

## Electronic Supplementary Information

# Mechanism of poly(*N*-isopropylacrylamide) cononsolvency in aqueous methanol solutions explored *via* oxygen K-edge X-ray absorption spectroscopy

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## Experimental and Theoretical Methods

### *Soft X-ray absorption spectroscopy*

Soft X-ray absorption spectroscopy (XAS) experiments were performed at the soft X-ray beamline BL-7A of the Photon Factory, Institute of Materials Structure Science, High Energy Accelerator Research Organization (KEK-PF).<sup>1</sup> The energy resolution  $\Delta E$  of soft X-rays was set to 0.29 eV at the O K-edge. The details of the transmission-type liquid cell for XAS have been described previously.<sup>2, 3</sup> The liquid cell was placed at ambient pressure conditions of helium gas, where a liquid layer was sandwiched between two Si<sub>3</sub>N<sub>4</sub> membranes with a thickness of 100 nm. The thickness of the liquid layer was controlled by adjusting the helium pressure around the liquid cell. The beam size of soft X-rays was  $200 \times 200 \mu\text{m}^2$  owing to the window size of the Si<sub>3</sub>N<sub>4</sub> membrane, which separates the chamber of the liquid cell in the atmospheric helium condition and the soft X-ray beamline under an ultrahigh vacuum condition. The XAS spectra were obtained using the Beer-Lambert law,  $\ln(I_0/I)$ , where  $I_0$  and  $I$  are the transmission signals of the bare Si<sub>3</sub>N<sub>4</sub> membranes and liquid samples confined by the Si<sub>3</sub>N<sub>4</sub> membranes, respectively. Liquid samples were exchanged using a syringe pump. The temperature of the liquid samples was controlled at 25 °C. The photon energies were calibrated precisely by measuring the XAS spectra of the polymer film before and after sample measurements.<sup>4</sup>

Poly(*N*-isopropylacrylamide) (PNIPAM) with Mn ~ 40,000 was purchased from Sigma-Aldrich. Methanol (MeOH) was purchased from FUJIFILM Wako Pure Chemical Co. Ultrapure water (H<sub>2</sub>O) with a resistivity above 18 M $\Omega$  was used in our experiments. PNIPAM with a concentration of 50 mg/mL was dissolved in aqueous MeOH solutions (MeOH)<sub>*x*</sub>(H<sub>2</sub>O)<sub>1-*x*</sub> with different MeOH molar fractions. As shown in Sec. S2, PNIPAM is insoluble at  $0.4 > x > 0.1$ , a behavior known as cononsolvency.

### *Molecular Dynamics simulation*

All-atom molecular dynamics (MD) simulations were performed using the GROMACS 2021.3 package.<sup>5</sup> The equation of motion was integrated using the Leapfrog algorithm with a time step of 2 fs. Under each condition, the initial conformation was first energy-minimized using a steep descent algorithm. A 2 ns equilibration was followed by a 1  $\mu\text{s}$  production run at 300 K and 0.1 MPa, both under an isothermal-isobaric (*NpT*) ensemble. Temperature and pressure for production runs were controlled using a Nose-Hoover thermostat and Parrinello-Rahman barostat with dumping constants being 0.5 ps and 1 ps, respectively, whereas Berendsen algorithm with a dumping constant of 2 ps was used for equilibration.

A single atactic (meso 47%) 40-mer PNIPAM chain was dissolved in 10,000 solvent molecules of MeOH and H<sub>2</sub>O. A cubic simulation box was used to apply periodic boundary conditions in all directions. PNIPAM was represented by the OPLS-AA model.<sup>6</sup> The default partial charges were multiplied by a factor of 1.31, as proposed in previous studies.<sup>7, 8</sup> H<sub>2</sub>O and MeOH molecules were represented by the SPC/E model<sup>9</sup> and KBFF model,<sup>10</sup> respectively. The nonbonded interactions were truncated at 1.4 nm without dispersion corrections. Long-range Coulombic interactions were evaluated using the particle-mesh Ewald algorithm. The Lennard-Jones parameters between different particles were determined using the Lorentz-Berthelot

combination rules:

$$\sigma_{ij} = (\sigma_{ii} + \sigma_{jj})/2, \epsilon_{ij} = \sqrt{\epsilon_{ii}\epsilon_{jj}} \quad (1)$$

LINCS constraints are applied for all bonds.

### ***Inner-shell calculation***

Inner-shell calculations were conducted using the program package GSCF3.<sup>11, 12</sup> The ground and core excited states were calculated through the Hartree-Fock method, namely,  $\Delta$ SCF (self-consistent field). The core-hole was frozen on a specified O atom in the SCF calculations for the core excited state. The relaxed Hartree-Fock potential for the O 1s ionized state was obtained by a partial SCF calculation within the orbital manifold orthogonalized to the valence excited state, whereas the Rydberg excited states were obtained using the improved virtual orbital method<sup>13</sup> to avoid spurious mixing between the valence and Rydberg orbitals. The present calculations do not include the zero-point vibrational energy. The scalar relativistic effect was not considered in the present calculation. The contracted Gaussian-type functions by Huzinaga et al. were used as primitive basis functions: (73/7) for C, N, and O, and (6) for H.<sup>14</sup> The contraction schemes were (3111121/3112/1\*) for C, N, and O, and (42) for H atoms. The d-type polarization functions were used at the C atoms ( $\zeta_d = 0.600$ ), N atoms ( $\zeta_d = 0.864$ ), and O atoms ( $\zeta_d = 1.154$ ).

The structural optimizations of *N*-isopropylacrylamide (NIPAM) and PNIPAM, which consists of five polymer units of NIPAM, were performed using Gaussian 16.<sup>15</sup> The structural optimization was performed with B3LYP/aug-cc-pVDZ, considering the previous result for the formamide-H<sub>2</sub>O complex.<sup>16</sup> The structures of MeOH and H<sub>2</sub>O were also optimized with B3LYP/aug-cc-pVDZ. The hydrogen bond (HB) models of NIPAM and PNIPAM with MeOH and H<sub>2</sub>O were determined from the results of the two-dimensional (2D) radial distribution function (RDF) obtained by MD simulations, as shown in Sec. S6.

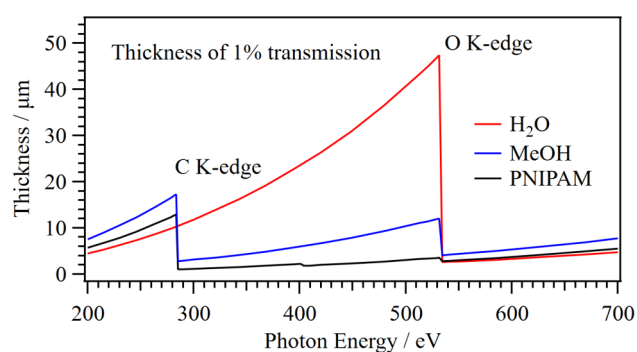
## Results and Discussion

### S1. Detection Limit of PNIPAM in XAS Spectra

In the 50 mg/mL PNIPAM solutions, the concentration of the polymer unit is about 440 mM. The detection limit of the C=O  $\pi^*$  peaks (532 eV) in the O K-edge XAS is 100 mM for separating the contribution of solvent MeOH and H<sub>2</sub>O, whose first peaks are around 535 eV. Note that metal oxides have peaks below 530 eV.<sup>17</sup> In this case, the detection limit of the metal oxide peaks becomes lower and closer to 10 mM because the metal oxide peaks are far away from the absorbance of solvent MeOH and H<sub>2</sub>O. The detection limit of C and N K-edge XAS in solvent H<sub>2</sub>O is also low and close to 10 mM because these energy regions are well known as the water window.<sup>18</sup> In the solvent MeOH, on the other hand, it is difficult to measure C K-edge XAS spectra of solute molecules because the first peak of MeOH at the C K-edge is around 288 eV.<sup>19</sup> The C=C  $\pi^*$  peaks (285 eV) such as benzene<sup>20</sup> and the C=N  $\pi^*$  peaks (286 eV) such as pyridine<sup>21</sup> can be observed in the solvent MeOH because these peaks are below the first peak of MeOH. Because the energetic position of the C=O  $\pi^*$  peaks are around 290 eV, the C=O  $\pi^*$  peaks cannot be observed in the solvent MeOH. Therefore, PNIPAM in solvent MeOH only shows the C=O  $\pi^*$  peaks in the O K-edge XAS, whose detection limit is above 100 mM. The detection limit of N K-edge XAS also becomes higher in the solvent MeOH. Soft X-rays at the N K-edge (400 eV) are also absorbed by solvent MeOH because the photon energy of the N K-edge is close to that of the C K-edge. In the soft X-ray beamline, monochromatic soft X-rays include not only first-order X-rays (400 eV) but also the high-order X-rays (800 eV, 1200 eV, *etc.*) due to the high order diffraction of a plane grating monochromator. Because of the absorption of the first-order X-rays (400 eV) by the solvent MeOH, the ratio of the high-order X-rays is increased compared to the first-order X-rays, resulting the increase of the detection limit of the N K-edge XAS.

As shown in Sec. S2, PNIPAM in aqueous MeOH solutions (MeOH)<sub>x</sub>(H<sub>2</sub>O)<sub>1-x</sub> at  $0.4 > x > 0.1$  becomes a cloud solution because PNIPAM is insoluble in these aqueous MeOH solutions. In these cases, PNIPAM aggregates are floating in aqueous MeOH solutions. For exploring the measurement conditions of the O K-edge XAS of PNIPAM in aqueous MeOH solutions, the thickness of the samples for 1% transmission of soft X-rays at different photon energy was calculated from the previous theoretical results,<sup>22</sup> as shown in Fig. S1. In the calculations, the density of liquid H<sub>2</sub>O was 1.0 g/cm<sup>3</sup>, that of liquid MeOH was 0.792 g/cm<sup>3</sup>, and that of solid PNIPAM was 1.386 g/cm<sup>3</sup>. The thickness of liquid H<sub>2</sub>O for 1% transmission of soft X-rays was 47.4  $\mu\text{m}$  before the absorbance of H<sub>2</sub>O and was 2.6  $\mu\text{m}$  after the absorbance of H<sub>2</sub>O at the O K-edge. The O K-edge XAS measurement of liquid H<sub>2</sub>O needs the liquid layer below 1  $\mu\text{m}$ . In the energy region of the C=O  $\pi^*$  peaks of PNIPAM, on the other hand, soft X-ray can transmit the liquid layer with the thickness of 40  $\mu\text{m}$ . The thick liquid layer should be prepared for the increase of the absorbance of the C=O  $\pi^*$  peaks of PNIPAM. The thickness of liquid MeOH for 1% transmission of soft X-rays was 12.1  $\mu\text{m}$  before the absorbance of MeOH and was 4.1  $\mu\text{m}$  after the absorbance of MeOH at the O K-edge. As a result, the thickness of aqueous MeOH solution at  $0.4 > x > 0.1$  would be between 10–40  $\mu\text{m}$  for the transmission of soft X-rays in the energy region of the C=O  $\pi^*$  peaks of PNIPAM. The thickness of solid PNIPAM for 1%

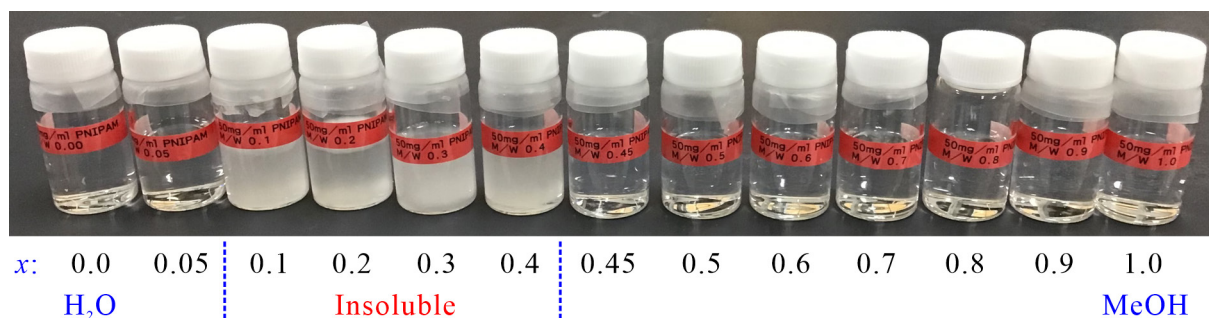
transmission of soft X-rays was  $3.5\ \mu\text{m}$  in the energy region of the  $\text{C}=\text{O}\ \pi^*$  peaks of PNIPAM. It means that soft X-rays cannot transmit PNIPAM aggregates when the sizes of the aggregates are above several  $\mu\text{m}$ . The  $\text{C}=\text{O}\ \pi^*$  peaks of PNIPAM are not observed when the concentrations of PNIPAM are below the detection limit (100 mM). Because the  $\text{C}=\text{O}\ \pi^*$  peaks of PNIPAM were not observed at  $0.4 > x > 0.1$  in the O K-edge XAS spectra of PNIPAM in aqueous MeOH solutions, the concentration of PNIPAM were below the detection limit and the sizes of the PNIPAM aggregates were above several  $\mu\text{m}$  in the aqueous MeOH solutions with the liquid thickness between  $10\text{--}40\ \mu\text{m}$ .



**Figure S1.** Thickness of liquid  $\text{H}_2\text{O}$ , liquid MeOH, and solid PNIPAM for 1% transmission of soft X-rays as a function of photon energy calculated from the previous theoretical results.<sup>22</sup>

## S2. Photos of PNIPAM in Aqueous MeOH Solutions

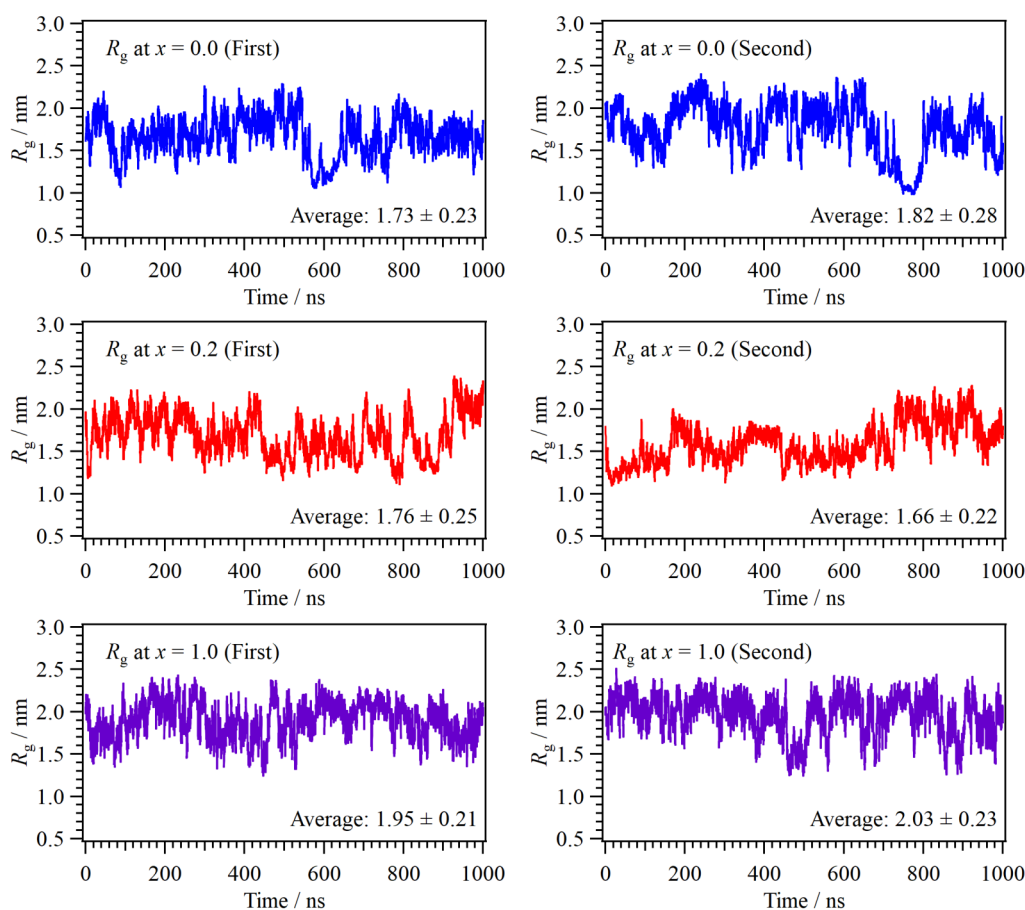
Figure S2 shows photos of PNIPAM in aqueous MeOH solutions at different molar fractions. The solvent in the left bottle is pure  $\text{H}_2\text{O}$  ( $x = 0.0$ ) and that in the right bottle is pure MeOH ( $x = 1.0$ ). PNIPAM is completely dissolved with the solvents in both the MeOH-rich region ( $x > 0.4$ ) and the  $\text{H}_2\text{O}$ -rich region ( $0.1 > x$ ). PNIPAM is insoluble at  $0.4 > x > 0.1$  and becomes a cloud solution. The deposition of PNIPAM is stronger by increasing the  $\text{H}_2\text{O}$  molar fraction and becomes a maximum value at  $x = 0.1$ .



**Figure S2.** Photos of PNIPAM in aqueous MeOH solutions  $(\text{MeOH})_x(\text{H}_2\text{O})_{1-x}$  at different molar fractions at a room temperature. PNIPAM is insoluble in the middle concentration region ( $0.4 > x > 0.1$ ), a behavior known as cononsolvency.

### S3. Time Evolution of Radius of Gyration

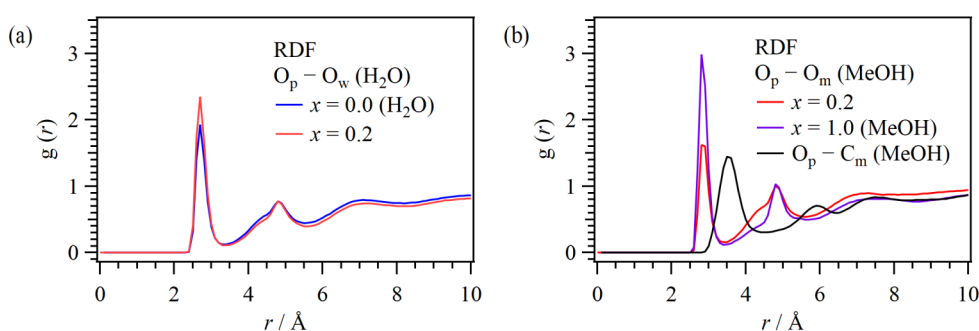
Figure S3 shows time evolution of radius of gyration ( $R_g$ ) in aqueous MeOH solutions at different molar fractions ( $x = 0.0, 0.2,$  and  $1.0$ ). Note that pure H<sub>2</sub>O ( $x = 0.0$ ) and pure MeOH ( $x = 1.0$ ) are good solvents for PNIPAM. We selected aqueous MeOH solution at  $x = 0.2$ , which represents the consolvency region ( $0.4 > x > 0.1$ ). Two independent trajectories were shown at each concentration. The average values of  $R_g$  at each trajectory were also shown. As averaging all the trajectories at each concentration, we found that  $R_g = 1.78 \pm 0.19$  at  $x = 0.0$ ,  $R_g = 1.71 \pm 0.17$  at  $x = 0.2$ , and  $R_g = 1.99 \pm 0.16$  at  $x = 1.0$ , respectively. The largest value of  $R_g$  at  $x = 1.0$  means that PNIPAM chains in pure MeOH spread due to the weak interactions between polymer units. The polymer units are close to each other due to the strong interactions at  $x = 0.0$ , resulting the decrease of  $R_g$ . The smallest value of  $R_g$  at  $x = 0.2$  means that the polymer units are aggregated and the deposition of PNIPAM is occurred.



**Figure S3.** Time evolution of radius of gyration ( $R_g$ ) of PNIPAM computed all atoms at  $x = 0.0, 0.2,$  and  $1.0$ . Two independent trajectories at each  $x$  are shown. Note that  $R_g$  were calculated from all the trajectories at different MeOH molar fractions.

#### S4. Coordination Numbers of C=O Group in PNIPAM

Figure S4(a) shows RDF of oxygen atom ( $O_p$ ) of the C=O group in PNIPAM with oxygen atom ( $O_w$ ) of H<sub>2</sub>O at  $x = 0.0$  (H<sub>2</sub>O) and  $x = 0.2$ . The position of the first coordination peak is 2.7 Å at  $x = 0.0$ , indicating the HB structures of the C=O group in PNIPAM with H<sub>2</sub>O. Since the position of the first coordination peak is not nearly changed at  $x = 0.2$ , the HB structure of PNIPAM with H<sub>2</sub>O at  $x = 0.2$  is nearly same as that at  $x = 0.0$ . Figure S4(b) shows RDF of  $O_p$  with oxygen atom ( $O_m$ ) of MeOH at  $x = 1.0$  (MeOH) and  $x = 0.2$ . The position of the first coordination peak is 2.8 Å at  $x = 1.0$ , indicating the HB structure of C=O group in PNIPAM with MeOH. The longer distance of the first coordination peak (3.5 Å) in RDF of  $O_p$  with carbon atom ( $C_m$ ) of MeOH also supports the HB structure of PNIPAM with MeOH. Since the position of the first coordination peak in RDF  $O_p - O_m$  at  $x = 0.2$  is nearly close to that at  $x = 1.0$ , the HB structure of PNIPAM with MeOH at  $x = 0.2$  is not changed compared to that at  $x = 1.0$ .



**Figure S4.** RDF of oxygen atom ( $O_p$ ) of PNIPAM with oxygen atom ( $O_w$ ) of H<sub>2</sub>O and oxygen atom ( $O_m$ ) of MeOH. (a)  $O_p - O_w$  in  $x = 0.0$  (H<sub>2</sub>O) and  $x = 0.2$ . (b)  $O_p - O_m$  in  $x = 0.2$  and  $x = 1.0$  (MeOH). RDF between  $O_p$  and carbon atom ( $C_m$ ) of MeOH at  $x = 1.0$  is also shown.

**Table S1.** Probability distributions of coordination numbers of  $O_p$  with  $O_w$  and  $O_m$  at  $x = 0.0$ ,  $x = 0.2$ , and  $x = 1.0$ .

	0	1	2	3
$O_p - O_w$ at $x = 0.0$	0.01	0.09	0.72	0.18
$O_p - O_w$ at $x = 0.2$	0.04	0.34	0.55	0.07
$O_p - O_m$ at $x = 0.2$	0.69	0.28	0.03	0.0
$O_p - O_m$ at $x = 1.0$	0.05	0.58	0.36	0.01

From the first coordination peaks in RDF, we have determined the probability distributions of coordination numbers of MeOH and H<sub>2</sub>O with the C=O group of PNIPAM, which are summarized in Table S1. In pure MeOH ( $x = 1.0$ ), one or two MeOH molecules form HB structures with the C=O group of PNIPAM. The probability of one MeOH coordination is larger than that of two MeOH coordination. In the MeOH-H<sub>2</sub>O

mixtures at  $x = 0.2$ , one or two H<sub>2</sub>O molecules form HB structures with C=O group of PNIPAM, whereas zero or one MeOH molecules form the HB structures. By increasing the H<sub>2</sub>O molar fraction from pure MeOH, total coordination numbers of MeOH and H<sub>2</sub>O are close to two. The HB structure of H<sub>2</sub>O with PNIPAM is easily formed compared to that of MeOH. In pure H<sub>2</sub>O ( $x = 0.0$ ), two H<sub>2</sub>O molecules form HB structures with the C=O group of PNIPAM.

## S5. Influences of Solvation Effects in Inner-Shell Calculations

Table S2 shows the photon energy of the C=O  $\pi^*$  peaks in calculated O K-edge inner-shell spectra of NIPAM and PNIPAM with 5 polymer units including the solvent effects of MeOH or H<sub>2</sub>O using the polarizable continuum model. The photon energy of the C=O  $\pi^*$  peak of NIPAM including a solvent effect of MeOH was set to 0 meV. The C=O  $\pi^*$  peak of NIPAM including a solvent effect of H<sub>2</sub>O is -3 meV and is close to that including a solvent effect of MeOH. The C=O  $\pi^*$  peak of NIPAM with no solvent effect shows a higher energy shift of 62 meV. Since the solvent effect of H<sub>2</sub>O in NIPAM is close to that of MeOH, the inner-shell calculations of NIPAM were performed by using the molecular structure of NIPAM including the solvent effect of MeOH.

The C=O  $\pi^*$  peak of PNIPAM including a solvent effect of H<sub>2</sub>O shows the energy shift of -5 meV compared to that including a solvent effect of MeOH. On the other hand, the C=O  $\pi^*$  peak of PNIPAM with no solvent effect shows a higher energy shift of 102 meV. Since the solvent effect of H<sub>2</sub>O in PNIPAM is close to that of MeOH, the inner-shell calculations of PNIPAM were performed by using the molecular structures of PNIPAM with 5 polymer units including the solvent effect of MeOH.

**Table S2.** The photon energy of the C=O  $\pi^*$  peaks in calculated O K-edge inner-shell spectra of NIPAM and PNIPAM including the solvent effects of MeOH or H<sub>2</sub>O. The photon energy is relative to the C=O  $\pi^*$  peak of NIPAM including a solvent effect of MeOH. The energy shifts ( $\Delta E$ ) of the C=O  $\pi^*$  peaks in NIPAM and PNIPAM from same molecules including solvent effects of MeOH are also shown.

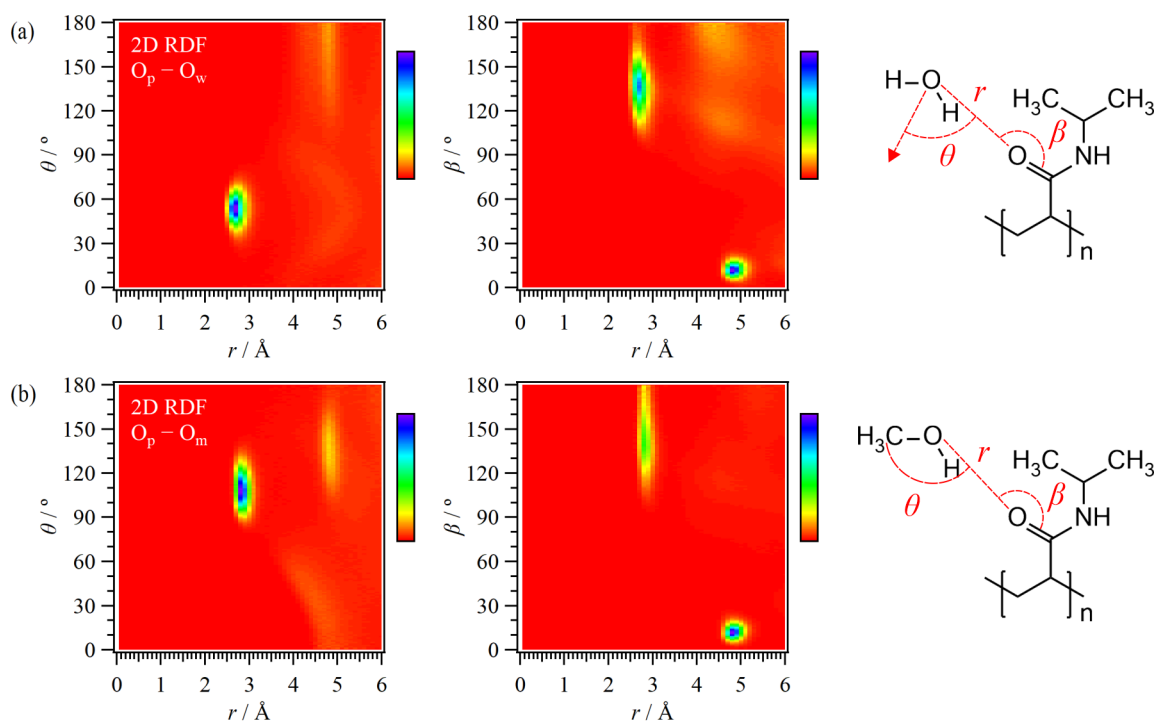
Molecule	Solvent	Energy / meV	$\Delta E$ / meV
NIPAM	–	62	62
	MeOH	0	0
	H <sub>2</sub> O	-3	-3
PNIPAM	–	318	102
	MeOH	216	0
	H <sub>2</sub> O	211	-5



## S6. Two-Dimensional RDF around C=O Group in PNIPAM

For building the HB model structures of PNIPAM with MeOH and H<sub>2</sub>O, we have calculated 2D RDF of O<sub>p</sub> with O<sub>m</sub> and O<sub>w</sub>, as shown in Figure S5. Figure S5(a) shows 2D RDF O<sub>p</sub> – O<sub>w</sub> in pure H<sub>2</sub>O. The horizontal axes of both panels are the distance  $r$  between O<sub>p</sub> and O<sub>w</sub>. The vertical axis of the left panel is the angle  $\theta$  between the normal vector of H<sub>2</sub>O and  $r$ . The vertical axis of the right panel is the angle  $\beta$  between the C=O molecular axis of PNIPAM and  $r$ . From the left panel, we found that  $r = 2.7 \text{ \AA}$  and  $\theta = 55^\circ$ . We also found that  $\beta = 136^\circ$  at  $r = 2.7 \text{ \AA}$  from the right panel. These results indicate that the HB structures of the C=O group in PNIPAM with H atom of H<sub>2</sub>O is stable in pure H<sub>2</sub>O. On the other hand, the peak around  $5 \text{ \AA}$  is assigned as the HB structure between the NH group in PNIPAM with oxygen atom of H<sub>2</sub>O.

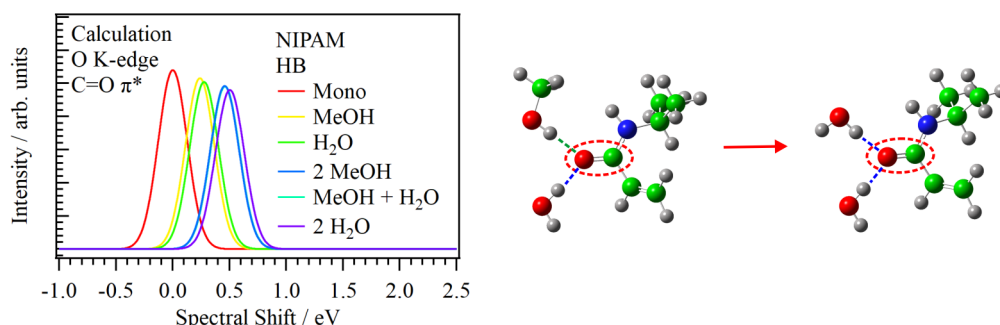
Figure S5(b) shows 2D RDF O<sub>p</sub> – O<sub>m</sub> in pure MeOH. The horizontal axes of both panels are the distance  $r$  between O<sub>p</sub> and O<sub>m</sub>. The vertical axis of the left panel is the angle  $\theta$  between the O-CH<sub>3</sub> molecular axis of MeOH and  $r$ . The vertical axis of the right panel is the angle  $\beta$  between the C=O molecular axis of PNIPAM and  $r$ . From the left panel, we found that  $r = 2.8 \text{ \AA}$  and  $\theta = 110^\circ$ . We also found that  $\beta = 139^\circ$  at  $r = 2.8 \text{ \AA}$  from the right panel. These results indicate that the HB structures of the C=O group in PNIPAM with H atom of MeOH is stable in pure MeOH. On the other hand, the peak around  $5 \text{ \AA}$  is assigned as the HB structure between the NH group in PNIPAM with oxygen atom of MeOH.



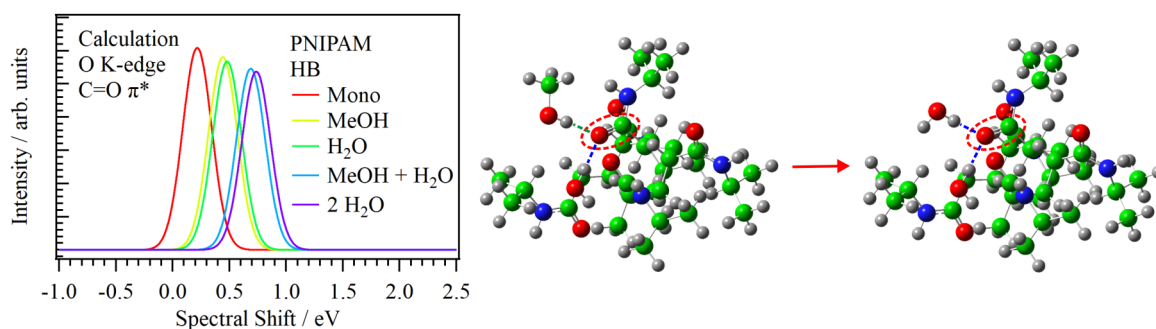
**Figure S5.** (a) 2D RDF O<sub>p</sub> – O<sub>w</sub> in pure H<sub>2</sub>O. The horizontal axes are the distance  $r$  between O<sub>p</sub> and O<sub>w</sub>. The vertical axes of the left and right panels are the angles  $\theta$  and  $\beta$ , respectively, whose definitions are described in the inset. (b) 2D RDF O<sub>p</sub> – O<sub>m</sub> in pure MeOH. The horizontal axes are the distance  $r$  between O<sub>p</sub> and O<sub>m</sub>. The vertical axes of the left and right panels are the angles  $\theta$  and  $\beta$ , respectively, whose definitions are described in the inset.

## S7. Inner-Shell Spectra of Different HB Structures of NIPAM and PNIPAM

Figure S6 shows inner-shell spectra of NIPAM with the HB structures of MeOH and H<sub>2</sub>O. As discussed in Sec. S5, the molecular structure of NIPAM is optimized with including the solvent effect of MeOH. The HB structures of NIPAM with MeOH and H<sub>2</sub>O were obtained from the distances and angles of the HB structures determined by 2D RDF described in Sec. S6. The C=O  $\pi^*$  peak in the HB structure of NIPAM with one MeOH molecule shows a higher energy shift of 239 meV compared to that of isolated NIPAM molecule. The C=O  $\pi^*$  peak in the HB structure of NIPAM with one H<sub>2</sub>O molecule shows a higher energy shift of 276 meV. The C=O  $\pi^*$  peaks in the HB structures of NIPAM with two MeOH, MeOH + H<sub>2</sub>O, and two H<sub>2</sub>O molecules are 457 meV, 460 meV, and 503 meV, respectively. Note that the energy shift of the C=O  $\pi^*$  peak in the HB structures of NIPAM from MeOH + H<sub>2</sub>O to two H<sub>2</sub>O is 43 meV. The C=O  $\pi^*$  peak shows a higher energy shift by increasing the coordination numbers of the HB structures of the C=O group in NIPAM with MeOH and H<sub>2</sub>O. The C=O  $\pi^*$  peak in the HB structure of NIPAM with H<sub>2</sub>O shows a higher energy shift than that with MeOH.



**Figure S6.** Calculated O K-edge inner-shell spectra of several HB structures of the C=O group in NIPAM with MeOH and H<sub>2</sub>O molecules. The photon energy is relative to the C=O  $\pi^*$  peak of an isolated NIPAM molecule. The inset shows HB structures of NIPAM with MeOH and H<sub>2</sub>O molecules.



**Figure S7.** Calculated O K-edge inner-shell spectra of several HB structures of the C=O group in PNIPAM with H<sub>2</sub>O and MeOH molecules. The photon energy is relative to the C=O  $\pi^*$  peak of an isolated NIPAM molecule. The inset shows HB structures of PNIPAM with MeOH and H<sub>2</sub>O molecules.

Figure S7 shows inner-shell spectra of PNIPAM with the HB structures of MeOH and H<sub>2</sub>O. As discussed in Sec. S5, the molecular structure of PNIPAM with 5 polymer units was optimized with including the solvent effect of MeOH. The HB structures of PNIPAM with MeOH and H<sub>2</sub>O were obtained from the distances and angles of the HB structures determined by 2D RDF described in Sec. S6. Note that the oxygen atom in the central polymer unit was only used in the inner-shell calculations of PNIPAM. The C=O  $\pi^*$  peak of isolated PNIPAM polymer is 216 meV. The C=O  $\pi^*$  peak in the HB structure of PNIPAM with one MeOH molecule is 442 meV. The C=O  $\pi^*$  peak in the HB structure of PNIPAM with one H<sub>2</sub>O molecule is 479 meV. The C=O  $\pi^*$  peaks in the HB structures of PNIPAM with MeOH + H<sub>2</sub>O and two H<sub>2</sub>O molecules are 688 meV and 736 meV, respectively. Note that the energy shift of C=O  $\pi^*$  peak in the HB structures of PNIPAM from MeOH + H<sub>2</sub>O to two H<sub>2</sub>O is 48 meV. Since the polymer units of PNIPAM were close to each other by the structural optimization, it is difficult to form the HB structures with two MeOH molecules. The PNIPAM chains spread in pure MeOH, and the C=O group in the polymer unit of PNIPAM can form the HB structures with two MeOH molecules in the actual environment. The C=O  $\pi^*$  peak shows a higher energy shift by increasing the coordination numbers of the HB structures of the C=O group in PNIPAM with MeOH and H<sub>2</sub>O. The C=O  $\pi^*$  peak in the HB structure of PNIPAM with H<sub>2</sub>O shows a higher energy shift than that with MeOH.

## S8. Discussion of Structures of Aqueous MeOH Solutions

To elucidate the cononsolvency mechanism of PNIPAM, it is important to study the molecular interactions of MeOH and H<sub>2</sub>O in aqueous MeOH solutions (MeOH)<sub>x</sub>(H<sub>2</sub>O)<sub>1-x</sub> with different molar fractions. Liquid MeOH ( $x = 1.0$ ) has one- or two-dimensional HB chain structures, because MeOH has only one donor site.<sup>23-28</sup> On the contrary, liquid H<sub>2</sub>O ( $x = 0.0$ ) shows tetrahedrally coordinated three-dimensional HB network.<sup>29</sup> Mass<sup>30</sup> and Raman spectroscopies<sup>31</sup> revealed that the local structures of aqueous MeOH solutions show three concentration regions with borders at  $x = 0.7$  and  $x = 0.3$ . This is because aqueous solutions are generally non-uniform and consist of microscopic cluster structures.<sup>32</sup> Using C K-edge XAS analysis, we also revealed that the hydrophobic interaction of the methyl group in MeOH exhibits three concentration regions with borders at  $x = 0.7$  and  $x = 0.3$ .<sup>19</sup> In the MeOH-rich region ( $x > 0.7$ ), the MeOH molecules formed chain structures with nearly zero influence from H<sub>2</sub>O. The hydrophobic clusters of MeOH are formed at  $0.7 > x > 0.3$ , which was also observed in the neutron scattering experiment.<sup>33</sup> In the H<sub>2</sub>O-rich region ( $0.3 > x$ ), the hydrophobic clusters are embedded in the three-dimensional HB network of H<sub>2</sub>O owing to the decrease in the cluster size. The formed structures in aqueous MeOH solutions were also investigated using MD simulations.<sup>34-36</sup>

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