fig. S17. Backscattered electron microscopy (BSE) images and corresponding elemental energy-dispersive X-ray spectroscopy (EDS) mapping of $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.55}\text{Bi}_{0.20}$ (a) after three cycles; (b) after quenching.

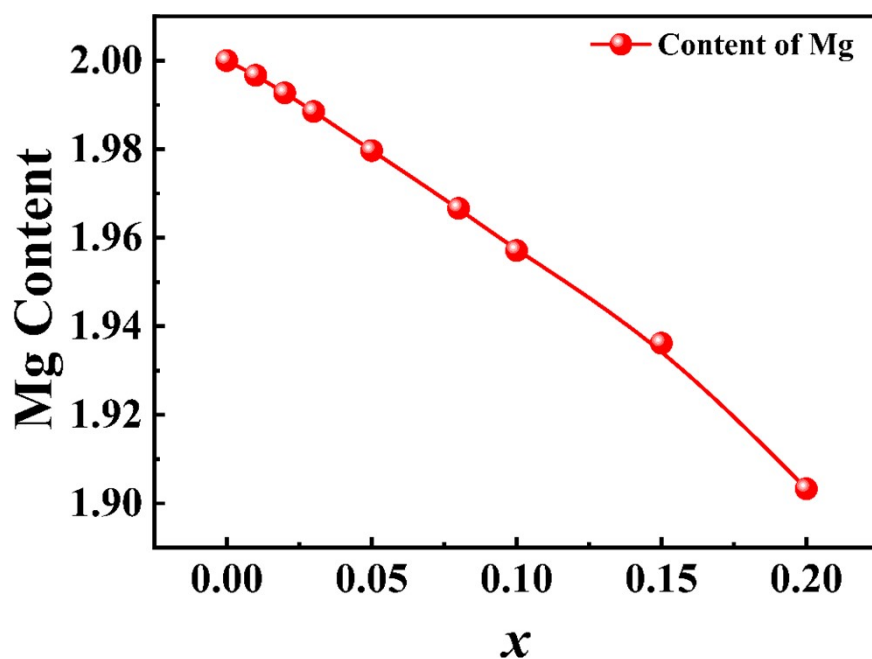


Fig. S18. Calculated Mg content as a function of Bi content x

Table S1. Atomic radius, electronegativity, and the most likely valence state of Mg, Si, Ge, Sn and Bi in $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.132}\text{Sn}_{0.748-x}\text{Bi}_x$.

Atom	Atomic radius (Å)	Electronegativity	Valence state
Mg	1.50	1.31	+2
Si	1.10	1.90	-4
Ge	1.25	2.01	-4
Sn	1.45	1.96	-4
Bi	1.60	2.02	-3

Table S2. Rietveld refined results of samples $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.75-x}\text{Bi}_x$

Composition	Atom	x	y	z	Occ.	U_{iso}	Site	Sym.
x=0	Mg	0.25	0.75	0.75	0.994	0.014	8c	-43m
	Si	0	0	0	0.120	0.006	4a	m-3m
	Ge	0	0	0	0.130	0.006	4a	m-3m
	Sn	0	0	0	0.750	0.006	4a	m-3m
x=0.01	Mg	0.25	0.75	0.75	1.000	0.011	8c	-43m
	Si	0	0	0	0.120	0.002	4a	m-3m
	Ge	0	0	0	0.130	0.002	4a	m-3m
	Sn	0	0	0	0.740	0.002	4a	m-3m
	Bi	0	0	0	0.010	0.002	4a	m-3m
x=0.02	Mg	0.25	0.75	0.75	0.981	0.014	8c	-43m
	Si	0	0	0	0.120	0.006	4a	m-3m
	Ge	0	0	0	0.130	0.006	4a	m-3m
	Sn	0	0	0	0.730	0.006	4a	m-3m
	Bi	0	0	0	0.020	0.006	4a	m-3m
x=0.03	Mg	0.25	0.75	0.75	0.983	0.014	8c	-43m
	Si	0	0	0	0.120	0.005	4a	m-3m
	Ge	0	0	0	0.130	0.005	4a	m-3m
	Sn	0	0	0	0.720	0.005	4a	m-3m
	Bi	0	0	0	0.030	0.005	4a	m-3m
x=0.05	Mg	0.25	0.75	0.75	0.973	0.014	8c	-43m
	Si	0	0	0	0.120	0.005	4a	m-3m

Bi	0	0	0	0.200	0.009	4a	m-3m
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Table S3. Quality factor of Rietveld refinement.

Composition	R_p	wR_p	R_{obs}	wR_{obs}	R_{all}	wR_{all}
x=0	4.63%	6.30%	1.51%	2.40%	1.51%	2.40%
x=0.01	5.04%	6.93%	1.52%	2.01%	1.52%	2.01%
x=0.02	5.09%	6.58%	1.49%	2.24%	1.49%	2.24%
x=0.03	3.64%	5.17%	1.20%	1.87%	1.20%	1.87%
x=0.05	5.10%	6.45%	1.45%	2.23%	1.45%	2.23%
x=0.08	4.16%	5.98%	1.19%	1.82%	1.19%	1.82%
x=0.10	6.97%	8.78%	1.29%	2.53%	1.29%	2.53%
x=0.15	6.13%	7.75%	1.11%	2.15%	1.11%	2.15%
x=0.20	5.26%	7.26%	1.33%	2.24%	1.33%	2.24%

Table S4. ADP and error of samples $Mg_{2-\delta}Si_{0.12}Ge_{0.13}Sn_{0.75-x}Bi_x$.

Composition	Mg U_{iso} (\AA^2)	Error (\AA^2)	X U_{iso} (\AA^2)	Error (\AA^2)
x=0	0.0144	4×10^{-4}	0.006	1.2×10^{-4}
x=0.01	0.015	4×10^{-4}	0.007	1.2×10^{-4}
x=0.02	0.0138	4×10^{-4}	0.00634	1.3×10^{-4}

$x=0.03$	0.014	3×10^{-4}	0.00541	9×10^{-5}
$x=0.05$	0.0138	4×10^{-4}	0.00526	1.4×10^{-4}
$x=0.08$	0.016	2×10^{-4}	0.0073	1×10^{-4}
$x=0.10$	0.0202	9×10^{-4}	0.0137	8×10^{-4}
$x=0.15$	0.0236	7×10^{-4}	0.0168	6×10^{-4}

Table S5. PDF fits results of samples $\text{Mg}_{2-\delta}\text{Si}_{0.12}\text{Ge}_{0.13}\text{Sn}_{0.65}\text{Bi}_{0.10}$.

Structure	Atom	x	y	z	Occ.	U_{iso}	Site	Sym.
Broken symmetry model	Mg1	0.234	0.766	0.766	0.960	3.520	1a	1
	Mg2	0.765	0.234	0.766	0.960	3.520	1a	1
	Mg3	0.765	0.766	0.234	0.960	3.520	1a	1
	Mg4	0.234	0.234	0.234	0.960	3.520	1a	1
	Mg5	0.765	0.234	0.234	0.960	3.520	1a	1
	Mg6	0.234	0.766	0.234	0.960	3.520	1a	1
	Mg7	0.765	0.766	0.766	0.960	3.520	1a	1
	Mg8	0.234	0.234	0.766	0.960	3.520	1a	1
	Si	0	0	0	0.120	2.180	4a	m-3m
	Ge	0	0	0	0.130	2.180	4a	m-3m
	Sn	0	0	0	0.650	2.180	4a	m-3m
	Bi	0	0	0	0.100	2.180	4a	m-3m
	Mg	0.25	0.75	0.75	0.961	0.045	8c	-43m

Si	0	0	0	0.120	0.021	4a	m-3m
Ge	0	0	0	0.130	0.021	4a	m-3m
Sn	0	0	0	0.650	0.021	4a	m-3m
Bi	0	0	0	0.100	0.021	4a	m-3m

Table S6. The atomic radius, ionic radius, and orbital energy of Mg, Si, Ge, Sn and Bi.

Elements	Atomic radius (Å)	Ionic radius (Å)	Orbital energy
Mg	1.50	0.86	-7.65 eV (s orbital)
Si	1.10	0.54	-7.78 eV (p orbital)
Ge	1.25	0.67	-7.54 eV (p orbital)
Sn	1.45	0.83	-7.01 eV (p orbital)
Bi	1.60	1.17	-8.15 eV (p orbital)

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