

Supplementary Material

Elemental ferroelectric topological insulator in ψ -bismuthene

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Note 1

To confirm the nontrivial topological order in ψ -bismuthene, the \mathbf{Z}_2 invariant is obtained by tracking the evolution of the Wannier charge center (WCC).^{1,2} The Wannier functions (WFs) with regard to the lattice vector \mathbf{R} is written as:³

$$|Rn\rangle = \frac{1}{2\pi} \int_{-\pi}^{\pi} dk e^{-ik(R-x)} |u_{nk}\rangle$$

The WFs depend on the gauge choice for the Bloch states $|u_{nk}\rangle$. Following Soluyanov and Vanderbilt's description, to optimally localize the WFs, the WCCs are defined as the mean value of the position operator defined by the mean value of $\bar{x}_n = \langle 0_n | \hat{X} | 0_n \rangle$, where the \hat{X} represent the position operator and $|0_n\rangle$ is the state corresponding to a WF in the cell with $\mathbf{R} = 0$.⁴ Then we obtain

$$\bar{x}_n = \frac{i}{2\pi} \int_{-\pi}^{\pi} dk \langle u_{nk} | \nabla_k | u_{nk} \rangle$$

Taking into account that $\sum_{\alpha} \bar{x}_{\alpha}^S = \frac{1}{2\pi} \int_{BZ} A^S$ for $S = \text{I or II}$, the summation in α is the occupied states and A is Berry connection, yielding the \mathbf{Z}_2 invariant following

$$Z_2 = \sum_{\alpha} [\bar{x}_{\alpha}^I\left(\frac{T}{2}\right) - \bar{x}_{\alpha}^{II}\left(\frac{T}{2}\right)] - \sum_{\alpha} [\bar{x}_{\alpha}^I(0) - \bar{x}_{\alpha}^{II}(0)]$$

here I and II are the Kramer partners. The identification of \mathbf{Z}_2 invariant can be obtained by counting the numbers of crossing between any arbitrary horizontal reference line and evolution of the WCCs, in which the odd number equals to $\mathbf{Z}_2 = 1$, revealing the nontrivial topological nature.

Since the bands near the Fermi level are primarily contributed by the p orbitals of Bi atoms, these orbitals are selected as projections to construct the maximally localized Wannier functions (MLWFs).⁴

⁶ A total of 48 bands, spanning discrete bands 17 to 64, along with 4 additional empty bands within the energy window of -6 to 6 eV relative to the Fermi level, are chosen for disentanglement, resulting in a set of well-localized WFs. Additionally, the spin Hall conductivity and Berry curvature calculations are performed using the Wannier interpolation method.^{7,8}

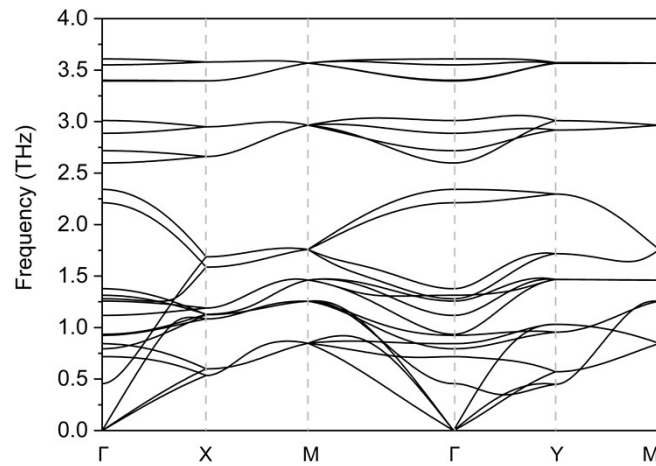


Fig. S1. Phonon spectrum for the 2D elemental ψ -bismuthene.

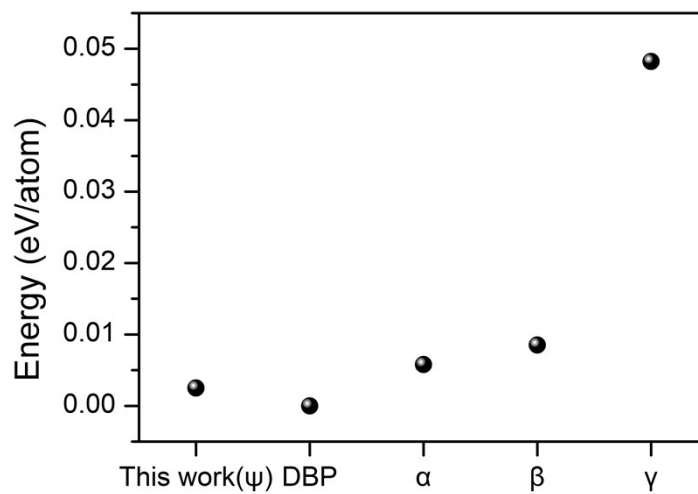


Fig. S2. Calculated binding energy difference between ψ , DBP, α , β , and γ bismuthene monolayers.

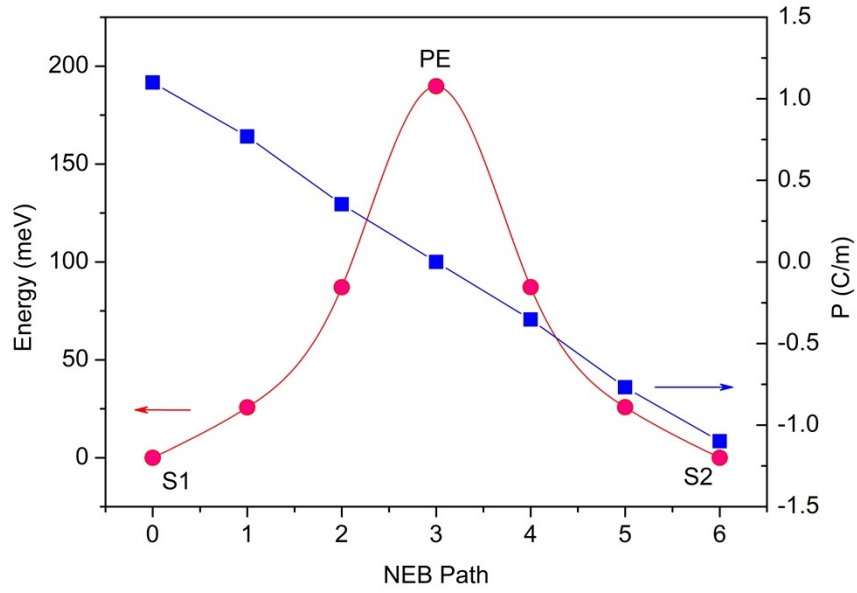


Fig. S3. The evolution of polarization P and energy landscape of the structural switching process between the energy-degenerate polarization states $S1$ and $S2$.

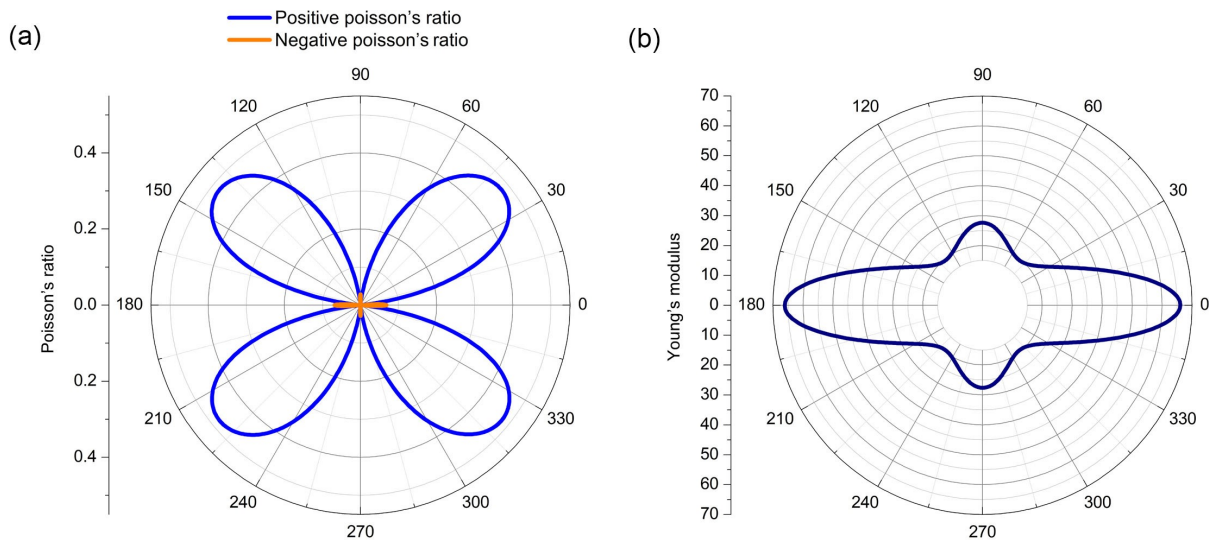


Fig. S4. (a) Poisson's ratio and (b) Young's modulus of ψ -bismuthene as a function of arbitrary angle.

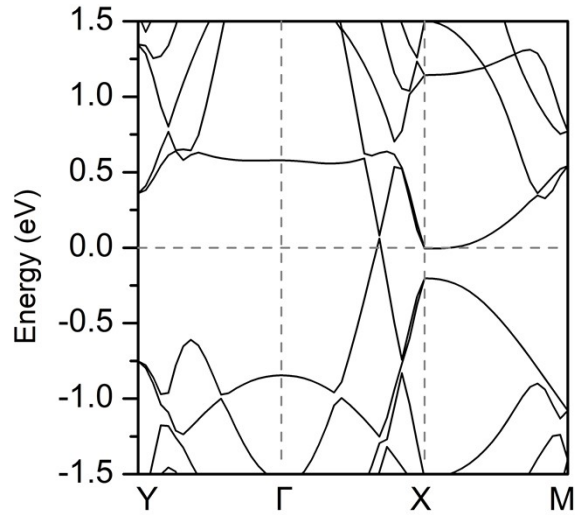


Fig. S5. Band structure of ψ -bismuthene calculated by Nanodcal using the mode-conservation pseudopotential and linear atomic orbit combination basis set.

References

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