

Electronic Supplementary Information (ESI)
Isovalent cation substitution drives structural
transformation and infrared nonlinear optical activity in
Eu-based chalcogenides

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1 Experimental Section

1.1 Materials and Instruments

All reagents used in the present experiments were purchased from commercial sources and directly used without further purification. All weighing processes were completed in an anhydrous and oxygen-free glove box. The semi-quantitative energy dispersive X-ray (EDX, Oxford INCA) spectra were measured with a field emission scanning electron microscope (FESEM, JSM6700F). Powder X-ray diffraction (PXRD) analysis was carried out in a Rigaku Mini-Flex II powder diffractometer (Cu- K_{α} , $\lambda = 1.5418 \text{ \AA}$). UV-vis-NIR absorption measurement was performed in the region of 200–2500 nm at room temperature using an UV-vis-NIR spectrometer (Perkin-Elmer Lambda 950). The reflectance spectrum of the BaSO₄ powder was collected as the baseline and the diffuse reflectance data were converted to absorbance internally by the instrument by use of the Kubelka-Munk function.¹ The IR transmittance of the single-crystal sample was measured using the PerkinElmer Spectrum One FT-IR Spectrometer, within the range of 400–4000 cm⁻¹. The thermal stability analyses were measured on a NETZSCH STA 449C simultaneous analyser.

1.2 Synthesis

All reactants were stored and processed in a glove box filled with high-purity argon. To prepare β -EuZnGeS₄ crystals, the following materials—Eu₂O₃ powder (4N, Aladdin), Zn powder (4N, Aladdin), Ge powder (5N, Aladdin), S powder (4N, Aladdin), and B powder (4N, Aladdin)—were carefully weighed according to the stoichiometric ratio, resulting in a total mass of 500 mg. Elemental B in the reactive system acts as a reducing agent due to its strong affinity for oxygen. It not only extracts oxygen from metal oxides (e.g., Eu₂O₃) but also helps prevent the detrimental

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effects of air, water, and silica tubes on the synthesis process. The powders were then ground, pressed into tablets, and placed in a quartz crucible. The crucible was flame-sealed under a vacuum of 10^{-4} Torr and placed in a computer-controlled muffle furnace. The heating procedure was as follows: the temperature was gradually increased to 1173 K over 50 hours and held at that temperature for 100 hours. Afterward, the furnace was cooled at a rate of 4 K/h, and the tube furnace was shut off once the temperature dropped to 573 K, allowing it to cool naturally to room temperature. The resulting target crystals were red-brown in color and were rinsed with distilled water and ethanol. The synthesized compound retained its weight and color at room temperature for over six months, demonstrating its excellent physical and chemical stability.

1.3 Single-Crystal Structure determination

Taking the high-quality crystal of β -EuZnGeS₄ with suitable sizes was selected for single-crystal X-ray diffraction (XRD) analysis. The single-crystal diffraction data collections were collected on a Saturn 724 install with graphite-monochromated Mo- K_{α} radiation ($\lambda = 0.71073$ Å) at room temperature. The absorption correction was performed by the multi-scan method.² Using direct methods and making further refinement by full-matrix least-square fitting on F^2 based on *SHELX-2014* software, the precise structure was determined successfully.³ The atomic coordinates and equivalent isotropic displacement parameters are given in the Tables S1 and S2. Bond distances (Å) and bond angles (°) of RbPbPS₄ are listed in Table S3. CIF of β -EuZnGeS₄ has been submitted with CCDC number 2385100.

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1.4 Second-Harmonic Generation (SHG) Measurements

The powder SHG property test was investigated by Kurtz-Perry method⁴ using a Q-switched laser radiation. The laser radiation at 2900 nm was selected as the laser sources with the laser energy of 10 mJ and AgGaS₂ were measured as the benchmark, respectively. Samples of β -EuZnGeS₄ and AgGaS₂ were ground and sieved into different granule sizes (30–46, 46–74, 74–106, 106–150, and 150–210 μ m) for the phase matching measurements. The frequency-doubled output signals were detected via photomultiplier tube and oscilloscope.

1.5 Laser Induced Damage Threshold (LIDT) Measurements

The LIDT of β -EuZnGeS₄ at the maximal scope of 150–210 μ m was carried out single pulse measurement method⁵ and similar scope of AgGaS₂ single crystal used to the reference. The whole measuring materials were foist into selfsame plastic holders (thickness: 1 mm and diameter: 8 mm), respectively. Using an optical microscope monitor the exterior change of sample under the 1064 nm laser radiation with pulse width τ_p of 10 ns. Nova II sensor with a PE50-DIF-C energy sensor and a Vernier caliper was used for measuring the power of laser beam and the damage spot radius.

2 Computational Details

The DFT calculations have been performed using the *Vienna ab initio simulation package* (VASP)^{6–8} with the Perdew-Burke-Ernzerhof (PBE)⁹ exchange correlation functional. The projected augmented wave (PAW)¹⁰ potentials with the valence states 6s, 6p and 4f for Eu, 4s, 4p and 3d for Zn, 4s and 4p for Ge, and 3s and 3p for S, respectively, have been used. A Γ -centered $7 \times 7 \times 5$ Monkhorst-Pack grid for the Brillouin zone sampling¹¹ and a cutoff energy of 750 eV for the plane wave expansion

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were found to get convergent lattice parameters. The linear and nonlinear optical calculation was performed in the condition of a Monkhorst-Pack k -point mesh of $9 \times 9 \times 7$.

The imaginary part of the dielectric function due to direct inter-band transitions is given by the expression:

$$\varepsilon_2(\hbar\omega) = \frac{2e^2\pi}{\Omega\varepsilon_0} \sum_{k,v,c} \left| \langle \psi_k^c | u \cdot r | \psi_k^v \rangle \right|^2 \delta(E_k^c - E_k^v - E) \dots\dots\dots (1)$$

where Ω , ω , u , v and c are the unit-cell volume, photon frequencies, the vector defining the polarization of the incident electric field, valence and conduction bands, respectively. The real part of the dielectric function is obtained from ε_2 by a Kramers-Kronig transformation:

$$\varepsilon_1(\omega) = 1 + \left(\frac{2}{\pi}\right) \int_0^{+\infty} d\omega' \frac{\omega'^2 \varepsilon_2(\omega')}{\omega'^2 - \omega^2} \dots\dots\dots (2)$$

The refractive index $n(\omega)$ can be obtained based on ε_1 and ε_2 .

In calculation of the static $\chi^{(2)}$ coefficients, the so-called length-gauge formalism¹² derived by Aversa and Sipe¹³ and modified by Rashkeev et al¹⁴ is adopted, which has proven to be successful in calculating the second order susceptibility for semiconductors and insulators. In the static case, the imaginary part of the static second-order optical susceptibility can be expressed as:

$$\begin{aligned} & \chi^{abc} \\ &= \frac{e^3}{\hbar^2\Omega} \sum_{nml,k} \frac{r_{nm}^a (r_{ml}^b r_{ln}^c + r_{ml}^c r_{ln}^b)}{2\omega_{nm} \omega_{ml} \omega_{ln}} [\omega_n f_{ml} + \omega_m f_{ln} + \omega_l f_{nm}] \\ &+ \frac{ie^3}{4\hbar^2\Omega} \sum_{nm,k} \frac{f_{nm}}{\omega_{mn}^2} [r_{nm}^a (r_{mn;c}^b + r_{mn;b}^c) + r_{nm}^b (r_{mn;c}^a + r_{mn;a}^c) + r_{nm}^c (r_{mn;b}^a + r_{mn;a}^b)] \\ &\dots\dots\dots(3) \end{aligned}$$

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where r is the position operator, $\hbar\omega_{nm} = \hbar\omega_n - \hbar\omega_m$ is the energy difference for the bands m and n , $f_{mn} = f_m - f_n$ is the difference of the Fermi distribution functions, subscripts a , b , and c are Cartesian indices, and $r_{mn;a}^b$ is the so-called generalized derivative of the coordinate operator in k space,

$$r_{nm;a}^b = \frac{r_{nm}^a \Delta_{mn}^b + r_{nm}^b \Delta_{mn}^a}{\omega_{nm}} + \frac{i}{\omega_{nm}} \times \sum_l (\omega_{lm} r_{nl}^a r_{lm}^b - \omega_{nl} r_{nl}^b r_{lm}^a) \dots\dots\dots (4)$$

where $\Delta_{nm}^a = (p_{nn}^a - p_{mm}^a) / m$ is the difference between the electronic velocities at the bands n and m .

As the nonlinear optical coefficients is sensitive to the momentum matrix, much finer k-point grid and large amount of empty bands are required to obtain a convergent $\chi^{(2)}$ coefficient. The $\chi^{(2)}$ coefficients here were calculated from PBE wave functions and a scissor operator has been added to correct the conduction band energy (corrected to the experimental gap), which has proven to be reliable in predicting the second order susceptibility for semiconductors and insulators.

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3 Figures and Tables

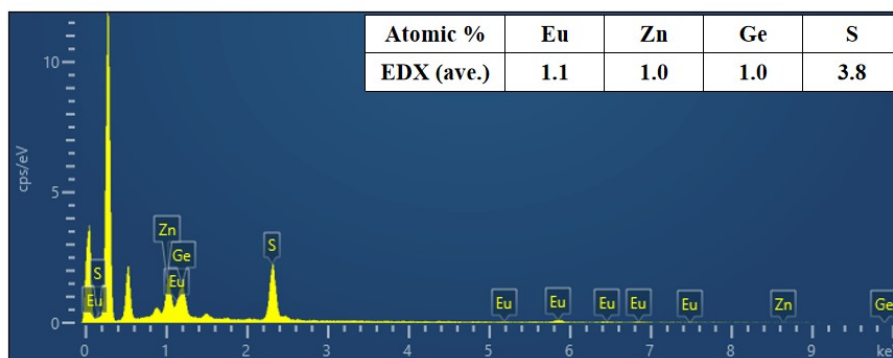


Figure S1. EDX results of β -EuZnGeS₄.

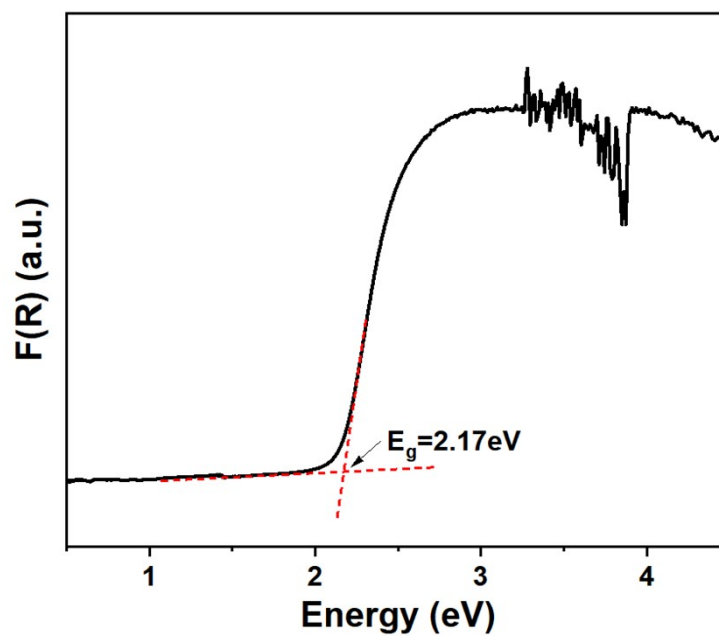


Figure S2. UV-vis-NIR diffuse reflectance spectrum of Eu₂GeS₄.

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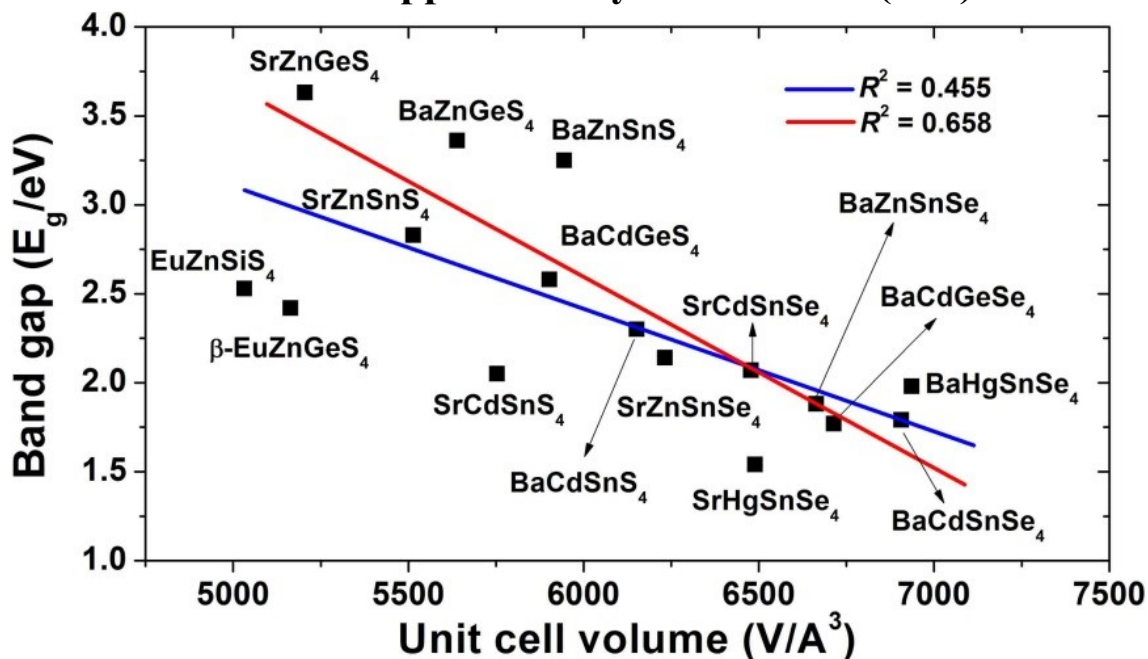


Figure S3. Plot of band gap energies in eV as a function of cell volumes in \AA^3 for the members of the $X^{II}\text{-}M^{II}\text{-}M^{IV}\text{-}Q_4$ family with the space group $Fdd2$.

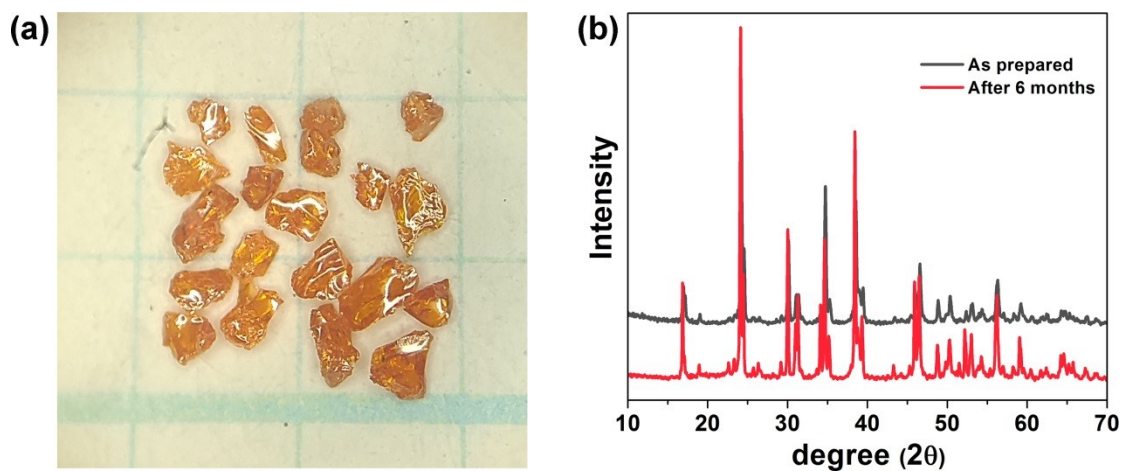


Figure S4. Crystal photo and powder XRD data of the $\beta\text{-EuZnGeS}_4$ compound stored for over 6 months.

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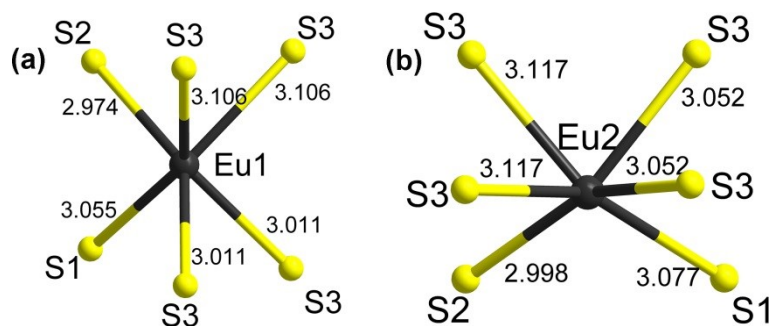


Figure S5. Coordination environment of (a) Eu1 and (b) Eu2 in ternary Eu_2GeS_4 .

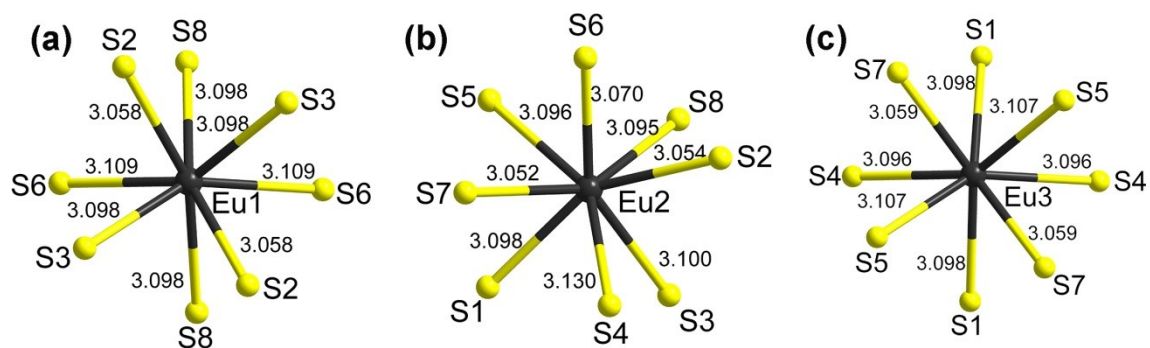


Figure S6. Coordination environment of (a) Eu1, (b) Eu2 and (c) Eu3 in quaternary $\beta\text{-EuZnGeS}_4$.

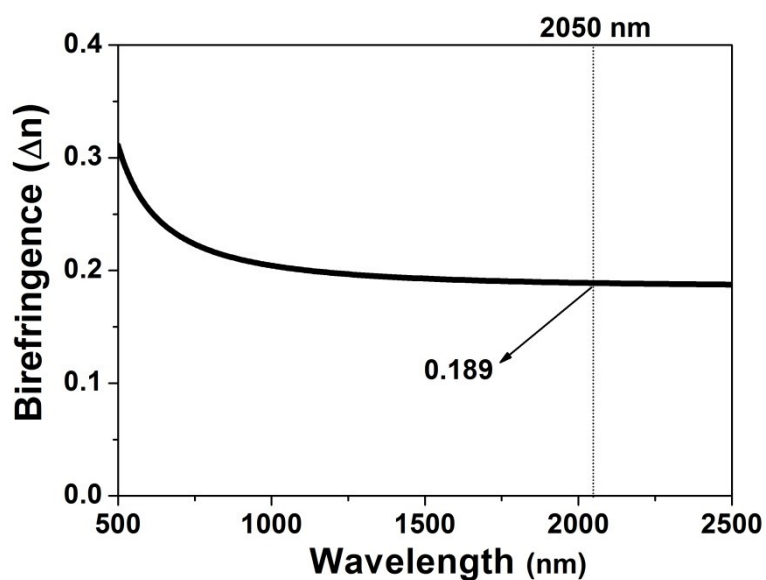


Figure S7. Calculated birefringence (Δn) of $\beta\text{-EuZnGeS}_4$.

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Table S1. Crystal data and structural refinement details for β -EuZnGeS₄.

Empirical formula	β -EuZnGeS ₄
Formula weight	418.16
Temperature/K	250 (2)
Crystal system	Orthorhombic
Space group	<i>Fdd2</i>
<i>a</i> /Å	20.6941(8)
<i>b</i> /Å	20.4059(6)
<i>c</i> /Å	12.2283(4)
α /°	90
β /°	90
γ /°	90
Volume/Å ³	5163.8(3)
<i>Z</i>	32
ρ_{calc} g/cm ³	4.30
μ /mm ⁻¹	19.1
Goodness-of-fit on F ²	1.083
Final R indexes [<i>I</i> >= 2 σ (<i>I</i>)]	$R_1 = 0.0187$, $wR_2 = 0.0450$
Final R indexes [all data]	$R_1 = 0.0263$, $wR_2 = 0.0474$
Largest diff. peak/hole / e Å ⁻³	0.64/-0.77
Flack parameter	-0.022(13)

^a: $R_1 = \Sigma||F_o| - |F_c||/\Sigma|F_o|$, $wR_2 = [\Sigma w(F_o^2 - F_c^2)^2/\Sigma w(F_o^2)^2]^{1/2}$

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Table S2. Atomic coordinates, equivalent isotropic displacement parameters and bond valence sums (BVS) for of β -EuZnGeS₄.

Atom	Wyckff	x	y	z	U_{eq} (Å) ^a	BVS*
Eu1	8a	0.75	0.25	0.89018(12)	0.0109(2)	2.04
Eu2	16b	0.50182(2)	0.24850(4)	0.89260(10)	0.0104(2)	2.06
Eu3	8a	0.75	0.25	0.39447(11)	0.0107(2)	2.04
Zn1	16b	0.62463(9)	0.32375(6)	0.10610(10)	0.0130(3)	1.88
Zn2	16b	0.62585(9)	0.36206(6)	0.63913(11)	0.0125(3)	1.91
Ge1	16b	0.62476(7)	0.38716(5)	0.39133(9)	0.0085(2)	4.00
Ge2	16b	0.62490(8)	0.18157(5)	0.68289(8)	0.0084(3)	4.03
S1	16b	0.6235(2)	0.20508(13)	0.5074(2)	0.0094(6)	2.01
S2	16b	0.6252(2)	0.27880(12)	0.7652(2)	0.0106(6)	2.12
S3	16b	0.70874(18)	0.12891(16)	0.7543(3)	0.0103(8)	2.47
S4	16b	0.70854(19)	0.37987(16)	0.5049(3)	0.0102(8)	1.90
S5	16b	0.54225(19)	0.37912(17)	0.5075(3)	0.0103(8)	1.92
S6	16b	0.54166(18)	0.12869(16)	0.7566(3)	0.0092(8)	1.94
S7	16b	0.6244(2)	0.29685(13)	0.2881(2)	0.0093(6)	2.10
S8	16b	0.6235(2)	0.46787(12)	0.2728(2)	0.0102(6)	2.02

U_{eq} is defined as one third of the trace of the orthogonalized U_{ij} tensor.

*The BVS was calculated by using empirical formula $V_i = \sum S_{ij} = \sum \exp[(r_0 - r_{ij})/0.37]$,

where S_{ij} is the bond valence associated with bond length r_{ij} , and r_0 .

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Table S3. Selected bond lengths (Å) and angle (°) of β -EuZnGeS₄.

Eu1–S2×2	3.058(4)	Zn1–S7	2.292(3)
Eu1–S8×3	3.098(4)	Zn1–S8×2	2.389(3)
Eu1–S3×2	3.098(3)	Zn1–S6×2	2.396(4)
Eu1–S6×2	3.109(3)	Zn1–S3×2	2.407(4)
Eu2–S7×2	3.053(4)	Zn2–S2	2.294(3)
Eu2–S2	3.054(4)	Zn2–S1×2	2.386(3)
Eu2–S6	3.070(3)	Zn2–S5	2.388(4)
Eu2–S8×2	3.095(4)	Zn2–S4	2.399(4)
Eu2–S5×2	3.096(3)	Ge1–S8	2.194(3)
Eu2–S1×2	3.098(4)	Ge1–S4	2.226(4)
Eu2–S3×2	3.100(3)	Ge1–S5	2.227(4)
Eu2–S4×2	3.130(3)	Ge1–S7	2.234(3)
Eu3–S7×2	3.059(4)	Ge2–S1	2.199(3)
Eu3–S4×2	3.096(3)	Ge2–S3	2.220(4)
Eu3–S1×2	3.098(4)	Ge2–S6	2.223(4)
Eu3–S5×2	3.107(3)	Ge2–S2	2.225(2)
\angle S2–Eu1–S8×2	66.76(7)	\angle S7–Eu3–S4×2	133.51(9)
\angle S2–Eu1–S8×2	156.64(5)	\angle S7–Eu3–S1×2	64.18(6)
\angle S2–Eu1–S3×2	79.13(10)	\angle S7–Eu3–S1×2	144.56(5)
\angle S2–Eu1–S3×2	69.66(8)	\angle S4–Eu3–S1×2	77.67(10)
\angle S8–Eu1–S3×2	132.62(9)	\angle S4–Eu3–S1×2	79.91(9)
\angle S8–Eu1–S3×2	83.38(8)	\angle S7–Eu3–S5	76.21(10)
\angle S2–Eu1–S6×2	128.70(9)	\angle S7–Eu3–S5	77.28(9)

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$\angle S2-Eu1-S6 \times 2$	80.35(8)	$\angle S4-Eu3-S5$	68.50(10)
$\angle S8-Eu1-S6 \times 2$	79.06(10)	$\angle S4-Eu3-S5$	146.80(6)
$\angle S8-Eu1-S6 \times 2$	73.16(9)	$\angle S1-Eu3-S5$	76.56(8)
$\angle S3-Eu1-S6 \times 2$	147.46(5)	$\angle S1-Eu3-S5$	135.47(9)
$\angle S3-Eu1-S6 \times 2$	69.40(10)	$\angle S7-Eu3-S5$	77.28(9)
$\angle S7-Eu2-S2$	124.28(10)	$\angle S7-Eu3-S5$	76.21(10)
$\angle S7-Eu2-S6$	76.20(10)	$\angle S4-Eu3-S5$	68.50(10)
$\angle S2-Eu2-S6$	70.14(8)	$\angle S4-Eu3-S5$	146.80(6)
$\angle S7-Eu2-S8$	150.18(8)	$\angle S1-Eu3-S5$	76.56(8)
$\angle S2-Eu2-S8$	66.84(7)	$\angle S1-Eu3-S5$	135.47(9)
$\angle S6-Eu2-S8$	83.92(8)	$\angle S7-Zn1-S8$	101.38(9)
$\angle S7-Eu2-S5$	71.52(8)	$\angle S7-Zn1-S6$	127.34(15)
$\angle S2-Eu2-S5$	130.39(10)	$\angle S8-Zn1-S6$	101.25(14)
$\angle S6-Eu2-S5$	69.83(11)	$\angle S7-Zn1-S3$	127.81(15)
$\angle S8-Eu2-S5$	80.92(10)	$\angle S8-Zn1-S3$	102.35(13)
$\angle S7-Eu2-S1$	64.26(6)	$\angle S6-Zn1-S3$	92.10(9)
$\angle S2-Eu2-S1$	150.36(7)	$\angle S2-Zn2-S1$	100.49(11)
$\angle S6-Eu2-S1$	135.56(10)	$\angle S2-Zn2-S5$	123.81(14)
$\angle S8-Eu2-S1$	121.77(10)	$\angle S1-Zn2-S5$	107.24(13)
$\angle S5-Eu2-S1$	78.75(9)	$\angle S2-Zn2-S4$	125.19(14)
$\angle S7-Eu2-S3$	132.07(10)	$\angle S1-Zn2-S4$	106.87(13)

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\angle S2–Eu2–S3	80.26(8)	\angle S5–Zn2–S4	91.92(11)
\angle S6–Eu2–S3	148.38(6)	\angle S8–Ge1–S4	118.12(14)
\angle S8–Eu2–S3	74.18(9)	\angle S8–Ge1–S5	117.88(14)
\angle S5–Eu2–S3	126.84(12)	\angle S4–Ge1–S5	101.20(12)
\angle S1–Eu2–S3	76.02(10)	\angle S8–Ge1–S7	104.23(11)
\angle S7–Eu2–S4	76.72(9)	\angle S4–Ge1–S7	107.45(13)
\angle S2–Eu2–S4	78.71(10)	\angle S5–Ge1–S7	107.31(13)
\angle S6–Eu2–S4	114.64(11)	\angle S1–Ge2–S3	119.96(15)
\angle S8–Eu2–S4	132.47(10)	\angle S1–Ge2–S6	119.42(15)
\angle S5–Eu2–S4	145.77(6)	\angle S3–Ge2–S6	102.20(10)
\angle S1–Eu2–S4	76.21(9)	\angle S1–Ge2–S2	104.30(9)
\angle S3–Eu2–S4	68.38(11)	\angle S3–Ge2–S2	104.59(15)
\angle S7–Eu3–S4 \times 2	71.50(8)	\angle S6–Ge2–S2	104.56(14)

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Table S4. The Δd (Å) and $\Delta\theta$ (deg) for the $X^{II}\text{-M}^{II}\text{-M}^{IV}\text{-Q}_4$ (X^{II} = Eu, Sr, Ba; M^{II} = Zn, Cd, Hg; M^{IV} = Si, Ge, Sn; Q = S, Se) family.

Number	Compound	Space group	Δd (Å)	$\Delta\theta$ (deg)	Ref.
1	BaZnSiSe ₄	<i>Ama2</i>	0.04	19.49	15
2	BaZnGeSe ₄	<i>Ama2</i>	0.06	19.65	15
3	SrZnGeS ₄	<i>Fdd2</i>	0.07	25.96	16
4	SrZnSnSe ₄	<i>Fdd2</i>	0.07	23.85	17
5	BaZnSnS ₄	<i>Fdd2</i>	0.08	26.02	18
6	SrZnSnS ₄	<i>Fdd2</i>	0.08	26.75	19
7	BaZnSnSe ₄	<i>Fdd2</i>	0.08	23.73	18
8	β-EuZnGeS₄	<i>Fdd2</i>	0.077	17.15	This work
9	EuZnSiS ₄	<i>Fdd2</i>	0.093	24.92	20
10	SrCdSnS ₄	<i>Fdd2</i>	0.1	28.57	21
11	SrCdSnSe ₄	<i>Fdd2</i>	0.1	25.95	21
12	BaCdSnSe ₄	<i>Fdd2</i>	0.1	24.97	22
13	BaCdSnS ₄	<i>Fdd2</i>	0.1	27.41	23
14	BaCdGeSe ₄	<i>Fdd2</i>	0.11	27.74	24
15	BaCdGeS ₄	<i>Fdd2</i>	0.16	31.34	25
16	SrHgSnSe ₄	<i>Fdd2</i>	0.16	15.9	26
17	BaHgSnSe ₄	<i>Fdd2</i>	0.16	16.8	26
18	SrCdSiS ₄	<i>Ama2</i>	0.13	31.61	27
19	SrCdGeSe ₄	<i>Ama2</i>	0.16	32.3	28
20	SrCdGeS ₄	<i>Ama2</i>	0.17	36.85	28
21	SrHgGeSe ₄	<i>Ama2</i>	0.18	29.27	29

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22	BaHgGeSe ₄	<i>Ama2</i>	0.19	28.73	29
23	SrZnSiSe ₄	<i>Ama2</i>	0.22	30.15	30
24	SrHgSiS ₄	<i>Ama2</i>	0.26	45.23	31
25	SrHgGeS ₄	<i>Ama2</i>	0.26	38.27	31
26	SrHgSnS ₄	<i>Ama2</i>	0.28	23.89	26
27	BaHgSiS ₄	<i>Ama2</i>	0.29	46.03	31
28	β -BaHgSnS ₄	<i>Ama2</i>	0.3	24.45	26
29	BaHgGeS ₄	<i>Ama2</i>	0.31	38.92	31
30	EuCdSiS ₄	<i>Ama2</i>	0.1718	44.314	20
31	EuCdGeS ₄	<i>Ama2</i>	0.172	37.11	32
32	EuCdGeSe ₄	<i>Ama2</i>	0.175	32.27	32
33	EuHgGeSe ₄	<i>Ama2</i>	0.221	34.43	33
34	EuHgGeS ₄	<i>Ama2</i>	0.248	38.983	34
35	EuHgSnS ₄	<i>Ama2</i>	0.282	23.87	33
36	α -BaHgSnS ₄	<i>Pnn2</i>	0.27	25.258	35

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