

SUPPORTING INFORMATION

for

Observation of Haldane Magnetism in Organically Templated Vanadium Phosphate (*enH*₂)_{0.5}VPO₄OH

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Synthesis and X-ray diffraction

Table S1. Chemicals used for the synthesis in deionized H₂O (5 μS cm⁻¹; Membrane-based deionizer)

Chemical	Purity, other notes	Supplier
Vanadium oxide (III) V ₂ O ₃	99.7% metals basis	Alfa Aesar
Phosphoric acid H ₃ PO ₄ (85 mass. % solution in H ₂ O)	ACS reagent, reducing substances, passes test ≤0.001% volatile acids (as CH ₃ COOH)	Sigma Aldrich
Ethylenediamine H ₂ N(CH ₂) ₂ NH ₂	99%	Ruskhim LLC

Table S2. Atomic coordinates and thermal parameters for (*enH*₂)_{0.5}VPO₄OH from Rietveld refinement

Atom	Wyckoff	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>
V1	4 <i>e</i>	0.0100(12)	0.1557(10)	0.2492(4)
P1	4 <i>e</i>	-0.5094(13)	-0.0934(19)	0.1515(4)
O1	4 <i>e</i>	0.149(4)	0.401(4)	0.1940(1)
O2	4 <i>e</i>	-0.318(3)	0.109(3)	0.1727(8)
O3	4 <i>e</i>	-0.743(3)	-0.119(3)	0.2112(9)
O4	4 <i>e</i>	-0.324(4)	-0.303(3)	0.1641(9)
O5	4 <i>e</i>	-0.674(2)	-0.091(3)	0.0769(6)
N1	4 <i>e</i>	-1.061(4)	0.233(4)	0.0446(10)
C1	4 <i>e</i>	-1.100(6)	-0.412(8)	0.0130(13)
H1	4 <i>e</i>	-1.21(3)	0.19(4)	0.020(8)
H2	4 <i>e</i>	-1.24(4)	-0.47(4)	0.05(1)

H3	4e	-1.18(6)	-0.35(4)	0.03(1)
H4	4e	-1.14(3)	0.27(3)	0.084(7)
H5	4e	-0.93(3)	0.12(3)	0.05(1)
H6	4e	0.31(4)	0.36(4)	0.180(9)

Bond lengths for the O-H (O1–H6 bond), N-H, C–H were fixed in the range 0.83-0.87 Å, 0.85-0.95 Å and 1.45-1.55 Å based on Refs. [1-3].

Comparison of $(enH_2)_{0.5}VPO_4OH$ and NH_4VPO_4OH crystal structures

The new compound, $(enH_2)_{0.5}VPO_4OH$, relates structurally to the previously described phase, NH_4VPO_4OH , but they are not isostructural. The first phase crystallizes in the $P2_1/c$ space group, while the second phase in the $P2_1/m$ space group. A comparative analysis of their crystal structures reveals similar two-dimensional structural fragments. These sheets are built up from the chains of edge-sharing VO_5OH octahedra strengthened by the PO_4 tetrahedra. The V-O chains run along the b axis in both crystal structures, as shown in the Fig. S1. The resulting V-P-O sheets are perpendicular to the a and c axes in NH_4VPO_4OH and $(enH_2)_{0.5}VPO_4OH$, respectively. In NH_4VPO_4OH , 2_1 -axes pass through vanadium atoms, so neighboring V-P-O sheets are connected by simple translation. In the crystal structure of $(enH_2)_{0.5}VPO_4OH$, the 2_1 -axes link symmetrically the neighboring V-P-O fragments, leading to a doubling of the c -axis compared to the a -axis of NH_4VPO_4OH , as shown in Fig. S1. This difference is due to the arrangement of larger ethylenediamine molecules in the interlayer space, compared to more compact ammonium groups.

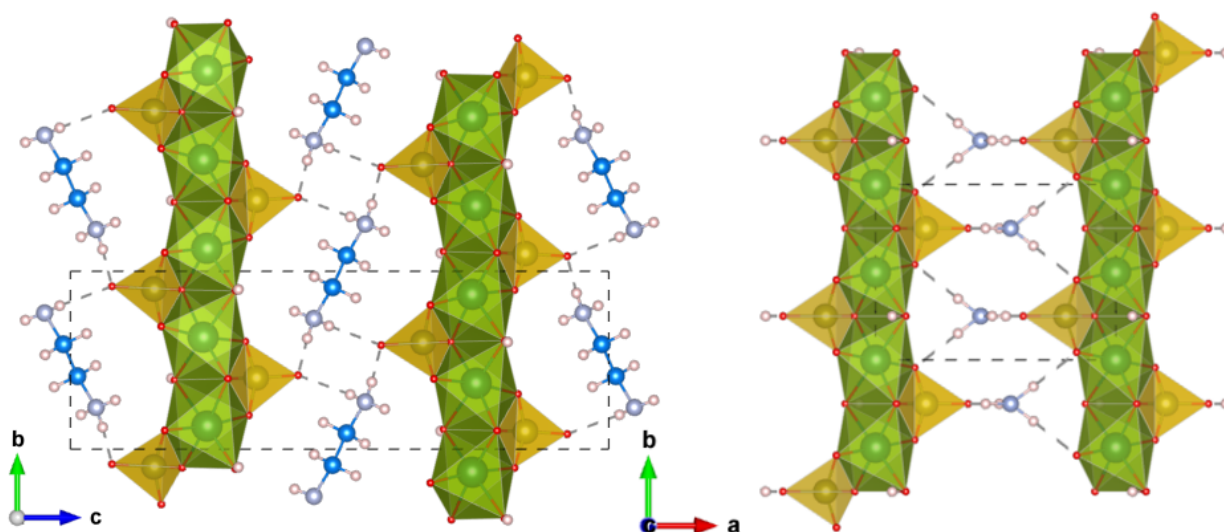


Fig. S1. Structurally related crystal structures of $(enH_2)_{0.5}VPO_4OH$ (left) and NH_4VPO_4OH (right), which are based on the 2D fragments of chains of edge-sharing VO_5OH octahedra strengthened by the PO_4 tetrahedra. V – light green, P – yellow, O – red, N – light blue, C – intense blue, H – pale pink.

Nuclear magnetic resonance

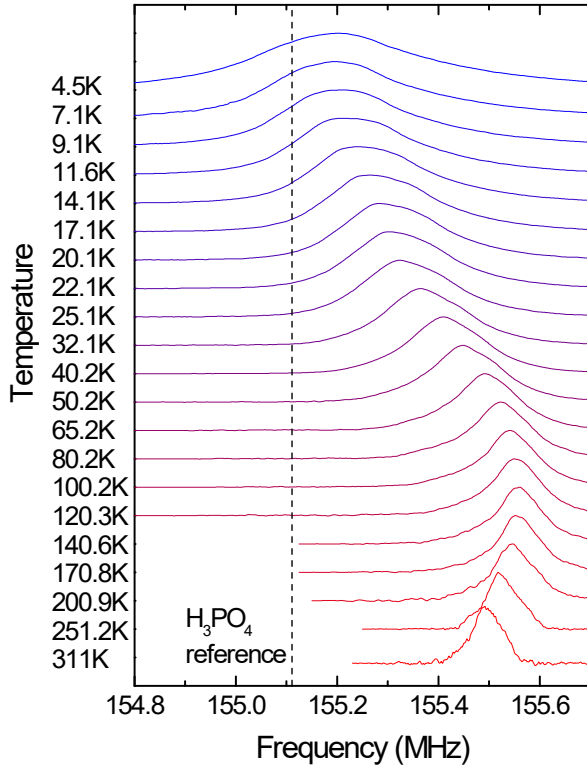


Figure S2. Evolution with temperature of the ^{31}P NMR spectra obtained at $\mu_0 H = 9$ T.

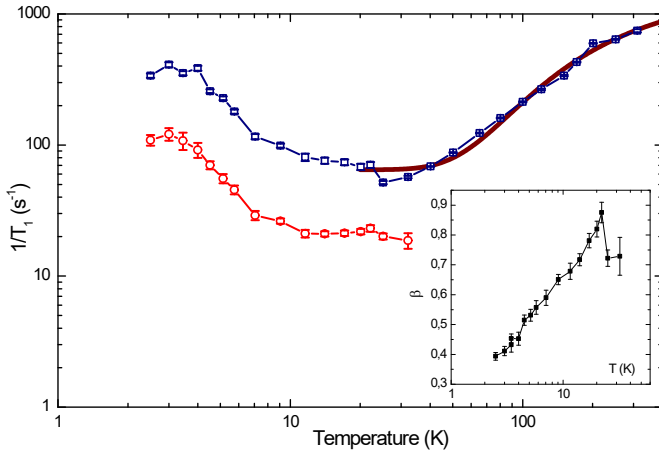


Figure S3. Temperature dependences of the fast and slow ^{31}P spin-lattice relaxation rates in $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$ measured at 9 T. Fast relaxation component is depicted by blue squares, with the slow component by red circles. Inset: stretching factor β of the slow relaxing component as a function of temperature.

As discussed in the main text, we attribute the fast-relaxing component $1/T_{1\text{fast}}$ of nuclear spin-lattice relaxation (NSLR) curves to the relaxation through the $S = 1$ spins of the Haldane chains. Above the spin-glass formation broad maximum, it would be reasonable to observe Haldane-like temperature behavior. In general, the NSLR rate due to magnetic centers can be written as [4]:

$$\frac{1}{T_1} = 2\gamma_n^2 k_B T \sum_q |A(q)|^2 \frac{\chi''(q, \omega_n)}{\omega_n},$$

where γ_n is the nuclear gyromagnetic ratio, $A(q)$ is the hyperfine coupling, and $\chi''(q, \omega_n)$ is the dissipative component of the dynamic susceptibility evaluated at the nuclear Larmor frequency ω_n and wave vector q .

Indeed, the data in the range of 20 – 311 K can be described by an activation law with an additional constant term $1/T_{\text{fast}} = 1/T_{\text{const}} + A \cdot \exp(-\Delta_R/k_B T)$ (the thick line on the **Fig. S3**), typical of gapped magnetic systems. The retrieved relaxation gap $\Delta_R = 229 \pm 18$ K is much higher than the value from the line shift data. However, the overestimation of the spin gap value in the NSLR experiments is common for Haldane systems; it arises at least from the dynamic susceptibility, involved in the relaxation process, as well as $A(q)$, while $K(T)$ and $\chi(T)$ deal only with the static one at $q = 0$ [5]. Various studies indicate different values for the relationship $\gamma = \Delta_R/\Delta$ between the relaxation gap Δ_R and the conventional gap $\Delta = 0.41J$ [3]. One of the most widely accepted value for γ is 1.5 [6-9]. Some theoretical studies estimate γ as 2 [10,11], and Pahari *et al.* even used the values of $3 = 1.5 \times 2$ [12]. However, even this approach leads to a slightly higher gap of $229/3$ K = 76 K (compared with 48.4 K obtained from ^{31}P NMR shift data). Some additional discrepancy may appear since most estimates of Δ_R are made for relatively low temperatures $k_B T < \Delta$, which can only be partially used in our case due to the spin-glass contribution. Nevertheless, it is important to note the observation of a spin gap also from the NSLR data.

The stretching factor β starts to decrease below 20 K, indicating an increase in the degree of disorder in the system (see the inset of Fig.S3). This result correlates with the increase in $1/T_1$ below 10 K and provides an additional piece of evidence for the emergence of a disordered spin glass state at temperatures around 3 K.

Density functional theory. Supplementary discussion

A. Spin exchanges J_1 and J_2 of $(en\text{H}_2)_{0.5}\text{VPO}_4\text{OH}$ and $\text{NH}_4\text{VPO}_4\text{OH}$

To evaluate the spin exchanges of $(en\text{H}_2)_{0.5}\text{VPO}_4\text{OH}$ and $\text{NH}_4\text{VPO}_4\text{OH}$ using the energy-mapping analysis [13-15], we carry out DFT+U calculations for the (a, 2b, c) supercell of $(en\text{H}_2)_{0.5}\text{VPO}_4\text{OH}$ and $\text{NH}_4\text{VPO}_4\text{OH}$ using the frozen core projector augmented plane wave (PAW) [16] encoded in the Vienna ab Initio Simulation Packages (VASP) [17] and the PBE potential [18] for the exchange-correlation functional. The electron correlation associated with the 3d states of Co was taken into consideration by DFT+U calculations with an effective on-site repulsion $U_{\text{eff}} = U - J = 2$ and 3 eV [19].

(1) Spin exchange paths

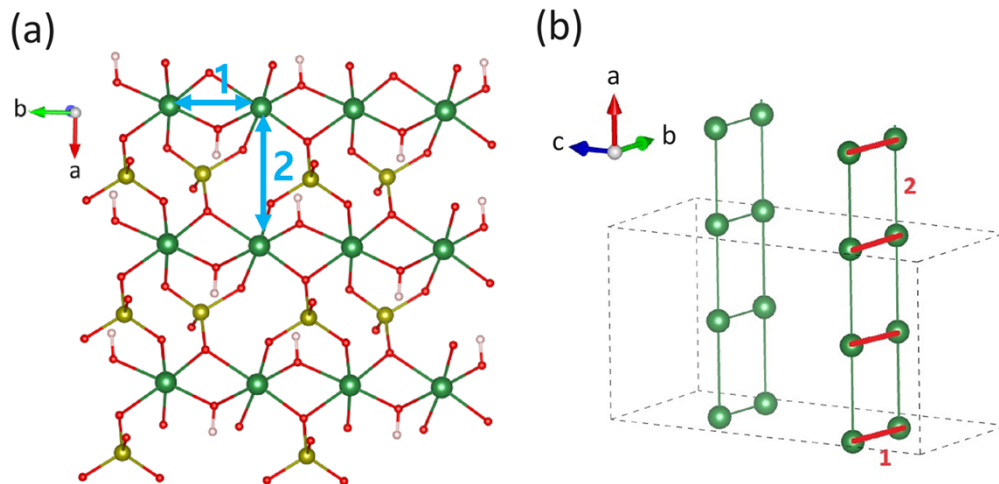


Figure S4. (a) A single VPO₄OH layer of $(en\text{H}_2)_{0.5}\text{VPO}_4\text{OH}$ showing how the edge-sharing ribbon chains are interconnected by the PO₄ groups to form a VPO₄OH layer, where the labels 1 and 2 refer to the intrachain and interchain spin exchange paths J_1 and J_2 , respectively. (b) Arrangement of the spin exchange paths J_1 and J_2 in $(en\text{H}_2)_{0.5}\text{VPO}_4\text{OH}$.

NOTE: The spin arrangement of **Fig. S4b** is also valid for $\text{NH}_4\text{VPO}_4\text{OH}$.

(2) Three ordered spin states used for the mapping analysis

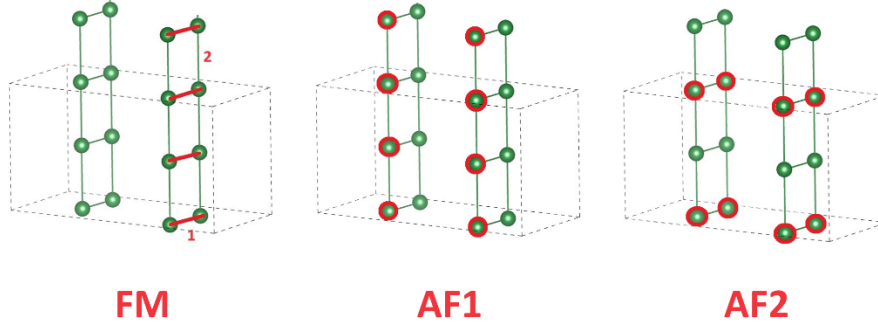


Figure S5. Spin arrangements of the ordered spin states FM, AF1, and AF2, where the green spheres without and with red circles represent up-spin and down-spin Co^{2+} ion sites, respectively.

(3) Energies of the ordered spin states in terms of the spin exchanges

$$\begin{aligned} E_{\text{FM}} &= (-8J_1 - 8J_2)S^2 \\ E_{\text{AF1}} &= (+8J_1 - 8J_2)S^2 \\ E_{\text{AF2}} &= (-8J_1 + 8J_2)S^2 \end{aligned}$$

(4) Spin exchanges in terms of the energies of the ordered spin states

$$\begin{aligned} J_1 &= (E_{\text{AF1}} - E_{\text{FM}})/(16S^2) \\ J_2 &= (E_{\text{AF2}} - E_{\text{AF1}})/(16S^2) \end{aligned}$$

(5) Relative energies of the ordered spin states in terms of DFT+U calculations

a) Computational parameters

Exchange-correlation functional: PBE
 Threshold for SCF convergence: 10^{-6} eV
 k-point: (6x8x2) for $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$
 (8x10x8) for $\text{NH}_4\text{VPO}_4\text{OH}$
 Plane wave cut-off energy: 450 eV

b) Results

Table S3. Relative energies (in meV/f.u.) obtained from DFT+U calculations

For $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$

	$U_{\text{eff}} = 2$ eV	$U_{\text{eff}} = 3$ eV
FM	25.31	20.91
AF1	0	0
AF2	25.37	20.85

For $\text{NH}_4\text{VPO}_4\text{OH}$

	$U_{\text{eff}} = 2 \text{ eV}$	$U_{\text{eff}} = 3 \text{ eV}$
FM	16.00	12.58
AF1	0	0
AF2	15.44	12.06

(6) Values of the spin exchanges

The values of J_1 and J_2 determined by the energy-mapping analyses are summarized in the main text.

B. Single-ion anisotropy of $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$ and $\text{NH}_4\text{VPO}_4\text{O}$

(1) Local Cartesian coordinates

To estimate the single-ion anisotropy of the V^{3+} ions in $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$ and $\text{NH}_4\text{VPO}_4\text{O}$, we first replace all but one V^{3+} ions of a unit cell with nonmagnetic ions Ga^{3+} to obtain a hypothetical structure that has one VO_6 octahedron containing a V^{3+} ion in a unit cell. Then we carry out DFT+U+SOC calculations [20] for this hypothetical structure using the local Cartesian coordinate defined for the VO_6 octahedron as defined in **Figure S6**, where we choose the local z-axis of the VO_6 octahedron along the shortest V-O bond (hereafter, the $\parallel z$ direction), and the $\perp z$ direction perpendicular to it.

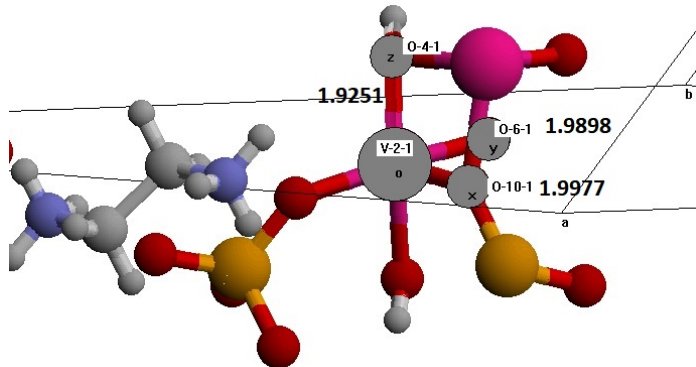


Figure S6. Local Cartesian coordinate system used for DFT+U+SOC calculations

(2) Computational parameters used

Exchange-correlation functional: PBE
 Threshold for SCF convergence: 10^{-6} eV
 k-point: (6x8x2) for $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$
 (6x8x10) for $\text{NH}_4\text{VPO}_4\text{OH}$
 Plane wave cut-off energy: 450 eV

(3) Results

Table S4. Relative energies ΔE (in K) of the $\parallel z$ and $\perp z$ spin orientations of $(\text{enH}_2)_{0.5}\text{VPO}_4\text{OH}$ and $\text{NH}_4\text{VPO}_4\text{O}$ obtained from DFT+U+SOC calculations as well as the spin and orbital moments μ_S and μ_L (in μ_B).

For $(enH_2)_{0.5}VPO_4OH$

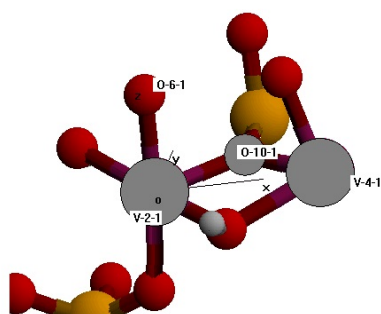
	$U_{\text{eff}} = 3 \text{ eV}$		$U_{\text{eff}} = 2 \text{ eV}$	
	//z	$\perp z$	//z	$\perp z$
ΔE	0	0.12	0	0.23
μ_S	1.890	1.890	1.874	1.874
μ_L	-0.019	-0.014	-0.023	-0.017

For NH_4VPO_4OH

	$U_{\text{eff}} = 3 \text{ eV}$		$U_{\text{eff}} = 2 \text{ eV}$	
	//z	$\perp z$	//z	$\perp z$
ΔE (K)	0	0.12	0	0.12
μ_S	1.891	1.891	1.875	1.875
μ_L	-0.023	-0.018	-0.028	-0.023

C. Partial density of states (PDOS) plots for $(enH_2)_{0.5}VPO_4OH$ and NH_4VPO_4OH showing only the up-spin t_{2g} states

(a)



(b)

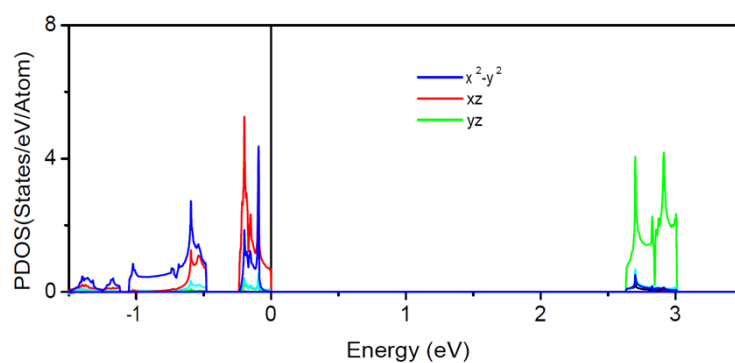


Figure S7. PDOS analysis for $(enH_2)_{0.5}VPO_4OH$: (a) The local Cartesian coordinates employed. (b) PDOS plots.

(a)

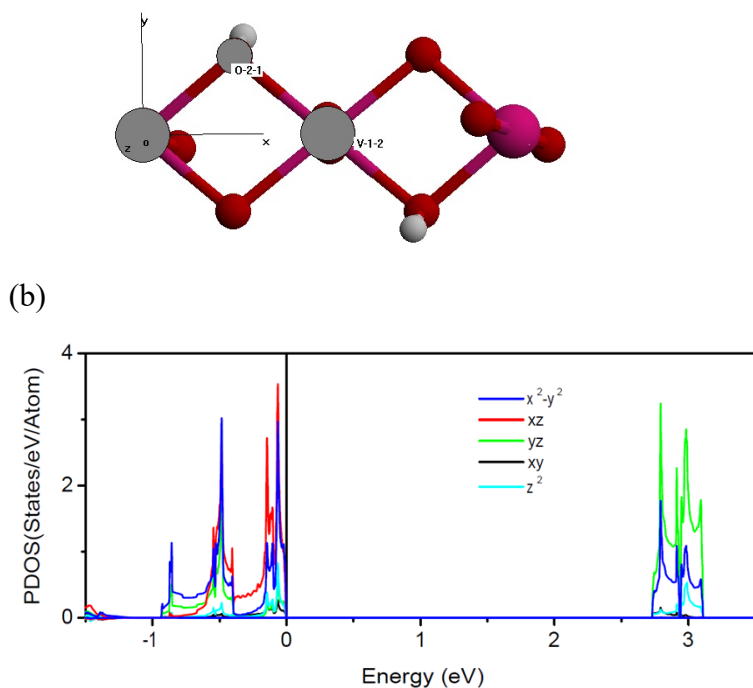


Figure S8. PDOS analysis for $\text{NH}_4\text{VPO}_4\text{OH}$: (a) The local Cartesian coordinates employed. (b) PDOS plots.

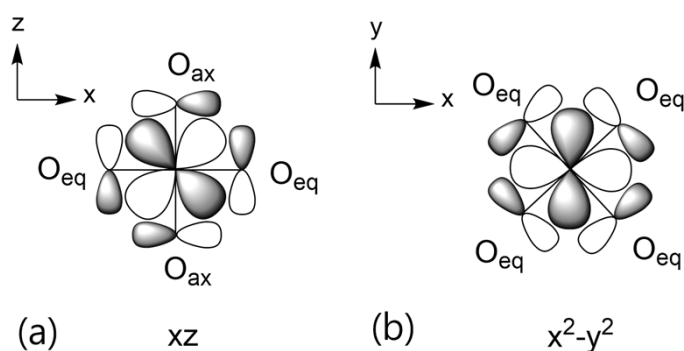


Figure S9. The p-orbitals of the O ligands that combine out-of-phase with the xz orbitals in (a), and with the x^2-y^2 orbitals in (b).

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