Supplementary Information

Synergistic carrier and phonon transport advance Ag dynamically-doped *n*-type PbTe thermoelectrics *via* Mn alloying

Wei Yuan, ‡^a Qian Deng, ‡^a Dong Pan, ^b Xiang An, ^a Canyang Zhao, ^a Wenjun Su, ^a Zhengmin He, ^a Qiang Sun^{*c,d} and Ran Ang^{*a,e}

^a Key Laboratory of Radiation Physics and Technology, Ministry of Education, Institute of Nuclear Science and Technology, Sichuan University, Chengdu 610064, China

^b State Key Laboratory of Superlattices and Microstructures, Institute of Semiconductors, Chinese Academy of Sciences, P.O. Box 912, Beijing 100083, China

^c State Key Laboratory of Oral Diseases, National Clinical Research Center for Oral Diseases, West China Hospital of Stomatology, Sichuan University, Chengdu, Sichuan 610041, China

^d Sichuan Provincial Engineering Research Center of Oral Biomaterials, Chengdu, Sichuan 610041, China

^e Institute of New Energy and Low-Carbon Technology, Sichuan University, Chengdu 610065, China

‡ W. Yuan and Q. Deng contributed equally to this work.

*Corresponding authors and Emails: <u>rang@scu.edu.cn (RA), qiangsun@scu.edu.cn</u> (QS)

1. Computational details

1.1 Single parabolic band model

The Pisarenko curves and the effective mass of *n*-type PbTe can be modeled by the single parabolic band (SPB) model, and the SPB model is based on the following equations:¹ Seebeck coefficient:

$$S = \pm \frac{k_B}{e} \left(\frac{(r+5/2)F_{r+2/3}(\eta)}{(r+3/2)F_{r+1/2}(\eta)} - \eta \right)$$
(eq. S1)

Hall carrier concentration:

$$n_{H} = 4\pi \left[\frac{2m^{*}k_{B}T}{h^{2}}\right]^{3/2} F_{1/2}$$
 (eq. S2)

Lorenz number:

$$L = \frac{k_B^2 3F_0 F_2 - 4F_1^2}{e^2 F_0^2}$$
(eq. S3)

where $k_{\rm B}$ is the Boltzmann constant, \hbar is the reduce Planck constant, η is the reduced Fermi level, and *r* denotes the scattering factor and equals -1/2 here assuming that the acoustic scattering mechanism dominates. The Fermi integral is given by:

$$F_n(\eta) = \int_0^\infty \frac{x^n}{1 + e^{x-\eta}} dx$$
 (eq. S4)

1.2 Single Kane band model

The thermoelectric-transport properties of *n*-type PbTe can be modeled by adopting the single Kane band (SKB) model. The transport coefficients are determined by the equations as follows:^{2, 3}

Hall carrier concentration:

$$n_{H} = \frac{1}{eR_{H}} = A^{-1} \frac{N_{V} (2m_{b}^{*} k_{B}T)^{3/2}}{3\pi^{2} \hbar^{3}} F_{0}^{3/2}$$
(eq. S5)

Hall factor:

$$A = \frac{3K(K+2)^{0} F_{-4}^{1/20} F_{0}^{3/2}}{(2K+1)^{2} ({}^{0}F_{-2}^{1})^{2}}$$
(eq. S6)

Hall Carrier mobility:

$$\mu_{H} = A \frac{2\pi\hbar^{4}eC_{l}}{m_{l}^{*} (2m_{b}^{*}k_{B}T)^{3/2} E_{def}^{-2} F_{0}^{3/2}}$$
(eq. S7)

Power factor:

$$PF = \frac{2N_V \hbar k_B^2 C_l 1}{\pi E_{def}^2 m_I^*} \left(\frac{{}^{1}F_{-2}^1}{{}^{0}F_{-2}^1} - \xi \right)^{20} F_{-2}^1$$
(eq. S8)

 ${}^{n}F_{k}^{m}$ has a similar form as the Fermi integral:

$${}^{n}F_{k}^{m} = \int_{0}^{\infty} \left(-\frac{\partial f}{\partial\varepsilon}\right) \varepsilon^{n} (\varepsilon + \alpha\varepsilon^{2})^{m} [(1 + 2\alpha\varepsilon)^{2} + 2]^{k/2} d\varepsilon \qquad (eq. S9)$$

Where ξ is the reduced chemical potential, N_v is the band degeneracy (N_v =4 for *n*-type PbTe),⁴ *K* is the band anisotropy (*K*=3.6 for *n*-type PbTe),⁵ k_B is the Boltzmann constant, \hbar is the reduce Planck constant, m_I^* is the inertial mass ($m_I^* = 3m_d^*/(N_V^{2/3}(2K^{1/3} + K^{-2/3})))$,⁶ C_{lis} the longitudinal elastic ($C_{l=7.1\times1010}$ Pa for *n*-type PbTe),^{7, 8} ε is the reduced energy of electronic state, α is the reciprocal reduced band gap ($\alpha = k_B T/E_g$), E_g is the band gap , $E_g = 0.18eV + 0.0004 eV/K \times T$, E_{def} is the deformation potential coefficient.

1.3 Calculation of lattice thermal conductivity κ_{lat}

The lattice thermal conductivity (κ_{lat}) of the Ag_{0.03}Pb_{1-x}Mn_xTe samples was calculated by the modified Debye-Callaway model, which can be expressed by:

$$\kappa_{lat} = \frac{k_B}{2\pi^2 \upsilon} \left(\frac{k_B}{\hbar}\right)^3 \int_0^{\Theta/T} \tau_{tot}(x) \frac{x^4 e^x}{(e^x - 1)^2} dx \qquad (eq. S10)$$

Here, $v = 3^{1/3}(v_l^{-3} + 2v_t^{-3})^{-1/3}$ is the average speed of phonon, v_l and v_t denote the longitudinal and transverse speeds of sound respectively, \hbar is the reduced Planck constant, Θ is the Debye temperature, *x* is the relation of $\hbar \omega / k_B T$, ω is the phonon frequency, and τ_{tot} is the total phonon scattering relaxation time. The total phonon relaxation time τ_{tot} for the Ag_{0.03}Pb_{1-x}Mn_xTe samples consists of Umklapp process, normal process, point defects, grain boundary, and nanoprecipitates (U+N+PD+GB+NP) by the following equation:⁹

$$\tau_{tot}^{-1} = \tau_U^{-1} + \tau_N^{-1} + \tau_{PD}^{-1} + \tau_{GB}^{-1} + \tau_{NP}^{-1}$$
(eq. S11)

The phonon scattering relaxation time for respective mechanism can be expressed as follows: Umklapp phonon scattering:

$$\tau_{U}^{-1} = A_{N} * \frac{2 k_{B} \bar{V}^{1/3} \omega^{2} \gamma^{2} T}{(6\pi^{2})^{\frac{1}{3}} \bar{M} \upsilon^{3}}$$
(eq. S12)

Normal phonon scattering:

$$\tau_N^{-1} = \frac{2 K_B V^{1/3} \omega^2 \gamma^2 T}{(6\pi^2)^3}$$
(eq. S13)

Point defect scattering:

$$\tau_{PD}^{-1} = \frac{V_{\omega}^{4}}{4\pi^{3}v^{3}} * \sum (1-x_{i}) \left[(\frac{M_{i}-M}{M})^{2} + \varepsilon (\frac{a_{i}-a}{a})^{2} \right]$$
(eq. S14)

Grain boundary scattering:

$$\tau_{GB}^{-1} = \nu/Gd \tag{eq. S15}$$

Nanoprecipitates scattering:

$$\tau_{NP}^{-1} = \upsilon [(2\pi R^2)^{-1} + (\frac{4}{9}\pi R^2 (\Delta D/D)^2 (\frac{\omega R}{\upsilon})^4)^{-1}]^{-1} N_p \qquad (\text{eq. S16})$$

In above equations, \overline{V} is the average atomic volume, \overline{M} is the average atomic mass, γ is the Grüneisen parameter, A_N is the ratio between normal process and Umklapp phonon scattering.

1.4 Calculation of quality factor **B**

The quality factor *B* of the $Ag_{0.03}Pb_{1-x}Mn_xTe$ samples was calculated by:¹⁰

$$B = \left(\frac{k_B}{e}\right)^2 \frac{8\pi e (2m_e k_B T)^{3/2} \mu_W}{3h^3 \kappa_{lat}} T$$
(eq. S17)

2. Supplementary Figures



Figure S1. ZT of the main components of dynamic doped PbTe systems.¹¹⁻¹⁶



Figure S2. κ_{lat} at 323 K of the main components of dynamic doped PbTe systems.¹¹⁻¹⁶ the Lorenz number calculation methods for representative samples of each system: single parabolic band (SPB) model: PbTe:Cr/Ag2Te, PbTe:Ag, PbTe:Cu/MnTe, PbTe:Ga; single Kane band (SKB) model: PbTe:In/I, PbTe:Cu.



Figure S3. EDS spectra obtained from Ag₂Te secondary phase.



Figure S4. The Hall carrier concentration $n_{\rm H}$ and the Hall carrier mobility $\mu_{\rm H}$ as a function of x in Ag_{0.03}Pb_{1-x}Mn_xTe.



Figure S5. DSC curve of $Ag_{0.03}Pb_{0.95}Mn_{0.05}Te$, and the peak occurs at 593 K



Figure S6. Temperature-dependent effective masses m^* of PbTe and Ag_{0.03}Pb_{1-x}Mn_xTe (x = 0.02, 0.04, 0.05), the black scatter values are from the reported I doped PbTe material.⁶



Figure S7. Temperature-dependent power factor *PF* for *n*-type $Ag_{0.03}Pb_{1-x}Mn_xTe$ samples.



Figure S8. Temperature-dependent electrical thermal conductivity κ_e for *n*-type Ag_{0.03}Pb_{1-x}Mn_xTe samples.



Figure S9. Temperature-dependent thermal diffusivity (*D*) for *n*-type $Ag_{0.03}Pb_{1-x}Mn_xTe$ samples.



Figure S10. Temperature-dependent heat capacity (C_p) for *n*-type Ag_{0.03}Pb_{1-x}Mn_xTe samples.



Figure S11. The calculated quality factor *B* at 323 K, 523 K, 623 K and 773 K for *n*-type $Ag_{0.03}Pb_{1-x}Mn_xTe$ samples.



Figure S12. Thermal stability and reproducibility of (a) electrical conductivity σ , (b) Seebeck coefficient *S*, (c) total thermal conductivity κ_{tot} , and (d) figure of merit *ZT* for the high-performance *n*-type Ag_{0.03}Pb_{0.95}Mn_{0.05}Te sample.



Figure S13. The average $ZT(ZT_{avg})$ for $Ag_{0.03}Pb_{1-x}Mn_xTe$ samples.

3. Supplementary tables

Table S1. The sample densities of $Ag_{0.03}Pb_{1-x}Mn_xTe$ (x = 0, 0.01, 0.02, 0.03, 0.04, 0.05, 0.06) measured by Archimedes method.

Sample No.	Measured density (g cm ⁻³)	Relative density (%)
<i>x</i> =0	8.17	99.15
<i>x</i> =0.01	8.099	98.29
<i>x</i> =0.02	8.098	98.28
<i>x</i> =0.03	8.003	97.12
<i>x</i> =0.04	8.087	98.14
<i>x</i> =0.05	8.037	97.53
<i>x</i> =0.06	8.167	99.11

Parameters	Symbol	Value	Ref.
Gruneisen parameter	γ	1.65	3
Debye temperature	$\theta_{\rm D}({\rm K})$	163	17
Average sound velocity	v (m/s)	1720	17
Average atomic mass for PbTe	$M(\mathrm{kg})$	2.78×10 ⁻²⁵	18
Ratio of N- to U- process	$A_{ m N}$	2.9	19
Phenomenological parameter		65	18
Mole mass of Pb	$M_{ m Pb} ({ m g/mol})$	207.2	-
Mole mass of Te	$M_{\rm Te}$ (g/mol)	127.6	-
Mole mass of Ag	$M_{\rm Ag}$ (g/mol)	107.8682	-
Mole mass of Mn	$M_{ m Mn}$ (g/mol)	54.938	-
Radius of Pb atom	$r_{\mathrm{Pb}}(\mathrm{\AA})$	1.2	-
Radius of Ag atom	$r_{\mathrm{Ag}}(\mathrm{\AA})$	1.26	-
Radius of Mn atom	$r_{\rm Mn}$ (Å)	0.8	-

Table S2. The input parameters based on the theoretical simulation of lattice thermal conductivity in this work.

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