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# Thermodynamic model for voltammetric response in conducting redox polymers 

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## Supplementary

Developing the thermodynamic model for a conducting redox polymer film.
Here we consider systems consisting of current collector with pdoped CRP film deposited on it, this all is immersed in 1, 1 binary electrolyte solution (see schem.2).
Equilibrium processes (eqs.10-12) can be represented in the form of the following thermodynamic equations:

$$
\begin{align*}
& \mu_{R e d}(f)-\mu_{O x}(f)-\mu_{e}(s b)=0 \\
& m \cdot \mu_{R}(f)-\mu_{P}(f)-\mu_{e}(s b)=0 \\
& \mu_{R e d}(f)+\mu_{P}(f)-\mu_{O x}(f) m \cdot \mu_{R}(f)=0
\end{align*}
$$

where $\tilde{\mu}_{i}(j)$ is the electrochemical potential of $i$ species within $j$ phase; subscript $P$ stays for polarons; subscript $R$ stays for reduced repeating units of the polymer backbone; subscript $O x$ denotes oxidized redox centers of pendants; subscript Red denotes reduced redox centers of pendants. Above equations (Eqs. 1s, 2s, 3s) can be rewritten in the more detailed form, which is given below for eq. 3 s :

$$
\begin{aligned}
& \mu_{R e d}^{\cup}(f)+R T \ln \left(\vartheta_{\text {Red }}\right)-\mu_{O x}^{\cup}(f)-R T \ln \left(\vartheta_{O x}\right)- \\
& -F \Phi-\mu_{e}^{0}(s b)+F E=0 ; \\
& \left(\mu_{\text {Red }}^{\cup}(f)-\mu_{o x}^{\cup}(f)-\mu_{e}^{\cup}(s b)\right)+R T \ln \left(\vartheta_{R e d} / \vartheta_{O x}\right)+ \\
& +F(E-\Phi)=0 ; \\
& \frac{R T}{}+\ln \left(\overline{\vartheta_{O x}}\right)^{+} \\
& +\frac{(E-\Phi) F}{R T}=0 ;
\end{aligned}
$$

Here we used definitions of the electrochemical potentials of mobile and immobile species ${ }^{i}$ within the phas ${ }^{\beta}{ }^{k}$, which are in accordance with previous works. ${ }^{70,75,90}$ Here $\mu_{i}$ is the standard

[^0]chemical potential of $i$ species in the film phase; $\vartheta_{O x}$ is, as previously, the fraction of oxidized redox centers; $\vartheta^{\text {Red }}$ is the fraction of reduced redox centers. Exponentiation of the final version of eq. 4 s and use of dimensionless electric potentials $\Phi=\Phi F / R T$ and $\mathrm{E}^{\prime}=E F / R T$, finally gives us following expression for $\vartheta_{O x}$ :
$$
K_{\text {Red } O x} \cdot \exp (E-\Phi)=\overline{1-\vartheta_{O x}} ;
$$
where $K_{\text {RedOx }}$ is the dimensionless equilibrium constant of the process of eq.10.
In the same manner as eq.1s resulted in eq.5s, equations (7, 8, 2 s ) can be rewritten in the more convenient forms:
\[

$$
\begin{align*}
& C_{A}=K_{A} C_{0} \cdot \exp (\Phi) \\
& C_{K}=K_{K} C_{0} \cdot \exp (-\Phi) \\
& \wedge_{P} \cdot \exp \left(L-\Psi ノ-\overline{\left(1-\vartheta_{P}\right)^{m^{\prime}}}\right.
\end{align*}
$$
\]

where $K_{K}$ is the equilibrium constant of the process of cation distribution between the phases of solution and polymer film (eq.6). $K_{A}$ is the equilibrium constant of the process of counterion distribution between the phases of solution and polymer film (eq.5). $K_{P}$ is the equilibrium constant of the process of electron transfer between the current collector and the polymer backbone (eq.11).
From equations 5 s and 8 s one can obtain thermodynamic description of the process of charge exchange between the redox-active pendants and the backbone (eq.12), which takes the following form:

$$
\begin{equation*}
\mathbf{N}_{\text {RedOx }} \cdot \overline{\left(1-\vartheta_{P}\right)^{m}}-\boldsymbol{\Lambda}_{P} \cdot \overline{1-\vartheta_{O x}} \tag{9s}
\end{equation*}
$$

Here we underline once again that polymer film is considered as distinct thermodynamic phase, which is at equilibrium with the other phases of the system. Thus, the electric potential within the film phase $\Phi$ differs from the current collector potential $E$ and differs from the electrolyte solution potential (latter set to be zero in this work).
To obtain the quasi-equilibrium voltammetric curves corresponding to the applied approach, we have to consider the electroneutrality condition (eq.3, 4); it means that the charge
carried out from the film through the current collector/film interface must be compensated by the charge $Q=F A L\left(C_{A}-C_{K}\right)$ injected into the film through the film/solution border. Therefore, the current I passing through the current collector/film interface can be expressed in the following way:

$$
\begin{equation*}
I=\frac{w_{q}}{d t}=F A L \frac{\cdot . \cdots}{d t}=\frac{\cdots \cdots}{R T \quad . .} \tag{10}
\end{equation*}
$$

where $v$ is the potential sweep rate $[\mathrm{V} / \mathrm{s}] ; A$ is the polymer film geometric surface area $\left[\mathrm{m}^{2}\right] ; t$ is time [s]; $E^{\prime}$ is, as previously, dimensionless electrode potential $\left(E^{\prime}=E F / R T\right)$.
From eq.9s one can express the fraction of oxidized pendants $\vartheta_{O x}$ through the fraction of oxidized repeating units of the organic conjugated backbone $\vartheta_{P}$ :

$$
\begin{equation*}
\overline{K_{\text {RedOx }}}\left(1-\vartheta_{P}\right)^{\prime \prime \prime}+\vartheta_{P} \tag{11}
\end{equation*}
$$

Next, one can express the film-bulk potential $\Phi^{\prime}$ via the fraction of oxidized repeating units of the backbone ${ }^{\vartheta_{P}}$. Substituting expressions for anion film-bulk concentration $C_{A}$ (eq.6s), cation film-bulk concentration ${ }^{C}{ }_{K}$ (eq.7s) and the fraction of oxidized redox centers $\vartheta_{O x}$ (eq.11s) into the electroneutrality condition (eq.4) one gets the following form of the electroneutrality condition:

$$
\begin{equation*}
\left.K_{A}{ }_{N}^{C} C^{0} \cdot \left\lvert\, \frac{\exp \left(\Phi^{\prime}\right)-\vartheta_{K} K_{K} C_{0} \cdot \exp \left(-\frac{\left.\Phi^{9}\right)}{\frac{K_{P}}{K_{\text {RedOx }}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}}\right.}{n \cdot m}\right.\right] \tag{12}
\end{equation*}
$$

Multiplication of the both sides of eq.12s by $\exp \left(\Phi^{\prime}\right) / K_{A}$ leads to the following result:

$$
=\frac{N}{K_{A} C_{0}} \cdot\left[\frac{\exp ^{\prime}(2 \Phi)}{\left[\frac{\vartheta_{P}}{K_{R}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}\right.}+\frac{\vartheta_{P}}{n \cdot m}\right] \exp \left(\Phi^{\prime}\right)
$$

Eq.13s is the quadratic equation for $\exp \left(\Phi^{\prime}\right)$, thus, it can be simply turned into the expression of $\exp \left(\Phi^{\prime}\right)$ vs. ${ }^{\vartheta_{P}}$ :

$$
\begin{equation*}
=\frac{2 \cdot \exp (\Phi)}{K_{A} C_{0}} \cdot\left[\frac{\vartheta_{P}}{\frac{K_{P}}{K_{\text {RedOx }}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}}+\frac{\vartheta_{P}}{n \cdot m}\right]+ \tag{14}
\end{equation*}
$$



Figure 1s The film bulk potential $\Phi$ dependence on the electrode potential $E$. Parameters are following: $n=1 ; m=2 ; N=2$ mol/l; $K_{A}=0.1$; $C_{0}=1_{\text {mol/l. Black line }} K_{K}=0.001_{\text {; Red line }} K_{K}=0.01$.

On the other hand, the maximum value of the film-bulk potential corresponds to the completely charged state of the film. In such a state, co-ion (cation) concentration inside the film is negligible. Thus, anion concentration must reach its maximum value $C_{\mathrm{A}, M A X}$ and be equal to the maximum concentration of charge of immobile charge carriers in the film $(N+M / m)$. Thus, parameters such as $N, m, n=M / N_{\text {influence the }} \Phi_{M A X}$. as they define the maximum concentration of charge of immobile charge carriers in the film.
By now we expressed all necessary variables via $\vartheta_{P}$, thus, we can continue with derivation of quasi-equilibrium current equation. Substituting $\left(C_{A}-C_{K}\right)$ in the expression for quasiequilibrium current (eq.10s) with the right side of the eq.12s, one gets the following form of the expression for quasiequilibrium current:

$$
\begin{equation*}
I=\frac{v F^{2} A L}{R T} ; \tag{17}
\end{equation*}
$$

S

One can take into account that $d \vartheta_{P} / d E=\left(d E / d \vartheta_{P}\right)^{-1}$ and
$d E / d \vartheta_{P}$ $d E / d \vartheta_{P}$ can be expressed from eq. 8 s via $\vartheta_{P}, \Phi$ and $d \Phi / d \vartheta_{P}$. Therefore, one can finally derive the analytical expression for the quasi-equilibrium current $I$ vs. $\vartheta_{P}$. Unfortunately, the derivative $d E / d \vartheta_{P}$ is so big, that makes it completely inappropriate to be included even in supplementary. So, we will leave the final expression for current in the following form:

$$
\begin{equation*}
\left.\left.+\frac{R T}{N \vartheta_{p}{ }^{2} n m\left(\left(\frac{K_{P}}{K_{R} K_{d o x}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}\right)\right.} \frac{n m}{K_{\text {Redox }}}\left(1-\vartheta_{P}\right)^{m-1}-1\right)\right) \left\lvert\, \div \frac{K_{P}}{\left[\left(\frac{K_{1}}{K_{\text {Redox }}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}\right) n m\right]^{2}}\right. ; \tag{18}
\end{equation*}
$$

Substituting expression for $\exp (\Phi)$ (eq.14s) into eq. 8 s gives the analytical expression for $E_{\text {vs. }} \vartheta{ }^{\text {. }}$

$$
\begin{equation*}
\cdot \frac{\left[\frac{N}{2 K_{A} C_{0}}\left(\frac{\sqrt{K_{P}\left(1-\vartheta_{H}\right)_{P}^{m}}}{\frac{K_{P}}{K_{\text {RedOx }}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}}+\frac{\vartheta_{P}}{n m}\right)+\frac{1}{2}\right.}{} \cdot \sqrt{\left.4 \frac{K_{K}}{K_{A}}+\left[\frac{N}{K_{A} C_{0}}\left(\frac{\vartheta_{P}}{\frac{K_{P}}{K_{\text {RedOx }}}\left(1-\vartheta_{P}\right)^{m}+\vartheta_{P}}+\frac{\vartheta_{P}}{n m}\right)\right]\right]^{2}} . \tag{19}
\end{equation*}
$$

Thus, quasi-equilibrium voltammetric curves (/ vs. E' plots) can be calculated with the following list of parameters: $N, C_{0}, n, m$, $K_{K}, K_{A}, v, A, L, T, K_{P}, K_{\text {Red } O x}$.


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