Supporting information for:

Single-Molecule Tip-Enhanced Raman Spectroscopy of C₆₀ on the Si(111)-(7×7) Surface

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1. Methods for sample preparation and SPM-TERS measurements

Sample preparation. All experiments were performed in ultra-high vacuum (UHV) chambers (base pressure $< 5 \times 10^{-10}$ mbar). We used an n-type Si(111) wafer (resistance: 0.03–0.08 Ω , thickness: 0.5 mm). The sample was degassed at ~800 K in the chamber overnight, then it was annealed several times by direct current heating to ~1470 K for 10 seconds. The C₆₀ molecule was purchased from Aldrich and used without further purification. The molecules were evaporated at 610 K from a K-cell evaporator onto the substrate held at room temperature.

STM measurement. We used a low-temperature STM from UNISOKU Co., Ltd. (modified USM-1400) operated with Nanonis SPM Controller (SPECS GmbH). The bias voltage (V_{bias}) was applied to the sample, and the tip was grounded. The tunnelling current (j_{STM}) was collected from the tip. We used Ag tips fabricated by focused ion beam milling^{S1}. Temperature of the system (80 or 10 K) is indicated in each figure caption.

TERS measurement. The excitation laser was focused to the STM junction with an insitu Ag-coated parabolic mirror (numerical aperture of ~0.6) mounted on the cold STM stage. The parabolic mirror was precisely aligned using piezo motors (Attocube GmbH), which allow three translational and two rotational motions. For the Raman measurements we used a HeNe laser for 633 nm (Laser 2000 GmbH) and a solid-state laser for 532 nm (Cobolt). The incident beam is linearly polarized along the tip axis (ppolarization). The scattered photons are collected by the same parabolic mirror and detected outside of the UHV chamber with a grating spectrometer (AndorShamrock 303i). Regarding the data processing, the Raman spectra shown in Figure 2 of the main text and **Figure S2** are baseline-corrected by using the asymmetrically reweighted penalized least squares (arPLS) method^{S2}.

2. Raman maps



Figure S1. (a) STM image under illumination for single C₆₀ on the faulted half unit cell of the Si(111)-7x7 surface (80 K, $V_{\text{bias}} = 0.3 \text{ V}$, $j_{\text{STM}} = 60 \text{ pA}$). **(b)** TERS spectrum of the C₆₀ in (a) where the three high frequency vibrational modes at 1382, 1441 and 1478 cm⁻¹ are observed (80 K, $V_{\text{bias}} = 0.3 \text{ V}$, $j_{\text{STM}} = 50 \text{ pA}$, $\lambda_{\text{ext}} = 532 \text{ nm}$, $P_{\text{ext}} = 5.5 \text{ mW}$). **(c-e)** TERS maps of the different peaks marked in (b). Each of the scale bars is independent and in arbitrary units for clarity.

3. Estimation of the enhancement factors

To estimate the enhancement factor of a single C_{60} TERS on the Si(111)-(7x7) surface, we compare the TERS intensity of the $A_g(2)$ mode to the Raman peak of the optical phonon mode of the bulk Si. Since the phonon peak remains constant regardless of the tip position^{S3}, it can be used as a reference of the Raman scattering intensity. The total effective enhancement factor g_{total} is given by:

$$g_{\text{total}} = \frac{\Gamma_{\text{NF}}}{\Gamma_{\text{FF}}} \cdot \frac{N_{\text{Si}}\sigma_{\text{Si}}}{N_{\text{C}_{60}}\sigma_{\text{C}_{60}}},$$
(1)

where $\Gamma_{\rm FF~(NF)}$ is the far-(near-) field Raman scattering rate, $N_{\rm Si}$ is the number of the Si atoms within the illuminated volume, and $N_{\rm C_{60}} = 1$ is the number of C₆₀, and $\sigma_{\rm Si~(C_{60})}$ is the Raman cross section of Si (C₆₀). As evident in Figures 1 and 2 of the main text, the Raman intensities of far-field bulk Si phonon peak and near-field C₆₀ $A_{\rm g}$ (2) peak are comparable; $\Gamma_{\rm FF} \approx \Gamma_{\rm NF}$.

Therefore, Equation (1) is rewritten as follows:

$$g_{\text{total}} \approx \frac{\sigma_{\text{Si}}}{\sigma_{\text{C}_{60}}} \cdot \rho_{\text{Si}} \cdot 2\pi \left(\frac{d_{\text{laser}}}{2}\right)^2 d_{\text{p}},$$
 (2)

where $\rho_{\rm Si} = 4.992 \times 10^{22}$ atoms cm⁻³ is the Si atom density^{S4}, $d_{\rm laser} = 3 \,\mu {\rm m}$ is the diameter of the Gaussian-beam spot^{S1}, $d_{\rm p} = 1.3 \times 10^{-4}$ cm is the penetration depth in Si at 530 nm (Ref. **S5**).

References **S4** and **S6** show the cross section of Si at 785 nm is $(1.0 \pm 0.2) \times 10^{-27} \text{ cm}^2$ and the differential cross section of C₆₀ at 752 nm is $(2.1 \pm 0.3) \times 10^{-29} \text{ cm}^2/\text{sr}$, respectively. According to the Rayleigh scattering formula, the cross section is proportional to the minus fourth power of the wavelength. We derive $\sigma_{\text{Si}} = 4.7 \times 10^{-27} \text{ cm}^2$ and $\sigma_{\text{C}_{60}} = 1.0 \times 10^{-27} \text{ cm}^2$ at 532 nm. From Equation (2), thus, we get:

$$g_{\rm total} \approx 10^{12}$$

The total enhancement factor g_{total} is the product of the chemical enhancement factor g_{Chem} and the electromagnetic enhancement factor g_{EM} :

$$g_{\text{total}} = g_{\text{chem}} \cdot g_{\text{EM}}.$$
 (3)

Taking $g_{EM} \propto \left(\frac{|E_z|}{|E_0|}\right)^4 \approx 10^9$ from Reference **S7**, where the EM field in an Ag-tip–Si-surface junction in the tunnelling regime was simulated, we obtain $g_{\text{chem}} \approx 10^3$.

4. TERS spectra in the tunnelling regime with the Si signal subtracted



Figure S2. Si-bulk-phonon-peak-subtracted version of Figures 2a–c in the main text for better visualization of low frequency region. The intensity of each spectrum was normalized based on the bulk phonon peak intensity (Figure 2d) before subtraction.

5. Comparison of the observed $C_{\rm 60}$ vibrational frequencies

Table S1. TERS observed peak frequencies (in cm^{-1}) for single C₆₀ molecules on Si(111) of the present experiment (blue columns) and previous experiments (grey columns)

TERS, sir	TERS, single C ₆₀ on Si(111) ^{a}	
Mol #1	Mol #2	Mol #3
(Fig. 2a)	(Fig. 2b)	(Fig. 2c)
247 w	256 w	
418 w		410 vw
448 w	443 vw	
492 sh		487 sh
577 m		569 w
704 m	707 m	
728 m	734 sh	742 w
797 m	760 sh	766 w
821 w		835 vw
921 w	927 w	934 vw
1054 w	1054 w	1041 m
1028 w		
	1082 w	
1156 w	1131 w	1164 m
	-	1187 w
1225 w		
	1402 s	
1416 m	1419 sh	
	1437 211	1377 s
1446 s	1446 s	1441 s
1-1-TO 5	1401	1402 m
	1481 W	1483 M
	1503 m	
1564 s	1564 s 1583 sh	1559 w

^{*a*} vw: very weak, w: weak, m: medium, s: strong, sh: shoulder.

^b Ref. S8.

^c Ref. S9.

^{*d*} R: Raman active, IR: infrared active.

^e Ref. S10.



6. Detection of adsorption-site switching of a C₆₀ molecule

Figure S3. (a) Waterfall plot of the time evolution of the TERS signal in the tunnelling regime of a single C₆₀ molecule adsorbed on the fault half unit cell (red dot, see inset below), stable over time at large tunnelling currents (5 s per spectrum, 10 K, $V_{\text{bias}} = 0.3 \text{ V}$, $j_{\text{STM}} = 5 \text{ nA}$, $\lambda_{\text{ext}} = 532 \text{ nm}$, $P_{\text{ext}} = 5.5 \text{ mW}$). **(b)** The two TERS spectra (blue and green arrows in a) showing stability over time. **(c)** Same as in a for another C₆₀ molecule, but in this case the time evolution clearly shows the transition between two bistable states (5 s per spectrum, 10 K, $V_{\text{bias}} = 0.3 \text{ V}$, $j_{\text{STM}} = 100 \text{ pA}$, $\lambda_{\text{ext}} = 532 \text{ nm}$, $P_{\text{ext}} = 5.5 \text{ mW}$). **(d)** STM image of two single C₆₀ molecules adsorbed on the fault half unit cell, with the right one showing brighter strips corresponding to subtle displacements originated with the tip scanning motion. The two TERS spectra (blue and green arrows in c) correspond to each of the bistable states observed in the waterfall plot.

7. Observation of the overtones and combination bands in the MPC regime in the faulted and unfaulted half unit cells



Figure S4. (a) Waterfall plot of TERS as a function of the tip height recorded during tip-approach and retraction over single C_{60} on the faulted half unit with the overtones and combination bands at higher frequencies (80 K, $V_{\text{bias}} = 0.3 \text{ V}$, $j_{\text{STM}} = 60 \text{ pA}$, $\lambda_{\text{ext}} = 532 \text{ nm}$, $P_{\text{ext}} = 5.5 \text{ mW}$, 5s/spectrum, Scale bar: arbitrary units). **(b)** TERS spectra in the tunnelling and contact regimes acquired at the positions of the corresponding arrow in a, highlighting that the overtones are only observed in the contact regime.



Figure S5. (a) Waterfall plot of TERS as a function of the tip height recorded during tip-approach and retraction over single C_{60} on the unfaulted half unit with the overtones and combination bands at higher frequencies (80 K, $V_{\text{bias}} = 0.3 \text{ V}$, $j_{\text{STM}} = 10 \text{ pA}$, $\lambda_{\text{ext}} = 532 \text{ nm}$, $P_{\text{ext}} = 5.5 \text{ mW}$, 5s/spectrum, Scale bar: arbitrary units). **(b)** TERS spectra in the tunnelling and contact regimes acquired at the positions of the corresponding arrow in a, highlighting that the overtones are only observed in the contact regime.

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