Theoretical insight into the influence of different molecular design strategies on photovoltaic properties for a series of POM-based dyes applied in dye-sensitized solar cells

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## 1. Theoretical background

The power conversion efficiency  $(\eta)$  of DSSCs can be expressed as [1]:

$$\eta = FF \frac{V_{\rm OC} J_{\rm SC}}{I_{\rm S}} \tag{1}$$

where  $I_S$  is the intensity of the incident light and *FF* is the fill factor of the cell. The  $J_{SC}$  can be defined as follows:

$$J_{\rm SC} = \int LHE(\lambda) \Phi_{\rm inject} \eta_{\rm collect} I_{\rm S}(\lambda) d\lambda$$
<sup>(2)</sup>

where  $LHE(\lambda)$ , the light harvesting efficiency at a specific wavelength, is determined by  $LHE(\lambda) = 1 - 10^{-f}$  and f is the oscillator strength of dyes corresponding to the maximum absorption  $\lambda_{\text{max}}$ .  $\eta_{\text{collect}}$  is the charge collection efficiency, which is assumed to be a constant for the same DSSCs.  $\Phi_{\text{inject}}$  is the electron injection efficiency, which is closely connected with  $k_{\text{inj}}$ , and can be described as follows [2]:

$$k_{\rm inj} = \frac{2\pi}{h} V_{\rm DA}^2 \left(\frac{\pi}{\chi K_{\rm B}T}\right)^{1/2} \exp\left[-\frac{(\Delta G_{\rm inj} + \chi)}{4\chi K_{\rm B}T}\right]$$

where  $V_{DA}$  is the coupling between the donor and acceptor states and  $\chi$  is the reorganization energy.  $\Delta G_{inj}$  can be determined by the following equation [3]:

$$\Delta G_{\rm inj} = E^*_{\rm dye} - E_{\rm CB} \tag{3}$$

where  $E_{dye}^*$  is the oxidation potential of dye in the excited state, which is related to the oxidation potential of the dye in the ground state ( $E_{dye}$ ) and the vertical transition energy ( $E_v$ ), i.e.,  $E_{dye}^* = E_{dye} - E_v$  [3].  $E_{CB}$  is the reduction potential of the TiO<sub>2</sub> conduction band (CB) and the experimental value -4.00 eV (vs vacuum) is used [4].  $\chi$ can be obtained using the following formula:

$$\chi = [E_0^+ - E_+] + [E_+^0 - E_0] \tag{4}$$

Where  $E_0^+$  is the energy of the cation calculated based on the optimized structure of

the neutral molecule,  $E_+$  represents the energy of the cation calculated based on the optimized cation structure,  $E_+^0$  corresponds to the energy of the neutral molecule calculated at the cationic state and  $E_0$ , is the energy of the neutral molecule at the ground state.

In order to evaluate the ICT abilities of dyes, the ICT parameters including the amount of transferred charges ( $q^{\text{CT}}$ ), the corresponding effective charge transfer distance ( $d^{\text{CT}}$ ) and the *t* index that assess the degree of separation between  $\rho^+$  (*r*) and  $\rho^-$  (*r*) based on the total densities for ground and excited states were calculated [5,6].  $\rho^+$  (*r*) and  $\rho^-$  (*r*) are defined as the points in space where the density increment and depletion upon absorption are produced. The larger *t* is, the little overlap between the electron density depletion and increment regions exists. The difference of electronic densities related to the electronic transition is given by:

$$\Delta \rho (r) = \rho_{\rm ES} (r) - \rho_{\rm GS} (r) \tag{5}$$

 $\rho_{\text{ES}}(r)$  and  $\rho_{\text{GS}}(r)$  are proposed to represent the electronic densities of excited and ground states, respectively.  $\rho^+(r)$  and  $\rho^-(r)$  are defined as the points in space where the density increment and depletion upon absorption are produced. And  $q^{\text{CT}}$  can be proposed:

$$q^{\rm CT} = \int \rho^+(r) \,\mathrm{d}r \tag{6}$$

The barycenters  $(r^+ \text{ and } r^-)$  of density distributions defined by  $\rho^+$  (r) and  $\rho^-$  (r) are written as the following equations:

$$r^{+} = (\mathbf{x}^{+}, \mathbf{y}^{+}, \mathbf{z}^{+}) = 1/q^{\text{CT}} \int r \rho^{+}(r) \, \mathrm{d}r$$
(7)

$$r^{-} = (\mathbf{x}^{-}, \mathbf{y}^{-}, \mathbf{z}^{-}) = 1/q^{\text{CT}} \int r \rho^{-}(r) \, \mathrm{d}r$$
(8)

The difference between  $r^+$  and  $r^-$  is defined as  $d^{CT}$ .

To elucidate the intrinsic reasons of different interaction energy ( $\Delta E_{tot}$ ) for the dimers of dyes, simple energy decomposition analysis (EDA) was performed.  $\Delta E_{tot}$  can be decomposed as [7]:

$$\Delta E_{\text{tot}} = (\Delta E_{\text{els}} + \Delta E_{\text{ex}}) + \Delta E_{\text{orb}} = \Delta E_{\text{steric}} + \Delta E_{\text{orb}}$$
(9)

where  $\Delta E_{\text{els}}$  is electrostatic interaction term, and normally negative if two fragments are neutral;  $\Delta E_{\text{ex}}$  is exchange repulsion term, which comes from the Pauli repulsion effect and is invariably positive. For convenience,  $\Delta E_{\text{els}}$  and  $\Delta E_{\text{ex}}$  terms are usually combined as steric term ( $\Delta E_{\text{steric}}$ ).  $\Delta E_{\text{orb}}$  is orbital interaction term, which arises from the mix of occupied molecular orbital (MOs) and virtual MOs. If the combined wavefunction is used as initial guess for complex, then  $E_{\text{orb}}$  can be evaluated by subtracting the first SCF iteration energy from the last SCF iteration energy:

$$\Delta E_{\rm orb} = E_{\rm SCF, \, last} - E_{\rm SCF, \, 1st} \tag{10}$$

$$\Delta E_{\text{steric}} = \Delta E_{\text{els}} + \Delta E_{\text{ex}} = \Delta E_{\text{tot}} - \Delta E_{\text{orb}}$$
(11)

As for  $V_{OC}$ , it can be defined by [8]:

$$V_{\rm OC} = \frac{E_{\rm CB}}{q} + \frac{kT}{q} \ln(\frac{n_{\rm c}}{N_{\rm CB}}) - \frac{E_{\rm redox}}{q}$$
(12)

where  $n_c$  is the number of electrons in CB, while q and kT are constants, representing the unit charge and thermal energy, respectively.  $N_{CB}$  represents the accessible density of CB states. The dye absorbed on TiO<sub>2</sub> surface can lead to the energy shift of  $E_{CB}$ ( $\Delta E_{CB}$ ). Apparently, a dye with large  $n_c$  and  $\Delta E_{CB}$  will induce a significant increase of  $V_{OC}$ .

## 2. The multicenter bond index

The multicenter bond index is a measure of aromaticity, and the higher the number, the stronger the aromaticity. The multicenter bond index is defined as:

$$I_{ABCDEF...K} = \sum_{a \in A} \sum_{b \in B} \sum_{c \in C} \dots \sum_{k \in K} \left[ (PS)_{ab} (PS)_{bc} (PS)_{cd} \dots \right] (PS)_{ka}$$
(13)

where P is the single electron density matrix, S is the overlapping matrix, a, b, c... are the number of the basis function, A, B, C... represent the number of atoms in the ring, and the atoms are adjacent according to their connection in the ring.

## 3. Simulations of the electron injection

The time-dependent survival probabilities (TDSP) curves were defined as the probability of the photo-excited electron which is still in the adsorbed dye molecule at time *t*. Therefore, the TDSP can be computed by applying the time-evolved electronic wave function into the atomic orbitals of the adsorbed dye molecule.

The time-evolved wave function  $\varphi(t)$  can be written as a linear combination of atomic orbitals:

$$\varphi(t) = \sum_{i,j} B_{i,j}(t) \phi_{i,j} \tag{14}$$

where  $\phi_{i,j}$  represents the orbitals *j* of the *i*-th atom. The expansion coefficients  $B_{i,j}(t)$ , introduced in Eq. (13), can now be computed according to the follow equation:

$$B_{i,j}(t) = \sum_{k} Q_{i,j}^{k} C_{k} exp(-\frac{i}{\hbar} E_{k} t)$$
(15)

The coefficient  $C_k$  in Eq. (14) is defined by the expansion of the initial state in an orthonormal basis set of  $\phi(k)$ ,

$$\varphi(0) = \sum_{k} C_k \phi(k) \tag{16}$$

The coefficient  $Q_{i,j}^k$  in this equation is defined according to the expansion of  $\phi(k)$  as

a linear combination of the atomic orbitals:

$$\boldsymbol{\phi}(k) = \sum_{j} Q_{j}^{k} \boldsymbol{\phi}_{j} \tag{17}$$

The eigenvalue  $E_k$  in Eq. (14) can be obtained by solving the extend Hückel theory (EH) eigenvalue problem:

$$HQ^{k} = E_{k}SQ^{k}$$
<sup>(18)</sup>

where H is the EH matrix and S is the overlap matrix in the atomic orbital basis. The non-diagonal Hamiltonian matrix is here determined by the Wolfsberg-Helmholtz constant K and the overlapping matrix element:

$$H_{j,m} = KS_{jm} \frac{H_{jj} + H_{mm}}{2}$$
(19)

Usually, the constant K is set to be 1.75, and the  $S_{jm}$  is defined according to the overlap of the atomic orbitals:

$$S_{jm} = \langle \varphi_j | \varphi_m \rangle \tag{20}$$

Therefore, the projection of the time-evolved electronic wave function onto the atomic orbitals of the adsorbed dye molecular can be obtained as follows:

$$\rho_{\text{MOL}}(t) = \left|\sum_{j}^{MOL} \sum_{m} B_{j}^{*}(t) B_{m}(t) S_{jm}\right|$$
(21)

Note that the sum over m includes all of the atoms in the research object, whereas the sum over j only includes the atoms in the adsorbed dye molecular.

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Figure S1. Initial structure of the  $(TiO_2)_{48}$  cluster in side and top view.



Figure S2. Frontier molecular orbital distributions of studied dyes.



Figure S3. Electrostatic potential maps for dimers of studied dyes.



Figure S4. Optimized geometrical structures of dyes adsorbed on TiO<sub>2.</sub>

**Table S1.** The calculated excitation energies  $E_v$  (eV), maximum absorption wavelengths  $\lambda_{max}$  (nm), oscillator strengths f and major electronic compositions of studied dyes.

dye	$E_{ m v}$	$\lambda_{\max}$	f	Major electronic compositions
1	1.99	623	0.99	H→L(84%)
	2.86	434	0.40	H−1→L(71%)
2	1.76	703	0.55	H→L(88%)
	2.39	519	1.34	H−1→L(85%)
3	1.66	746	0.54	H→L(91%)
	2.48	501	1.47	H−1→L(81%)
4	1.58	786	0.51	H→L(91%)
	2.44	508	1.26	H−1→L(79%)