

### Electronic supplementary information for

#### **First-principles study of the surface energies and electronic structures of $\gamma$ -CsSnI<sub>3</sub> surfaces**

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#### **Calculating the surface energies of polar surfaces**

Sn-rich condition is good for suppressing the formation of Sn-vacancy in CsSnI<sub>3</sub>, which is considered as the most detrimental defect, promoting high luminescent efficiency.<sup>1</sup> For Sn-perovskites, Sn-rich condition may be achieved by the presence of excess Sn sources such as Sn metal and SnF<sub>2</sub> additive.<sup>2,3</sup> Nakamura et al. found that Sn metal nanoparticles are produced from the reaction of a dihydropyridine derivative with SnF<sub>2</sub> (II) additive.<sup>3</sup> Assuming Sn-rich condition (i.e.,  $\Delta\mu_{\text{Sn}}=0$  eV), we obtain  $\Delta\mu_{\text{Cs}} + 3\Delta\mu_{\text{I}} = -4.720$  eV for CsSnI<sub>3</sub>. Combining Sn-rich with CsI-rich ( $\Delta\mu_{\text{Cs}} + \Delta\mu_{\text{I}} = -3.081$  eV) conditions of stable CsSnI<sub>3</sub>, we obtain:  $\Delta\mu_{\text{Cs}} = -2.261$  eV and  $\Delta\mu_{\text{I}} = -0.820$  eV. Similarly, we obtain:  $\Delta\mu_{\text{Cs}} = -2.536$  eV and  $\Delta\mu_{\text{I}} = -0.728$  eV at SnI<sub>2</sub>-rich condition ( $\Delta\mu_{\text{Cs}} + \Delta\mu_{\text{I}} = -3.264$  eV). The surface energies ( $E_{\text{surf}}$ ) of (100), (010) and (101) polar surfaces are calculated at these conditions. In particular, the  $E_{\text{surf}}$  of (101) surface, either for CsI<sub>3</sub>- or Sn-termination, only depends on  $\Delta\mu_{\text{Cs}} + 3\Delta\mu_{\text{I}}$ :

$$E_{\text{surf}}^{\text{CsI}_3} = \lambda_{\text{CsI}_3} - \frac{1}{S}(\Delta\mu_{\text{Cs}} + 3\Delta\mu_{\text{I}}) \quad (1)$$

$$E_{\text{surf}}^{\text{Sn}} = \lambda_{\text{Sn}} + \frac{1}{S}(\Delta\mu_{\text{Cs}} + 3\Delta\mu_{\text{I}}) \quad (2)$$

For Sn-rich condition, the  $E_{\text{surf}}$  of (101) surfaces can be calculated according to  $\Delta\mu_{\text{Cs}} + 3\Delta\mu_{\text{I}} = -4.720$  eV and thus is independent of CsI-rich or SnI<sub>2</sub>-rich condition.

#### **Effect of dispersion force on the structures and surface energies**

To show the effect of London dispersion force on the structures and surface energies, we use PBEsol+D3<sup>4</sup> to calculate the structural parameters of CsSnI<sub>3</sub> bulk and the surface energies of (001) surfaces. It is seen from Table S2 that the structural parameters calculated by PBEsol and PBEsol+D3 are consistent with each other and agree well with experimental values. From Table S3, one can find that the surface energies of the CsI-(001) and SnI<sub>2</sub>-(001) surfaces calculated from PBEsol+D3 are overall increased by 0.07 and 0.11 J/m<sup>2</sup>, respectively, in comparison with the values obtained from PBEsol. It is worth noting that the difference in the surface energies between SnI<sub>2</sub>-(001) and CsI-(001) surfaces increase about 0.04 J/m<sup>2</sup>, indicating that CsI-termination remains the most stable compared with SnI<sub>2</sub>-termination with the consideration of London dispersion force. Therefore, including London dispersion force in the calculations does not change the main conclusions based on PBEsol.

**Table S1** Surface energies ( $E_{\text{surf}}$ , J/m<sup>2</sup>) of the CsSnI<sub>3</sub> (001) surfaces calculated by using the cutoff values of 350 eV and 500 eV.

$E_{\text{surf}}$	350 eV		500 eV	
	CsI-rich	SnI <sub>2</sub> -rich	CsI-rich	SnI <sub>2</sub> -rich
CsI-(001)	0.073	0.113	0.071	0.110
SnI <sub>2</sub> -(001)	0.196	0.156	0.194	0.154

**Table S2** Structural parameters of  $\gamma$ -CsSnI<sub>3</sub> calculated by using PBEsol and PBEsol+D3. Experimental (Expt.) values<sup>5</sup> are also listed here for a comparison.

	a (Å)	b	c
PBEsol	8.779	8.344	12.238
PBEsol+D3	8.525	8.510	12.225
Expt.	8.688	8.643	12.378

**Table S3** Surface energies ( $E_{\text{surf}}$ , J/m<sup>2</sup>) of  $\gamma$ -CsSnI<sub>3</sub> (001) surfaces calculated by using PBEsol and PBEsol+D3.

$E_{\text{surf}}$	PBEsol		PBEsol+D3	
	CsI-rich	SnI <sub>2</sub> -rich	CsI-rich	SnI <sub>2</sub> -rich
CsI-(001)	0.073	0.113	0.141	0.185
SnI <sub>2</sub> -(001)	0.196	0.156	0.309	0.265

**Table S4** Surface energies ( $E_{\text{surf}}$ , J/m<sup>2</sup>) of the (100), (010) and (101) polar surfaces. Assuming Sn-rich condition (i.e.,  $\Delta\mu_{\text{Sn}}=0$  eV), the  $E_{\text{surf}}$  are calculated at CsI-rich:  $\Delta\mu_{\text{Cs}} = -2.261$  eV and  $\Delta\mu_{\text{I}} = -0.820$  eV, and SnI<sub>2</sub>-rich condition:  $\Delta\mu_{\text{Cs}} = -2.536$  eV and  $\Delta\mu_{\text{I}} = -0.728$  eV.

$E_{\text{surf}}$	(100)		(010)		(101)	
	CsSnI	I <sub>2</sub>	CsSnI	I <sub>2</sub>	CsI <sub>3</sub>	Sn
CsI-rich	0.467	0.312	0.449	0.312	0.243	0.418
SnI <sub>2</sub> -rich	0.496	0.283	0.477	0.285	0.243	0.418

**Table S5**  $E_{\text{surf}}$  (J/m<sup>2</sup>) of polarity-compensated (010) surface. Three terminations (Cs, SnI and I) are considered for (010) surface. The  $E_{\text{surf}}$  values are calculated at the CsI-rich ( $\Delta\mu_{\text{Cs}}+\Delta\mu_{\text{I}}=-3.081$  eV) and SnI<sub>2</sub>-rich ( $\Delta\mu_{\text{Cs}}+\Delta\mu_{\text{I}}=-3.264$  eV) conditions.

$E_{\text{surf}}$	(010)		
	Cs	SnI	I
CsI-rich	0.120	0.231	0.118
SnI <sub>2</sub> -rich	0.175	0.177	0.118

**Table S6**  $E_{\text{surf}}$  ( $\text{J}/\text{m}^2$ ) of the (100) and (001) surfaces of  $\text{CsSnI}_3$  and  $\text{CsPbI}_3$ . The  $E_{\text{surf}}$  of the (100) and (001) surfaces of  $\text{CsPbI}_3$  are taken from a previous study.<sup>6</sup>

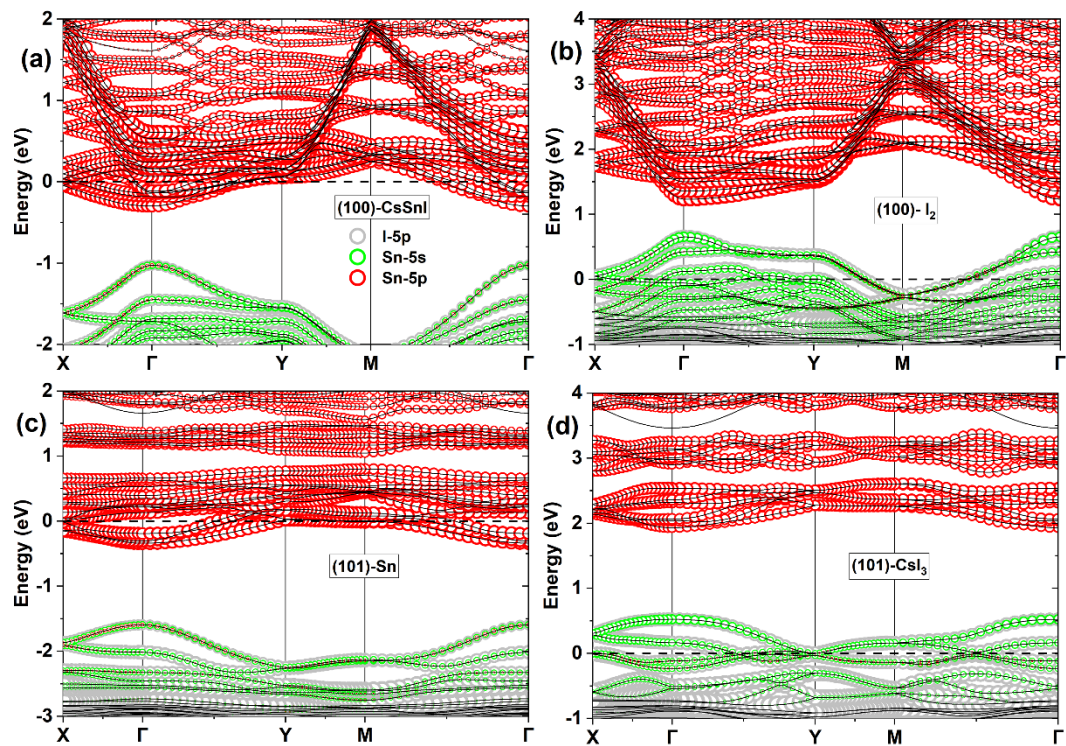
		(100)					(001)	
M	$E_{\text{surf}}$	Cs	MI	CsMI	I <sub>2</sub>	I	CsI	MI <sub>2</sub>
Sn	CsI-rich	0.119	0.246	0.467	0.312	0.118	0.073	0.196
	SnI <sub>2</sub> -rich	0.176	0.189	0.496	0.283	0.118	0.113	0.156
Pb	CsI-rich	0.112	0.252	0.465	0.532	0.130	0.069	0.193
	PbI <sub>2</sub> -rich	0.137	0.228	0.477	0.520	0.130	0.085	0.177

**Table S7** Band gaps ( $E_g$ ) of the (110), (001), (100) and (101) surfaces of  $\text{CsSnI}_3$ . The calculated band gap of the bulk is 0.53 eV. Nine-layers slabs are used for the (110) and (001) surfaces; thirteen-layers slabs for the (100)-Cs, (100)-SnI and (100)-I surfaces; fifteen-layers for the (101)-CsI<sub>2</sub> and (101)-I\*. These choices ensure a similar thickness of about 29 Å for all these slabs. Orbital-projected band structures and the effective masses in the main text and ESI are calculated using these slabs.

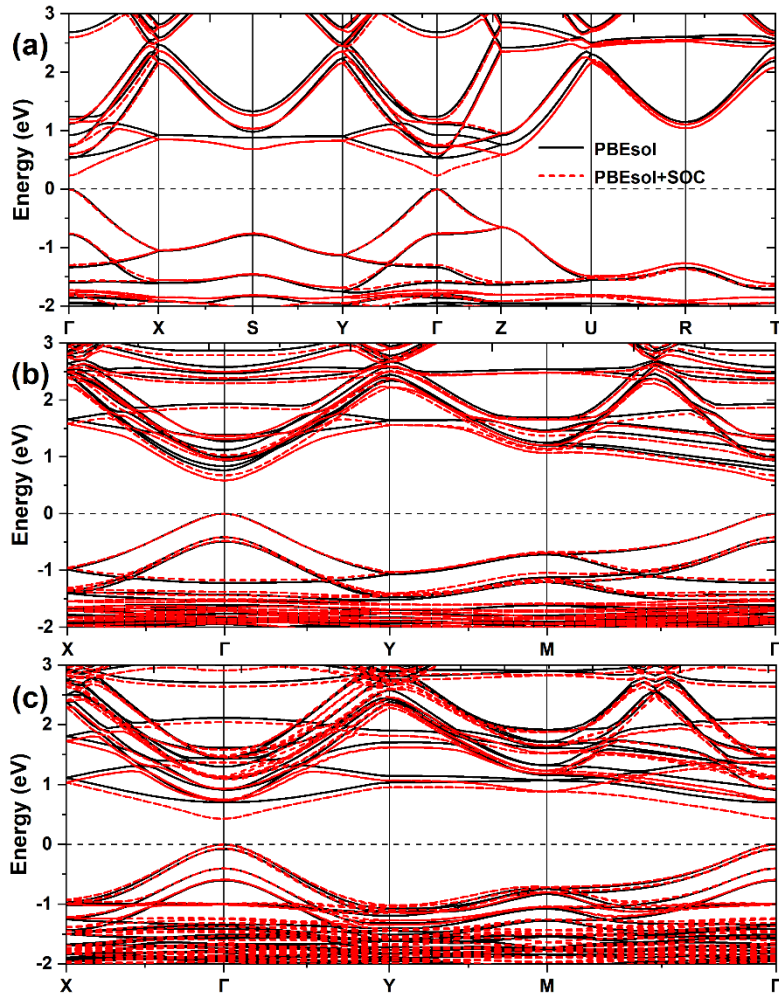
Surfaces	(110)		(001)		(100)			(101)	
Terminations	CsI	SnI <sub>2</sub>	CsI	SnI <sub>2</sub>	Cs	SnI	I	CsI <sub>2</sub>	I*
$E_g$	0.68	0.65	0.72	0.64	0.87	0.80	0.83	0.88	0.99

**Table S8** Calculated effective mass (in units of the free-electron mass  $m_0$ ) of hole ( $m_h$ ) and electron ( $m_e$ ) along  $\Gamma$ -M for the (110), (001), (100) and (101) surfaces.

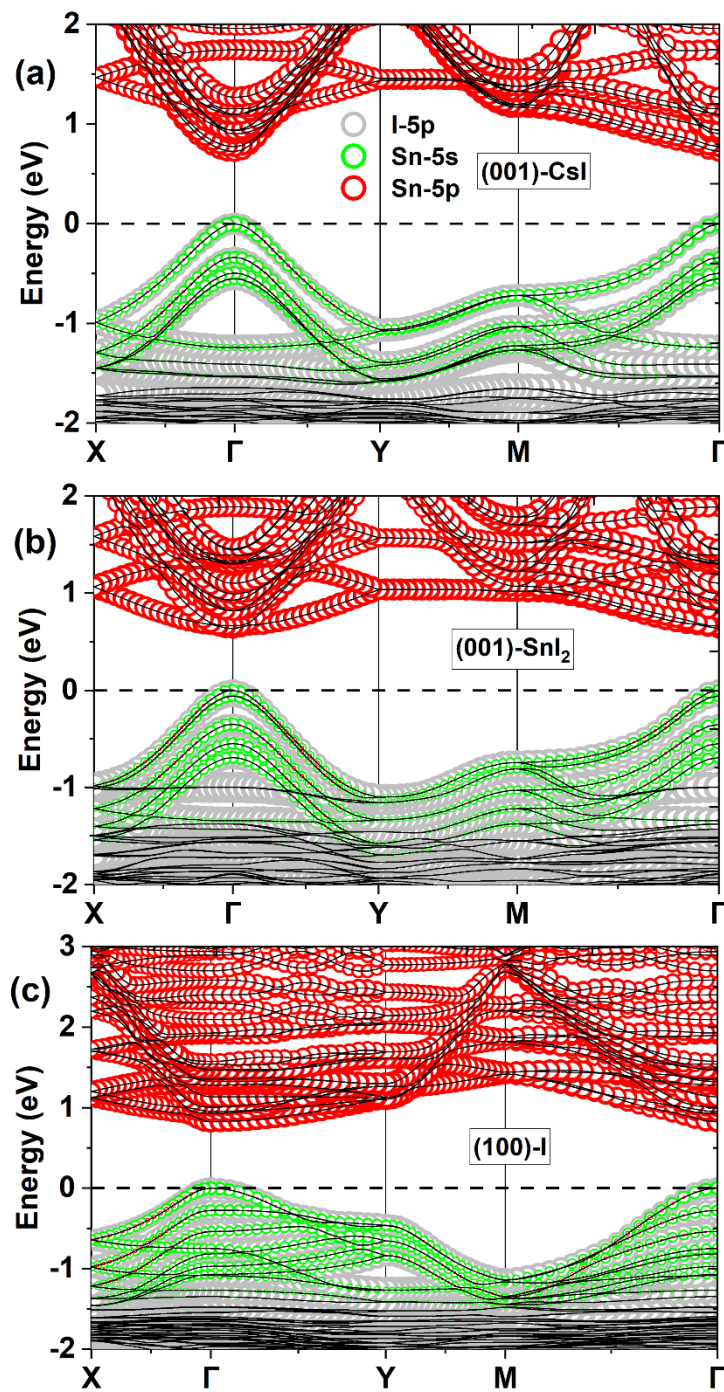
Surfaces	(110)		(001)		(100)			(101)	
Terminations	CsI	SnI <sub>2</sub>	CsI	SnI <sub>2</sub>	Cs	SnI	I	CsI <sub>2</sub>	I*
$m_h$	0.26	0.26	0.31	0.34	0.54	0.39	0.45	0.65	0.73
$m_e$	0.30	0.74	0.36	0.41	0.55	0.97	0.52	1.06	2.41



**Figure S1** Orbital-projected band structures of the CsSnI- and I<sub>2</sub>-terminated (100) surfaces, and the Sn- and CsI<sub>3</sub>-terminated (101) surfaces. The Fermi level is indicated by the horizontal dashed line.



**Figure S2** Band structures calculated without (PBEsol) and with SOC (PBEsol+SOC) for: (a) the bulk; (b) CsI-(001) surface; (c) SnI<sub>2</sub>-(001) surface. The energy of the VBM is set as zero. Seven-layers slabs are used for the (001) surfaces.



**Figure S3** Orbital-projected band structures of the (001)-CsI, (001)-SnI<sub>2</sub> and (100)-I surfaces. The energy of the VBM is set as zero.

### References

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