## **Supplementary information**

## **TiO2-Modified MoS<sup>2</sup> Monolayer Films Enables Sensitive NH<sup>3</sup> Sensing at Room Temperature**

Lun Tan,<sup>1,2#</sup>, Xianzhen Liu<sup>1#</sup>, Peng Wu<sup>2</sup>, Liwei Cao<sup>2</sup>, Wei Li<sup>2</sup>, Ang Li<sup>2,\*</sup>, *Zhao Wang1,\* , Haoshuang Gu1, \**

<sup>1</sup>Hubei Engineering Research Center for Safety Detection and Control of Hydrogen Energy - Hubei Key Laboratory of Micro-Nanoelectronic Materials and Devices, School of Microelectronics, Hubei University, Wuhan 430062, P.R. China.

2 Institute of Microstructure and Properties of Advanced Materials, Beijing University of Technology, Beijing, 100124, P.R. China

\*Corresponding authors, E-mail: [wangzhao@hubu.edu.cn](mailto:wangzhao@hubu.edu.cn); [ang.li@bjut.edu.cn;](mailto:ang.li@bjut.edu.cn) guhsh@hubu.edu.cn

# Both authors contribute equally to this work.



**Fig.S1** The growth of monolayer  $MoS<sub>2</sub>$  films. (a) Schematic of the CVD apparatus. (b) Temperature profile and Ar flow during the growth process. (d-f) The SEM images of MoS2 grown at 650℃, 700 $\degree$ C, 750 $\degree$ C and 800 $\degree$ C, respectively. The morphology of MoS<sub>2</sub> at different growth temperature in 10 min is shown in **Fig S1.** When the temperature is below 700°C, many separate irregular polygonal-like crystals can be seen in **Fig S1(c)**. The Raman spectrum of those small crystals is corresponding to the  $MoO<sub>2</sub>$  and  $MoOS<sub>2</sub>$  which are the intermediate product during the growth of  $MoS<sub>2</sub>$  because of incomplete sulfurization at low temperature as shown in Fig S2. When the temperature reaches to  $700\text{°C}$ , there are many separate  $MoS<sub>2</sub>$  crystals on substrate. The Raman spectrum of those small crystals in **Fig S2(a)** shows the  $E_{2g}$  and  $A_{1g}$  peaks of MoS<sub>2</sub>, the frequency difference ( $\Delta k$ ) between the  $E^{I}_{2g}$  and  $A_{Ig}$  peaks is ~ 20 cm<sup>-</sup> <sup>1</sup>, which is corresponding to the monolayer  $MoS<sub>2</sub>$ . The  $MoS<sub>2</sub>$  crystals gradually connects into continuous polycrystalline films, when temperature rises to  $750^{\circ}$ C due to the accelerated reaction rate at high

temperature. At 800°C,  $\Delta k$  of films is ~ 24 cm<sup>-1</sup> in Fig S2, which indicated that the films are multilayer  $MoS<sub>2</sub>$  films. The rapid evaporation of  $MoO<sub>3</sub>$  makes the pressure of  $MoO<sub>3</sub>$  higher on the surface of substrates, which results in the growth of multilayer  $MoS<sub>2</sub>$ as reported in ref.1.[1]



**Fig.S2** The Raman spectrum of production grown at 650℃, 750℃ and 800°C, respectively.



Fig.S3 SEM images of MoS<sub>2</sub> films at different growth time. When the growth time is shorter than 5 min, there are separated triangle-like  $MoS<sub>2</sub>$ nanoflakes on substrate. When the growth time is extended to 10 min, the separated  $MoS<sub>2</sub>$  nanoflakes connected into polycrystalline films. The new nuclei will form on the surface of  $MoS<sub>2</sub>$  films with the further increase of growth time, which lead to the increase of the thickness of  $MoS<sub>2</sub>$  films. Once the growth time exceeds 20 min, there are some pores on the  $MoS<sub>2</sub>$ films because of the poor thermal-stability of  $MoS<sub>2</sub>$  films in S insufficient environment.



**Fig. S4** Fabrication and photography of  $TiO<sub>2</sub>$  modifying  $MoS<sub>2</sub>$  monolayer films sensors. (a) the device was annealed at 250℃ in Ar for contacting tightly. (b) Photograph of a monolayer  $MoS<sub>2</sub>$  films with interdigitated electrodes.



Fig. S5 AFM images of 10 nmol/cm<sup>2</sup> TiO<sub>2</sub> modifying monolayer. (a) and (b) represent the 2D and 3D AFM images of sensors. (c) particle size distribution of  $TiO<sub>2</sub>$  nanoparticles. The  $TiO<sub>2</sub>$  nanoparticles are uniformly distributed on the surface of monolayer  $MoS<sub>2</sub>$  films except for several big cluster (>100 nm). The average size of about those nanoparticles is about 55 nm.



**Fig. S6** HADDF images and EDS elements mappings of hybrid. It is clearly that the  $MoS<sub>2</sub>$  is wide distributed. the sample was obtained by scraping from the substrate, which led to the cluster and wrinkle of  $MoS<sub>2</sub>$ films. Therefore, there are few large area uniform flat  $MoS<sub>2</sub>$  films on carbon films. The size of  $TiO<sub>2</sub>$  nanoparticle is about 50 nm, which is consistent with the result obtained by AFM. **Fig S6 (e)** is the selected area electron diffraction (SAED) pattern of the **Fig 2(a)**.



Fig. S7 The XPS spectrum of TiO<sub>2</sub> nanoparticles modifying monolayered  $MoS<sub>2</sub>$  films.



**Fig. S8** Comparison of sensor response and response time of different sensors towards 500 ppm of ammonia in air. The response time varied from 5 min to 12 min with the increase amount of  $TiO<sub>2</sub>$  nanoparticles. -

**Table. S1** Linear fit of each sensor towards  $50-1000$  ppm. the  $R^2$  is greater than 0.94, indicating that the response of each sensor has a good linear relationship with the concentration of  $NH<sub>3</sub>$  from 50-1000 ppm.

<b>Sample</b>	$R^2$	slope	
$\bf S0$	0.94	0.0020	
S <sub>1</sub>	0.95	0.0101	
S <sub>2</sub>	0.96	0.0111	
<b>S3</b>	0.98	0.0150	
S <sub>4</sub>	0.96	0.0082	





Fig. S9 The dynamic response of 30 nmol/cm<sup>2</sup> TiO<sub>2</sub> nanoparticles modifying monolayer  $MoS<sub>2</sub>$  films towards 50 ppm  $NH<sub>3</sub>$  at different humidity.



Fig. S10 DFT calculation details. (a)Top view of a monolayer  $MoS<sub>2</sub>$ showing hole top, S top and Mo top for  $N(NH_3)$  adsorption. (b) and (c) are isosurface plot of the electron charge density difference for  $NH_3$  on monolayer  $MoS<sub>2</sub>$  and  $TiO<sub>2</sub>-MoS<sub>2</sub>$  respectively. To study the adsorption of NH<sub>3</sub> on the MoS<sub>2</sub>,  $3 \times 3 \times 1$  supercell of 2H-MoS<sub>2</sub> was built as shown in **Fig S10(a)**. A large vacuum layer of 15 Å was used to avoid interlayer interactions. A plane wave cutoff of 450 eV was consistently used during the whole process. In the process of geometry optimization, the convergency of energy and atomic force were  $10^{-5}$  eV and  $0.03$  eV/Å, respectively. A Monkhorst-Park mesh of  $3 \times 3 \times 1$  for the Brillouin zone integration was employed. The adsorption energy of  $NH<sub>3</sub>$  molecules on  $MoS<sub>2</sub>$  was calculated using  $E<sub>ads</sub> = E<sub>total</sub> - (E<sub>MoS2</sub> or TiO2-MoS2} + E<sub>NH3</sub>)$ , where

 $E_{total}$  is the total energy of the NH<sub>3</sub> molecules on MoS<sub>2</sub> (or TiO<sub>2</sub>-MoS<sub>2</sub>) surface system,  $E_{MoS2}$  ( $E_{TiO2-MoS2}$ ) is the energy of monolayered  $MoS2$  $(TiO<sub>2</sub>-MoS<sub>2</sub>), E<sub>NH3</sub>$  is the energy of isolated NH<sub>3</sub> molecule.

To explore the underlying mechanism of the significant resistance changes upon the adsorption of  $NH_3$  molecules on  $MoS_2$ , DFT calculation was carried to analysis the  $NH_3$  adsorption on  $MoS_2$  as shown in **Fig. S10**. There are seven absorption sites are listed in **Fig S10(a)** which are depended on the sites of N and orientation of N-H bond. Table S2 shows the absorption energy of  $NH_3$  molecules on each site. It is clearly that the absorption energy of H atoms near S (M1 and H1) are lower than far one  $(S1$  and S2), which indicate that the absorption site of H in NH<sub>3</sub> molecule atom close S atom is more stable. this is consistent with the charge density difference of absorption system in **Fig. S10(b)** because the obvious charge transfer between H and S atoms. The blue area and red area represent electrons depletion and accumulation, respectively. It is clearly shown that the electrons were transferred from  $NH<sub>3</sub>$  to  $MoS<sub>2</sub>$ , which resulted in the decrease of the resistance when  $MoS<sub>2</sub>$  exposed to  $NH<sub>3</sub>$  and this is consistent with the experimental results of bare  $MoS<sub>2</sub>$  films towards  $NH<sub>3</sub>$ . Obviously, the N atom on the Mo top and the H atom orienting to S atom has small absorption energy, which is consistent with the result of the easily charge transferring between S and H atoms in **Fig S10 (a)**. Isosurface plot of the electron charge density difference for  $NH_3$  on  $TiO_2$ -MoS<sub>2</sub> is shown in the

**Fig. S10(c).** It shows that the  $NH_3$  molecule interact with O atoms of  $TiO_2$ , which may cause the decrease of sensor resistance. The table S2 shows the absorption energy of  $NH_3$  molecule absorbed on the surface of  $MoS_2$  and  $TiO<sub>2</sub>-MoS<sub>2</sub>$ , respectively. The adsorption energy is smaller when ammonia molecule adsorbed on the  $TiO<sub>2</sub>$ -MoS<sub>2</sub> surface, indicating that ammonia molecule is more easily adsorbed on the  $TiO<sub>2</sub>$  surface than on  $MoS<sub>2</sub>$  surface since  $TiO<sub>2</sub>$  has more active sites.

Table. S2 The adsorption energy of NH<sub>3</sub> molecule on different site of  $MoS<sub>2</sub>$  and TiO<sub>2</sub>-MoS<sub>2</sub>.

Absorption	H1			$H2$ $H3$ $S1$ $S2$	M1	M <sub>2</sub>	$TiO2$ -
sites							MoS <sub>2</sub>
$E_{\text{abs}}$ (eV) -0.170 -0.166 -0.167 -0.110 -0.107 -0.172 -0.166 -0.69							



Fig. S11. The LDOS (local density of states) of before and after NH<sub>3</sub> absorption on  $MoS<sub>2</sub>$  (a) and TiO<sub>2</sub> (b). The absorption of  $NH<sub>3</sub>$  molecules result in several distinct states at the valence bands, which is close to that of pristine  $MoS<sub>2</sub>$  without  $NH<sub>3</sub>$  absorption. Therefore, the NH<sub>3</sub> absorption don't have a substantial effect on the electronic structures of  $MoS<sub>2</sub>$ . Before the adsorption of the NH<sub>3</sub> molecules, the band gap of the TiO<sub>2</sub> is ~ 2 eV according to the calculation results, as shown in the **Fig S11**. After the  $NH<sub>3</sub>$  absorption, new electron states can be found in the band gap of  $TiO<sub>2</sub>$  according to the calculation results, which can be attributed to that  $NH<sub>3</sub>$  molecule interact with O atoms and finally lead to the narrowing of band gap to 0.7 eV. According to the first principle calculation results, the surface adsorption of  $NH_3$  on TiO<sub>2</sub> will have a strong impact on the electron structure of the  $TiO<sub>2</sub>$ , which may lead to the variation of the electrical transportation behavior of the  $TiO<sub>2</sub>$  host materials.

## **Reference**

[1] J. Wei, J.-K. Huang, J. Du. Effect of the geometry of precursor crucibles on the growth of MoS2 flakes by chemical vapor deposition. *New Journal of Chemistry*. 2020, **44**, 21076-21084.