

## Supporting Information

### Prediction of high- $T_c$ superconductivity in $H_6SX$ under submegabar pressure

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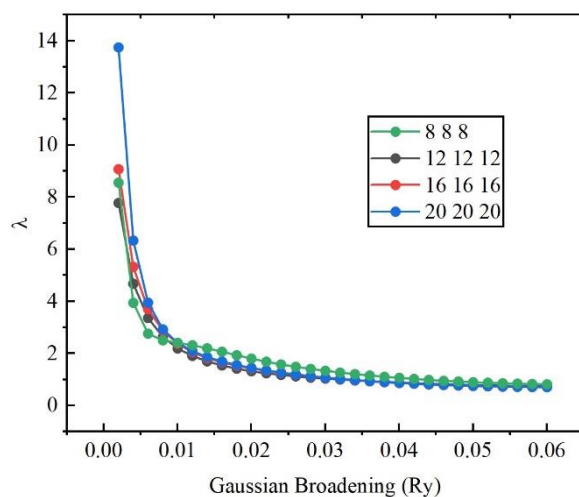


FIG. S1 Energy Electron-phonon coupling constant  $\lambda(q)$  as a function of the- Gaussian broadening  $\sigma$  for  $H_6SCl$  at  $\mathbf{q} = (0, 0, 0)$  at c200 GPa. We believe that  $12 \times 12 \times 12$   $k$ -meshes are converged in  $H_6SCl$ .

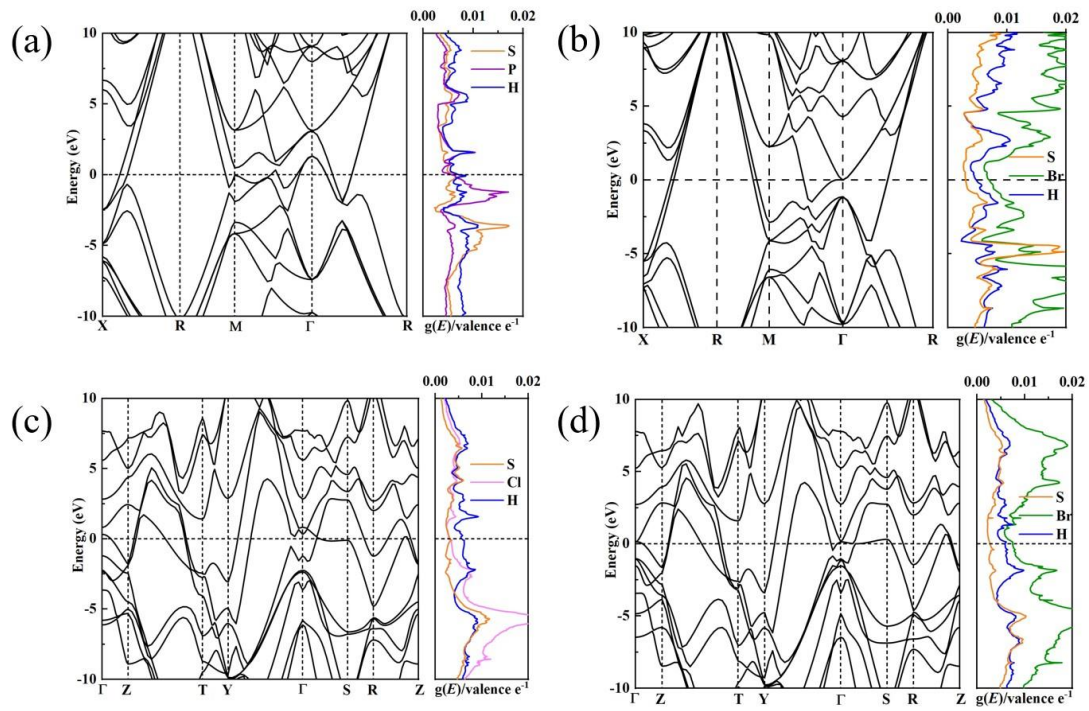


FIG. S2 Electronic band structure and DOS (states/eV/f.u.) of (a)  $Pm\bar{3}m$ -H<sub>6</sub>SP at 200 GPa, (b)  $Pm\bar{3}m$ -H<sub>6</sub>SBr at 200 GPa, (c)  $Cmmm$ -H<sub>6</sub>SCl at 200 GPa and (d)  $Cmmm$ -H<sub>6</sub>SBr at 200 GPa.

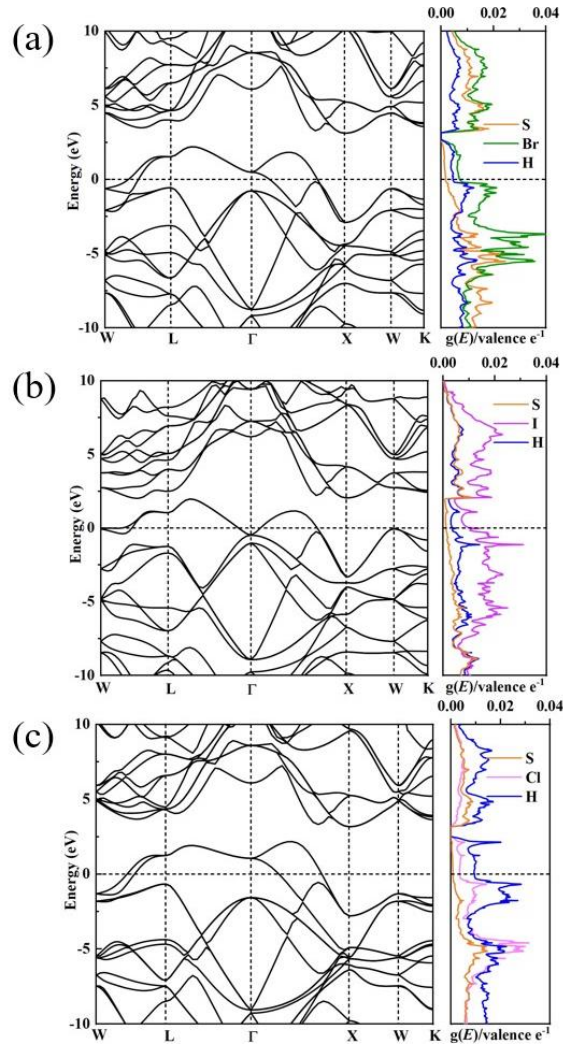


FIG. S3 Electronic band structure and DOS of (a)  $Fd\bar{3}m$ -H<sub>6</sub>SBr at 200 GPa, (b)  $Fd\bar{3}m$ -H<sub>6</sub>SI at 200 GPa, and (c)  $Fd\bar{3}m$ -H<sub>6</sub>SCl at 200 GPa.

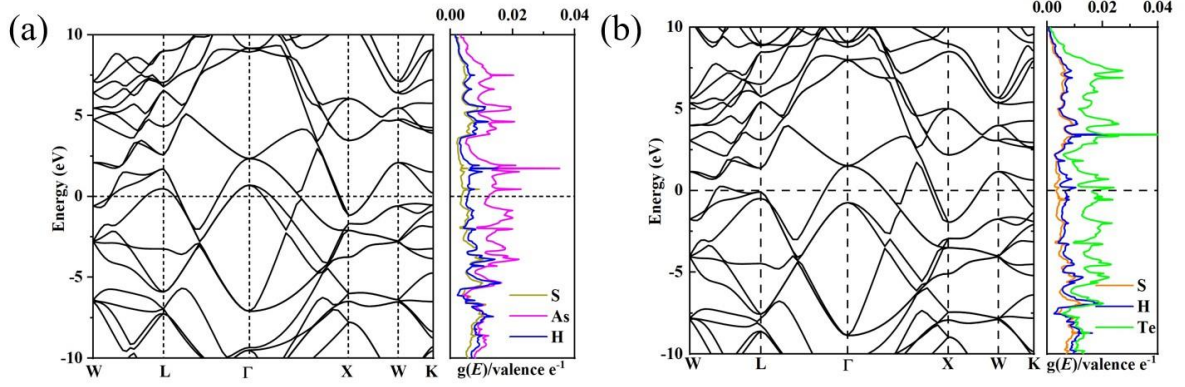


FIG. S4 Electronic band structure and DOS of (a)  $Fd\bar{3}m$ -H<sub>6</sub>SAs at 200 GPa and (b)  $Fd\bar{3}m$ -H<sub>6</sub>STe at 200 GPa.

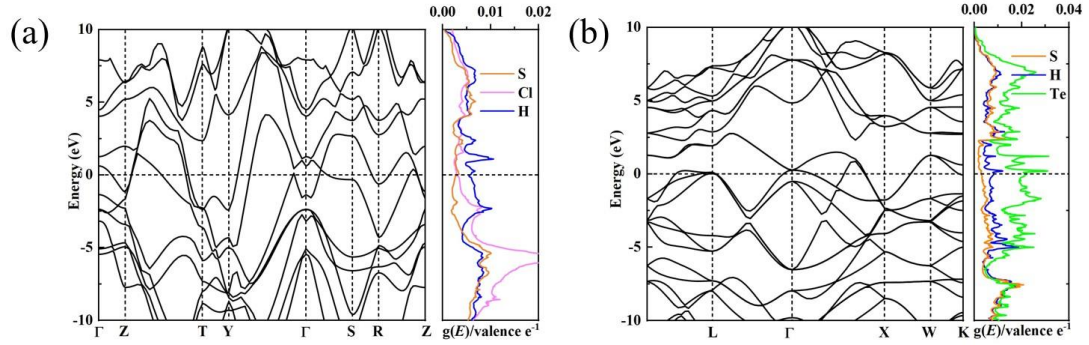


FIG. S5 Electronic band structure and DOS of (a)  $Cmmm$ -H<sub>6</sub>SCl at 130 GPa and (c)  $Fd\bar{3}m$ -H<sub>6</sub>STe at 100 GPa.

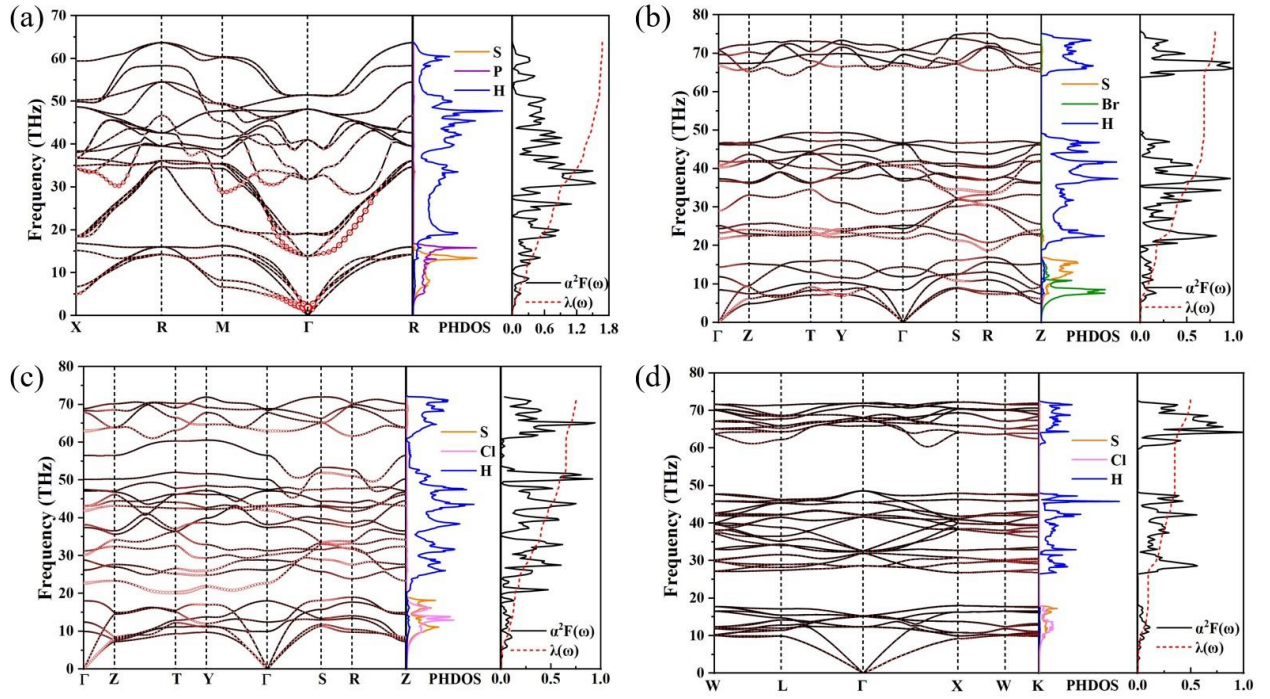


FIG. S6 Phonon spectra, phonon density of states, Eliashberg phonon spectral function  $\alpha^2F(\omega)$ , and the electron-phonon integral  $\lambda$  for (a)  $Pm\bar{3}m$ -H<sub>6</sub>SP at 200 GPa, (b)  $Cmmm$ -H<sub>6</sub>SBr at 200 GPa, (c)  $Cmmm$ -H<sub>6</sub>SCl at 200 GPa and (d)  $Fd\bar{3}m$ -H<sub>6</sub>SCl at 200 GPa.

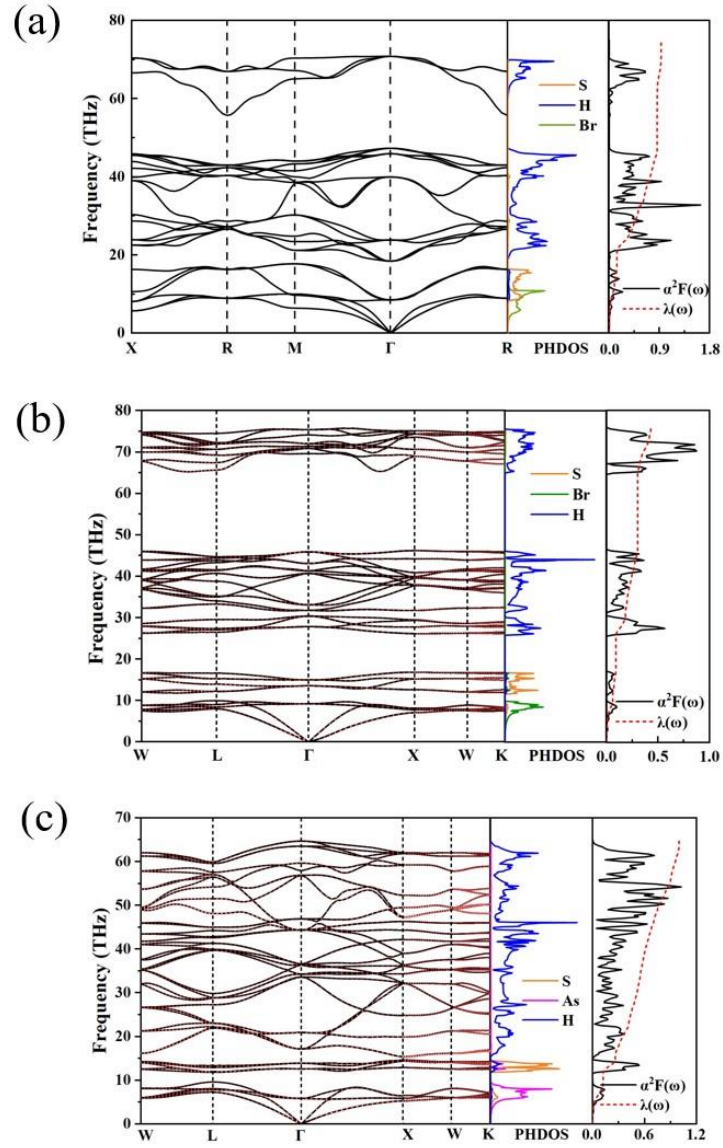


FIG. S7 Phonon spectra, phonon density of states, Eliashberg phonon spectral function  $\alpha^2F(\omega)$ , and the electron-phonon integral  $\lambda$  for (a)  $Pm\bar{3}m$ -H<sub>6</sub>SBr at 200 GPa, (b)  $Fd\bar{3}m$ -H<sub>6</sub>SBr at 200 GPa and (c)  $Fd\bar{3}m$ -H<sub>6</sub>SAs at 200 GPa.

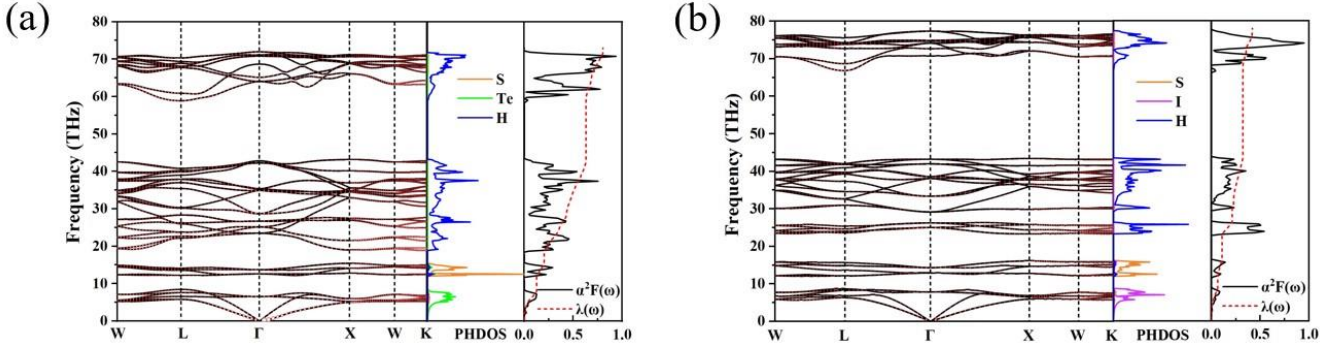


FIG. S8 Phonon spectra, phonon density of states, Eliashberg phonon spectral function  $\alpha^2F(\omega)$ , and the electron-phonon integral  $\lambda$  for (a)  $Fd\bar{3}m$ -H<sub>6</sub>STe at 200 GPa and (b)  $Fd\bar{3}m$ -H<sub>6</sub>SI at 200 GPa.

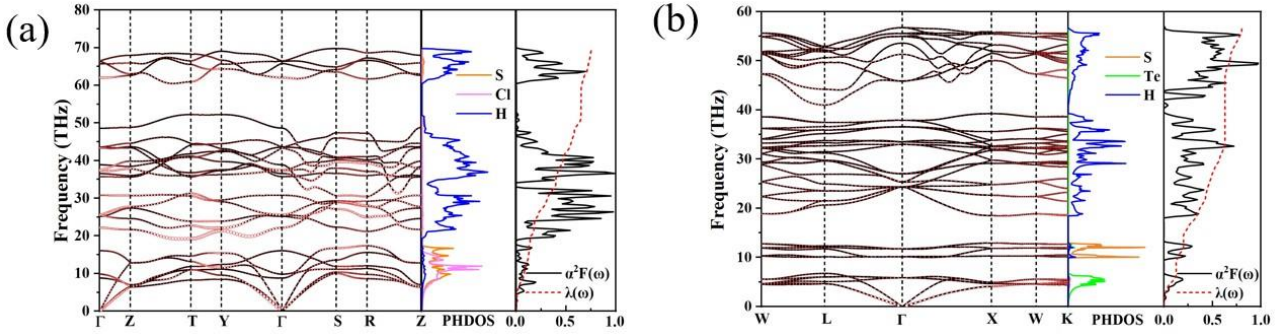


FIG. S9 Phonon spectra, phonon density of states, Eliashberg phonon spectral function  $\alpha^2F(\omega)$ , and the electron-phonon integral  $\lambda$  for (a)  $Cmmm$ -H<sub>6</sub>SCl at 130 GPa and (b)  $Fd\bar{3}m$ -H<sub>6</sub>STe at 100 GPa.

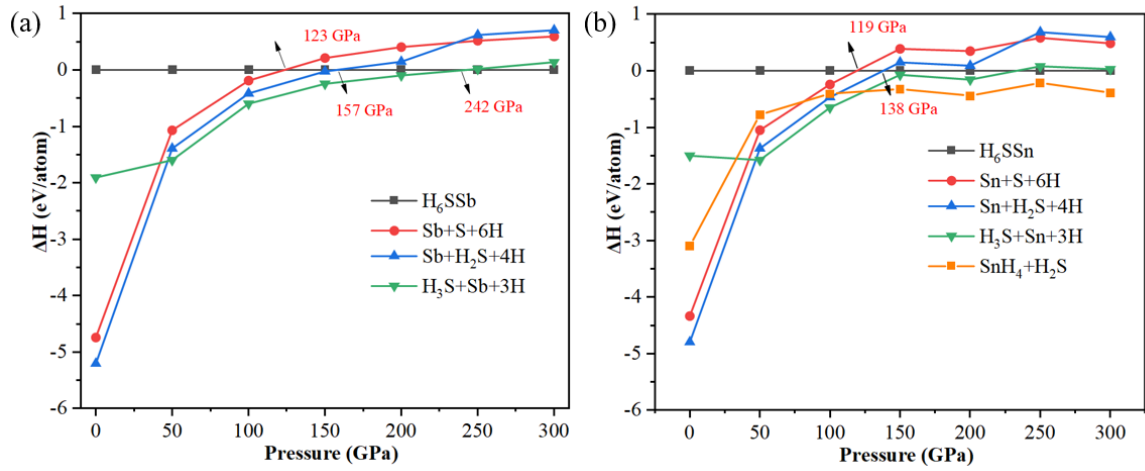


FIG. S10 Enthalpy difference vs pressure for (a)  $H_6SSb$ , referenced to the decomposition enthalpy into  $H_6SSb$ . The decomposition enthalpies into  $Sb + S + 6H$ ,  $Sb+H_2S+4H$  and  $H_3S+Sb+3H$  were also plotted, respectively. (b)  $H_6SSn$ , referenced to the decomposition enthalpy into  $H_6SSn$ . The decomposition enthalpies into  $Sn + S + 6H$ ,  $Sn+H_2S+4H$ ,  $H_3S+Sn+3H$  and  $SnH_4+H_2S$  were also plotted, respectively.

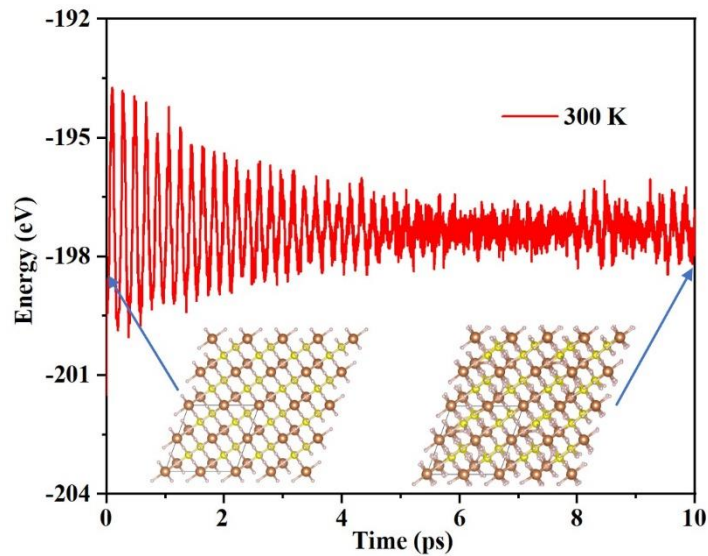


FIG. S11 the total free energy as a function of the ab initio molecular dynamics (AIMD) simulation time at 300 K under 200 GPa. The inset presents the structure snapshots of  $Fd\bar{3}m$   $H_6SSb$  at 0 and 10 ps, respectively.



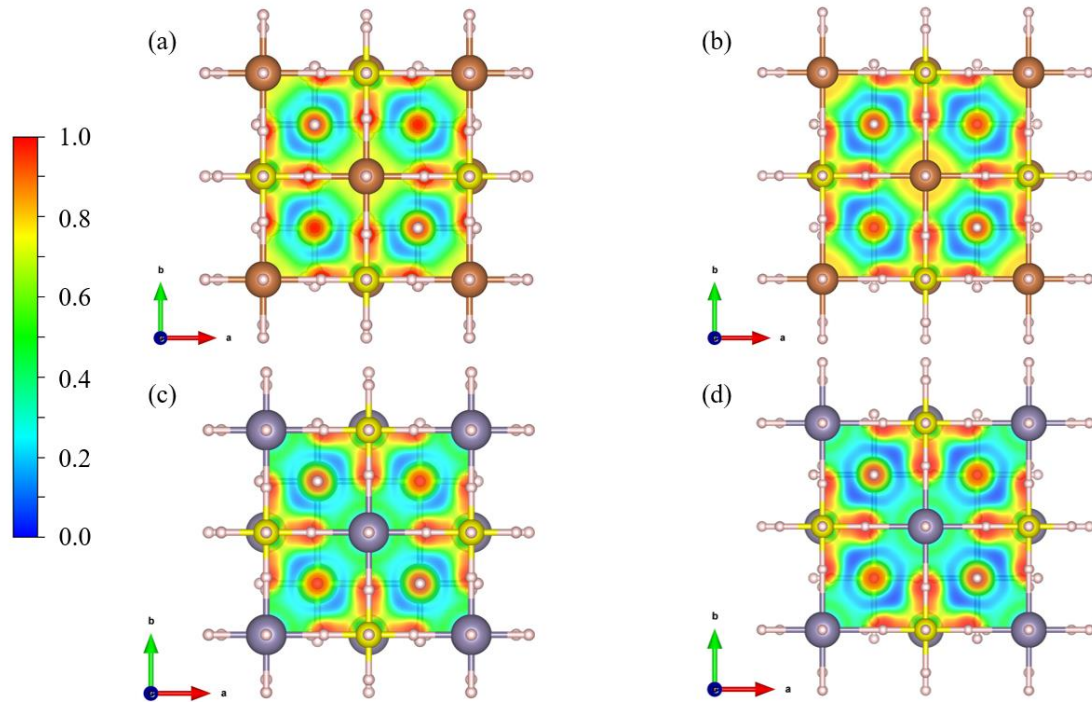


FIG. S12 Calculated electron localization function (ELF) of (a)  $Fd\bar{3}m$ -H<sub>6</sub>SSb at 200 GPa, (b)  $Fd\bar{3}m$ -H<sub>6</sub>SSb at 40 GPa, (c)  $Fd\bar{3}m$ -H<sub>6</sub>SSn at 200 GPa and (d)  $Fd\bar{3}m$ -H<sub>6</sub>SSn at 40 GPa.

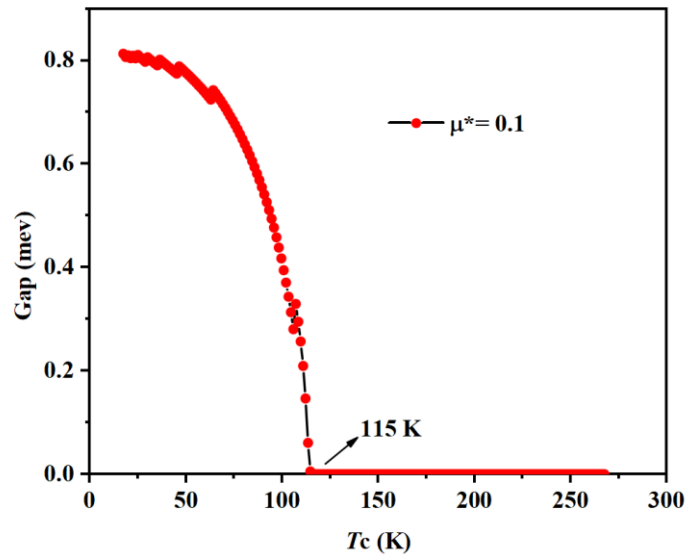


FIG. S13 Calculated anisotropic superconducting gap of H<sub>6</sub>SSb in the  $Fd\bar{3}m$  structure at 40 GPa.

Table. S1 The valence electrons of the atoms

<b>S</b>	<b>H</b>	<b>Sn</b>	<b>P</b>	<b>As</b>	<b>Sb</b>	<b>Te</b>	<b>Cl</b>	<b>Br</b>	<b>I</b>
$s^2p^4$	$S^1$	$s^2p^2$	$s^2p^3$	$s^2p^3$	$s^2p^3$	$s^2p^4$	$s^2p^5$	$s^2p^5$	$s^2p^5$