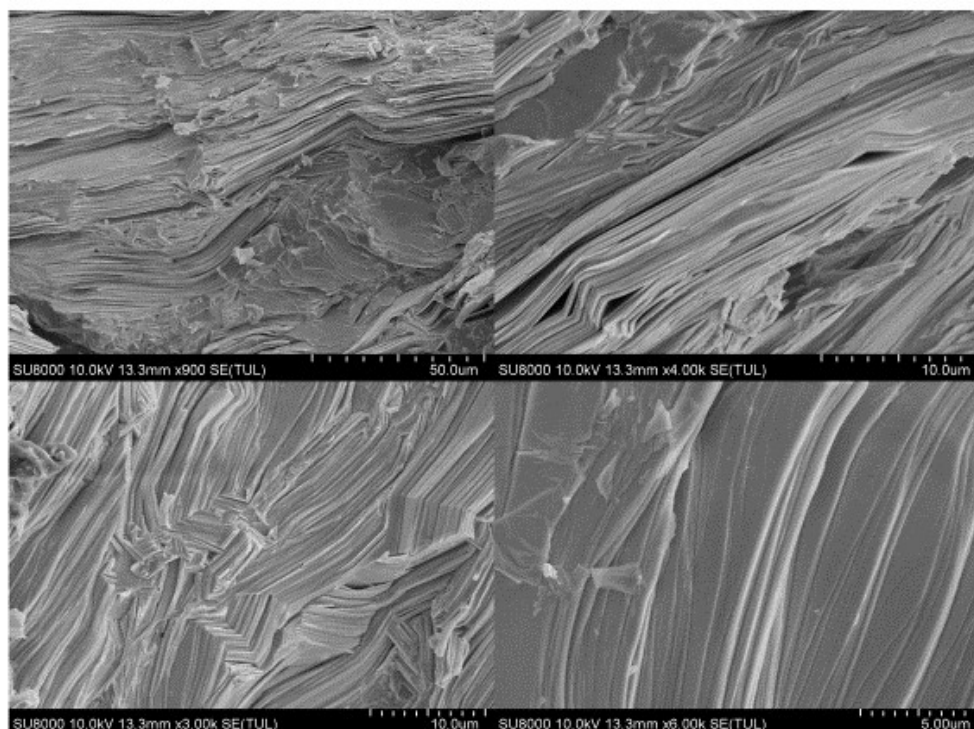


## Support Information

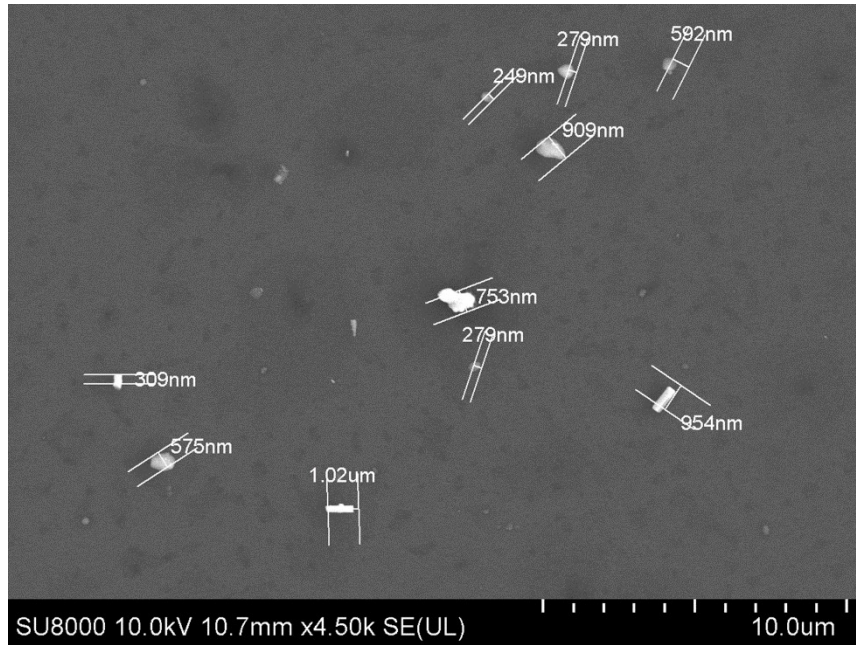
**Table S1** The density( $d$ ) and the relative density( $d_r$ ) of BTS- $x$ Ag<sub>2</sub>Se ( $x=0, 0.1, 0.15, 0.25$  and  $0.35$  vol%).

BTS- $x$ vol% Ag <sub>2</sub> Se	$d$ (g/cm <sup>3</sup> )	$d_r$ (%)
$x=0$	7.70	98.6%
$x=0.10$	7.58	97.1%
$x=0.15$	7.60	97.3%
$x=0.25$	7.50	96.0%
$x=0.35$	7.49	95.9%

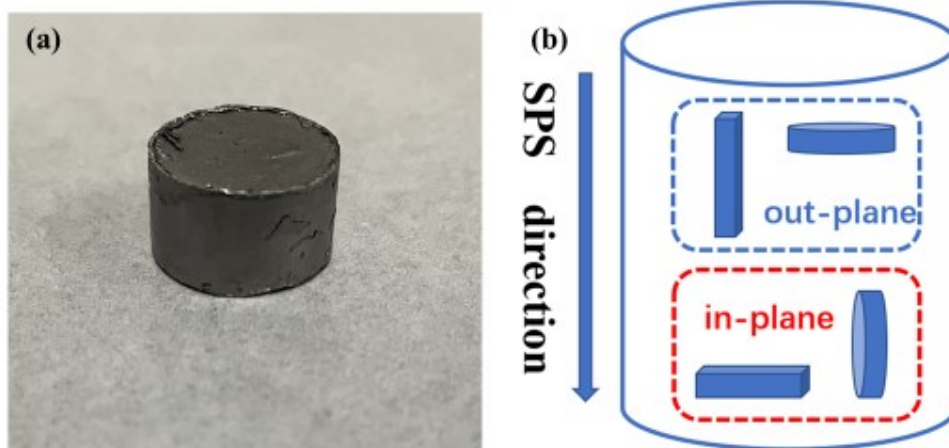
where  $d_r$  is relative density, defined as  $d_r=d/d_o$ , here  $d$  is the measured density and  $d_o$  (7.81g/cm<sup>3</sup>) is the theoretical density of Bi<sub>2</sub>Te<sub>2.5</sub>Se<sub>0.5</sub>. For composite specimens Bi<sub>2</sub>Te<sub>2.5</sub>Se<sub>0.5</sub>- $x$ vol%Ag<sub>2</sub>Se, the theoretical density is modified as:  $d_o=(1-x)d_1 + xd_2$ , where  $d_1=d_o$  for BTS matrix and  $d_2$  (8.05 g cm<sup>-3</sup>) is the theoretical density of Ag<sub>2</sub>Se.



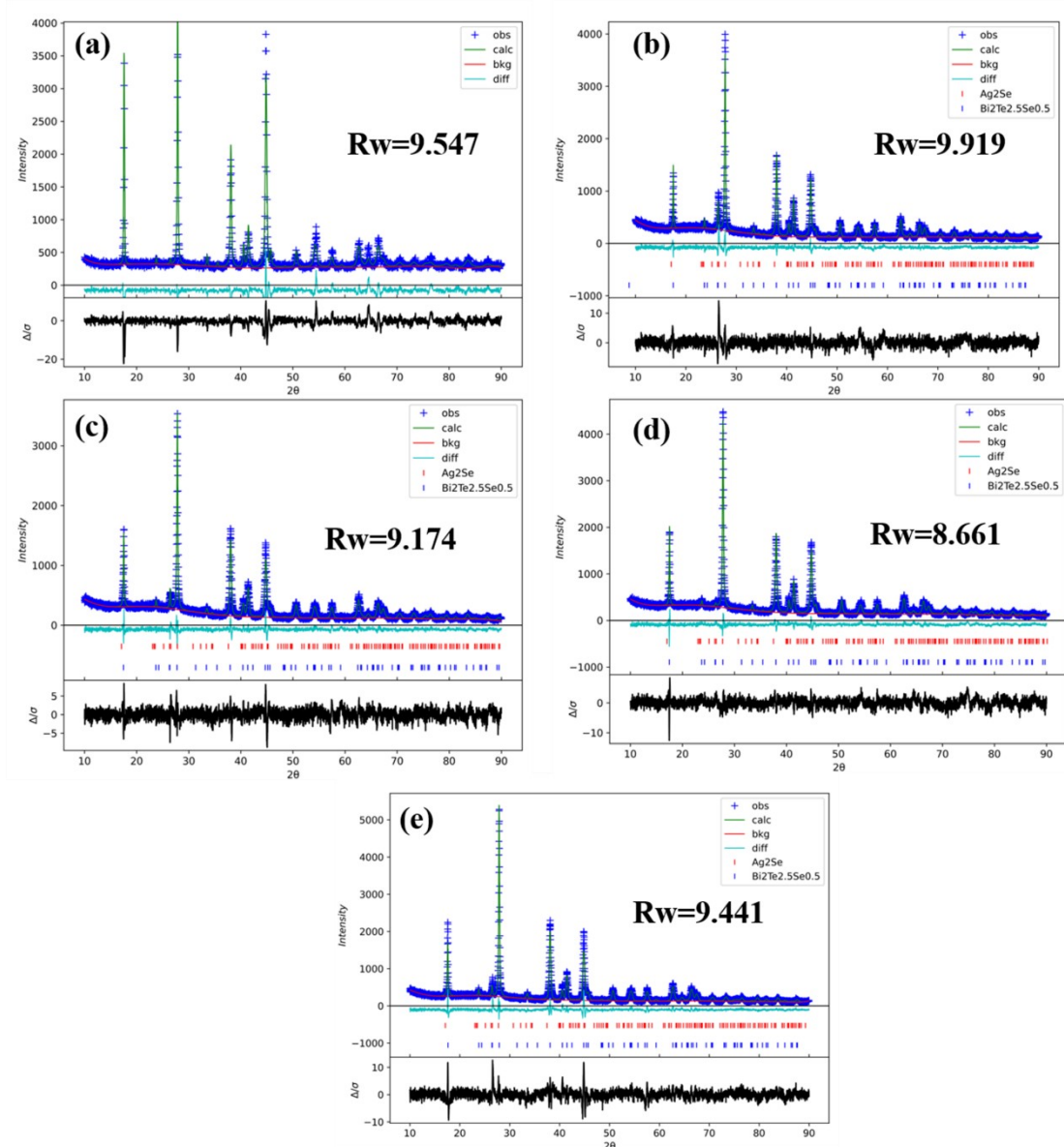
**Fig. S1.** SEM image of the fractured surface of the BTS.



**Fig. S2. SEM image of the Ag<sub>2</sub>Se particle**



**Fig. S3. Schematic illustration of the preparation of samples for measuring TE properties**



**Fig. S4.** The Rietveld refinement results for BTS- $x$ Ag<sub>2</sub>Se ( $x=0$  (a), 0.10 (b), 0.15 (c), 0.25 (d) and 0.35 (e) vol%)

## Calculation of $L$ number

For metals or heavily doped degenerate semiconductors,  $L$  is constant and equals to  $2.44 \times 10^{-8} \text{W}\Omega\text{K}^{-2}$ , whereas it approaches  $1.5 \times 10^{-8} \text{W}\Omega\text{K}^{-2}$  for nondegenerate semiconductors. On the basis of the single parabolic band (SPB) model with acoustic

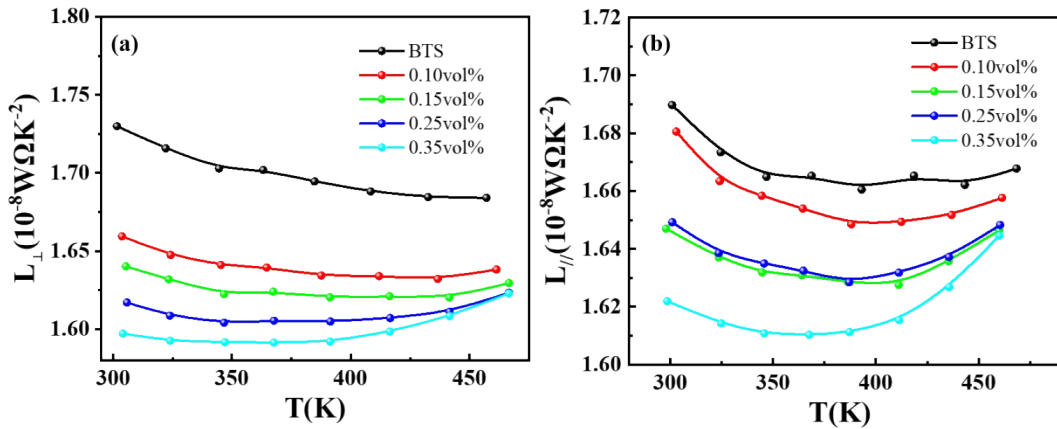
phonon scattering dominating, the Lorenz number,  $L$ , can be derived from the following equations<sup>[1]</sup>:

$$L = \left(\frac{k_B}{e}\right)^2 \left( \frac{\left(\lambda + \frac{7}{2}\right) F_{r+\frac{3}{2}}(\eta)}{\left(\lambda + \frac{3}{2}\right) F_{r+\frac{1}{2}}(\eta)} - \left[ \frac{\left(\lambda + \frac{5}{2}\right) F_{r+\frac{3}{2}}(\eta)}{\left(\lambda + \frac{3}{2}\right) F_{r+\frac{1}{2}}(\eta)} \right]^2 \right) \quad (S1)$$

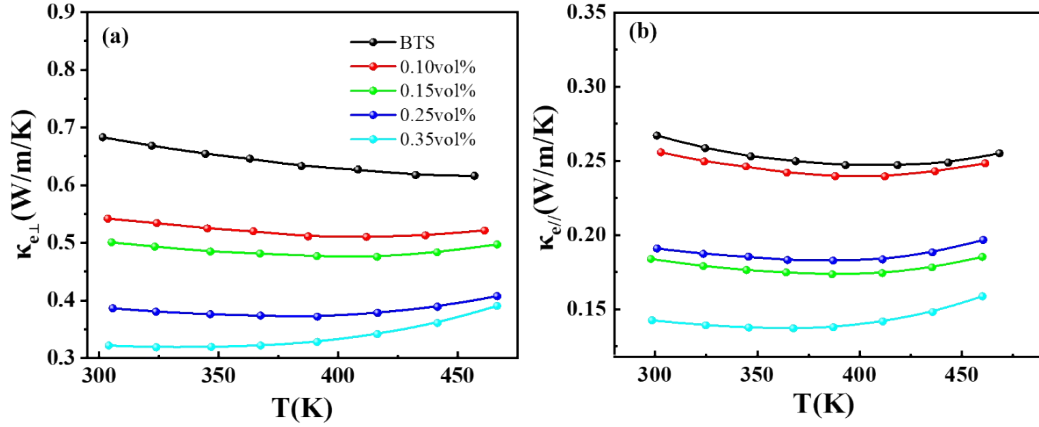
$$F_n(\eta) = \int_0^{\infty} \frac{x^n}{1 + e^{x-\eta}} dx \quad (S2)$$

$$S = \pm \frac{k_B}{e} \left( \frac{\left(\lambda + \frac{5}{2}\right) F_{\lambda+\frac{3}{2}}(\eta)}{\left(\lambda + \frac{3}{2}\right) F_{\lambda+\frac{1}{2}}(\eta)} - \eta \right) \quad (S3)$$

where  $k_B$  is the Boltzmann constant,  $e$  is the electron charge,  $\lambda$  is the scattering parameter equaling to -0.5 for acoustic phonon scattering,  $F_n(\eta)$  is the  $n$ th order Fermi integral and  $\eta$  is the reduced Fermi energy, which can be calculated from the measured Seebeck coefficient according to Equation S3. The results are shown in the Figure S5.



**Fig. S5.** The Lorenz number normal (a) and parallel (b) to the SPS direction for BTS- $x$ Ag<sub>2</sub>Se ( $x=0, 0.1, 0.15, 0.25$  and  $0.35$  vol%).



**Fig. S6.** The electron thermal conductivity normal (a,  $\kappa_{e\perp}$ ) and parallel (b,  $\kappa_{e//}$ ) to SPS direction for BTS- $x$ Ag<sub>2</sub>Se ( $x=0, 0.1, 0.15, 0.25$  and  $0.35$  vol%).

## Calculation of Lattice Thermal Conductivity

The lattice thermal conductivity  $\kappa_L$  can be calculated by using the Debye-Callaway model<sup>[2]</sup>:

$$\kappa_L = \frac{4\pi k_B^4 T^3}{\nu h^3} \int_0^{\frac{\theta}{T}} \tau(\zeta) \frac{\zeta^4 e^\zeta}{(e^\zeta - 1)^2} d\zeta \quad (\text{S4})$$

where  $h$ ,  $k_B$ ,  $\theta$ ,  $\nu$ ,  $\zeta$ , and  $\tau$  are the Planck constant, Boltzman constant, Debye temperature, phonon velocity, dimensionless variable, and total phonon relaxation time, respectively.  $\tau$  is associated with the multiple scattering centers, which can be expressed as<sup>[3]</sup>

$$\frac{1}{\tau} = \frac{1}{\tau_U} + \frac{1}{\tau_N} + \frac{1}{\tau_{PD}} + \frac{1}{\tau_B} + \frac{1}{\tau_{NP}} + \frac{1}{\tau_{DS}} \quad (\text{S5})$$

where  $\tau_U$ ,  $\tau_N$ ,  $\tau_{PD}$ ,  $\tau_E$ ,  $\tau_B$ ,  $\tau_{NP}$  and  $\tau_{DS}$  are relaxation time corresponding respectively to the scattering from phonon-phonon U- and N-process (U+N), point defects (PD), phase boundaries (B), and nanoparticles (NP), dislocation strain fields and dislocation cores (DS).

$$\tau_U^{-1} = \frac{\hbar\gamma^2 \omega^2 T}{M\nu^2 \theta_D} e^{-\theta_D/3T}$$

$$\tau_N^{-1} = \beta \tau_U^{-1} = \beta \frac{\hbar \gamma^2 \omega^2 T}{M v^2 \theta_D} e^{-\theta_D/3T}$$

$$\tau_{PD}^{-1} = \frac{V \omega^4}{4\pi v^3} \Gamma$$

$$\tau_B^{-1} = v/l$$

$$\tau_{NP}^{-1} = v \left[ (2\pi R)^2 + \left( \frac{4\pi R^2}{9} \left( \frac{\Delta D}{D} \right)^2 \left( \frac{\omega R}{v} \right)^4 \right)^{-1} \right]^{-1} N_P$$

$$\tau_{DS}^{-1} = 0.6 B_D^2 N_D \gamma^2 \omega \left\{ \frac{1}{2} + \frac{1}{24} \left( \frac{1-2r}{1-r} \right)^2 \left[ 1 + \sqrt{2} \left( \frac{v_L}{v_T} \right) \right] \right\} + N_D \frac{v^{-4/3}}{v_s^2} \omega^3$$

$$A = \frac{V \omega^4}{4\pi v^3} \Gamma$$

$$B = (1 + \beta) \frac{\hbar \gamma^2 \omega^2 T}{M v^2 \theta_D} e^{-\theta_D/3T}$$

$$C = 0.6 B_D^2 N_D \gamma^2 \left\{ \frac{1}{2} + \frac{1}{24} \left( \frac{1-2r}{1-r} \right)^2 \left[ 1 + \sqrt{2} \left( \frac{v_L}{v_T} \right) \right] \right\}$$

$$D = N_D \frac{v^{-4/3}}{v_s^2}$$

In the case of our sample, the total scattering rate  $\tau^{-1}$  can be shortened as

$$\frac{1}{\tau} = A \omega^4 + B \omega^2 \exp\left(-\frac{\theta_D}{3T}\right) + \frac{v}{l} + v \sigma N_p + C \omega + D \omega^3 \quad (\text{S6})$$

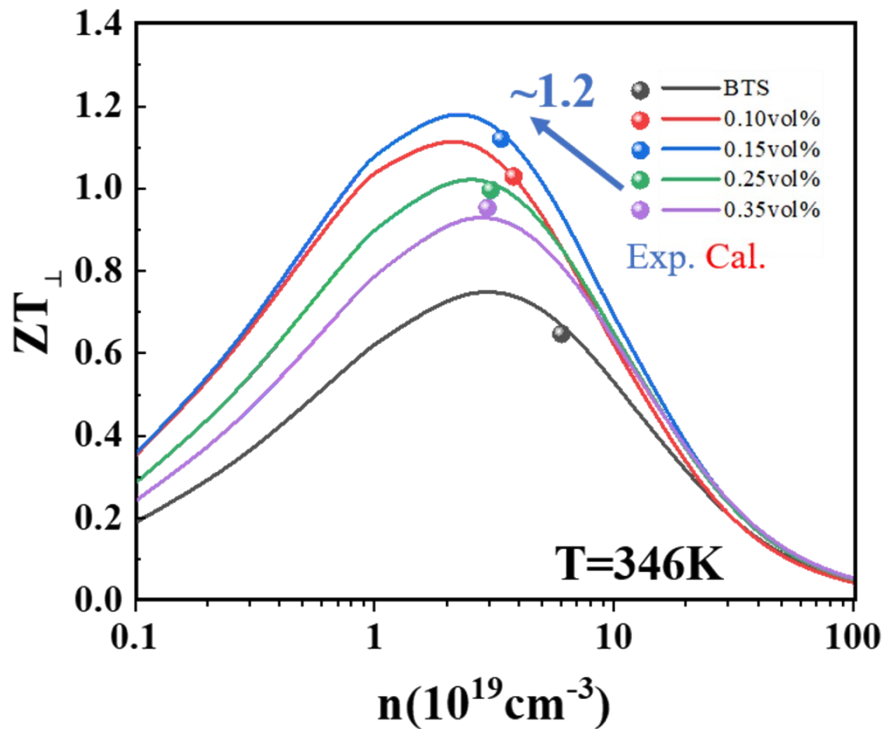
In the equation above,  $\omega$  is the phonon frequency and  $l$  is the average spatial distance among the nanoinclusions  $\text{Ag}_2\text{Se}$  in BTS while  $v/l$  corresponds to scattering from the phase boundaries.

To quantitatively calculate  $\kappa_L$ , the literature values<sup>[4]</sup> of Debye temperature  $\theta_D = 164\text{K}$  and sound velocity  $v = 1778\text{m s}^{-1}$  are used for  $\text{BTS}-x\text{Ag}_2\text{Se}$  ( $x=0, 0.1, 0.15, 0.25$  and

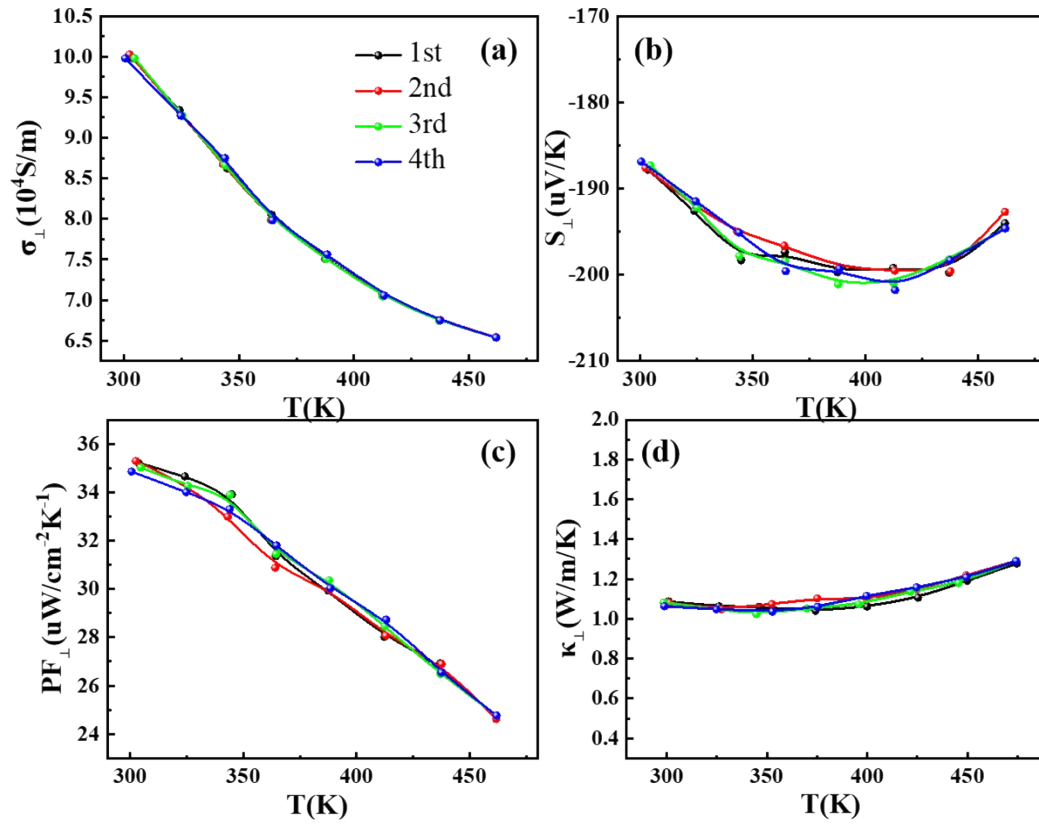
0.35 vol%).  $l$  is the average spatial distance between the Ag<sub>2</sub>Se nanoparticles in the BTS,  $\Delta D/D = 0.134$  and the average radius of Ag<sub>2</sub>Se (R) is 620nm; the pre-factors of point defect A, U- and N- process B, dislocation strain fields and dislocation cores C and D and the number density of nanoparticles  $N_p$  are provided in Table S2

**Table S2. Parameters used in Debye-Callaway model for BTS-xAg<sub>2</sub>Se samples (x=0, 0.1, 0.15, 0.25 and 0.35 vol%).**

	0vol%	0.10vol%	0.15vol%	0.25vol%	0.35vol%
A( $10^{-42}\text{S}^3$ )	5.45	5.45	5.45	5.45	5.45
B( $10^{-17}\text{Sk}^{-1}$ )	3.23	3.23	3.23	3.23	3.23
C	0	0.006	0.006	0.006	0.006
D( $10^{-30}\text{s}^2$ )	0	2.8	2.8	2.8	2.8
$l$	--	599	599	671	711
$N_p(10^{17}\text{m}^{-3})$	--	2.9	3	0.91	0.67

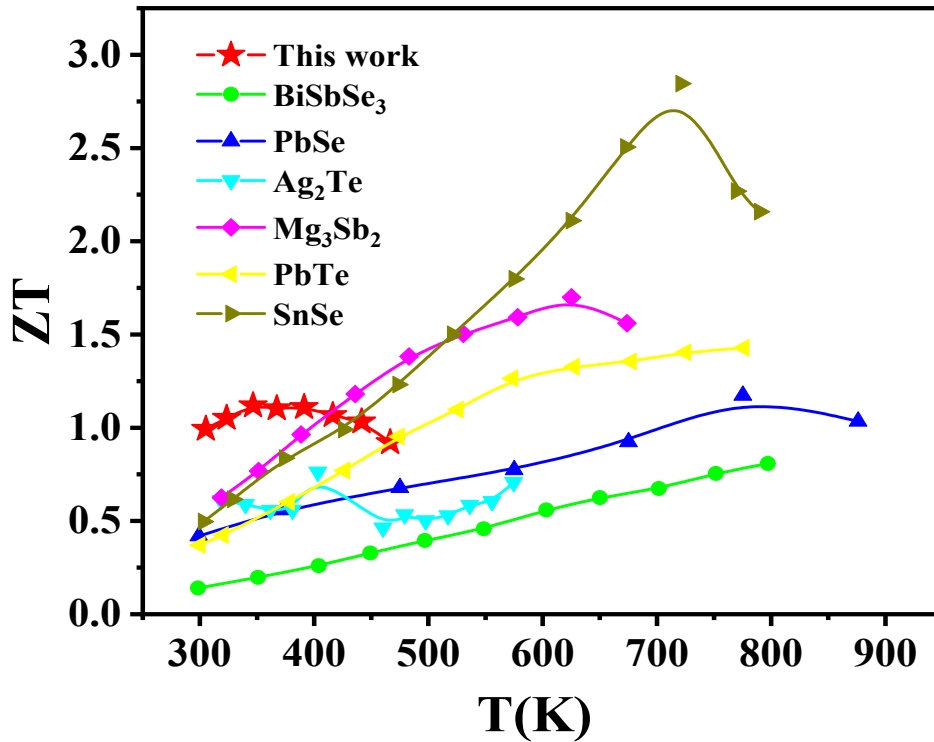


**Fig. S7. Calculated  $ZT_{\perp}$  as a function of  $n$  by a SPB model, the experimental  $ZT$  values are provided for comparison.**



**Fig. S8. Data repeatability of electrical conductivity (a), Seebeck coefficient (b), power factor (c) and thermal conductivity (d) normal to the SPS direction for BTS-0.15vol%Ag<sub>2</sub>Se sample**





**Fig. S9.** The temperature dependence of ZT for some typical N-type thermoelectric materials<sup>5-10</sup>.

From the Fig.S9, we can see the BTS-0.15vol%Ag<sub>2</sub>Se material has superior low-temperature thermoelectric performance in the low temperature region, which is better than other n-type systems.

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