RSC Advances

# **Electronic Supplementary Information for**

# A Formally Exact Theory to Construct Nonreactive Forcefields using Linear Regression to Optimize Bonded Parameters

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## Contents

1. Supplementary tables

2. Analytic derivatives of the damped nonbonded potential

3. Analytic derivatives of the Manz stretch potential

4. Analytic derivatives of my new angle-bending potential

## **S1.** Supplementary tables

Table S1: Comparison of quantum-mechanically-computed Born-Oppenheimer energies (FCI/aug-cc-pVQZ) to those of the fitted forcefield for the H<sub>2</sub> singlet molecule. Results are shown for different values of the LASSO regularization parameter  $\lambda$ . The last row shows the root-mean-squared error (RMSE) in hartree. Please see Section 2.6.1 of the main text for additional details and discussion.

length	Esinglet	Etriplet	$U^{FF}$ - $E_{singlet}$	$U^{FF}$ - $E_{singlet}$	$U^{FF}$ - $E_{singlet}$	$U^{FF}$ - $E_{singlet}$	U <sup>FF</sup> -E <sub>singlet</sub>	$U^{FF}$ - $E_{singlet}$
(pm)	(hartree)	(hartree)	(λ=10-8)	(λ=10-9)	(λ=10 <sup>-10</sup> )	(λ=10 <sup>-11</sup> )	(λ=10 <sup>-12</sup> )	(λ=10 <sup>-20</sup> )
50	-1.10342	-0.57329	-0.00003	-0.00001	0.00000	0.00000	0.00000	0.00000
60	-1.15536	-0.69147	0.00015	0.00008	0.00002	-0.00001	-0.00002	0.00000
70	-1.17259	-0.76151	0.00000	0.00001	0.00001	0.00001	0.00001	-0.00002
74.199	-1.17387	-0.78459	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
80	-1.17190	-0.81221	-0.00002	0.00001	0.00001	0.00002	0.00002	-0.00002
90	-1.16189	-0.85111	-0.00017	0.00001	0.00003	0.00006	0.00007	0.00000
100	-1.14725	-0.88188	-0.00029	-0.00007	-0.00002	0.00002	0.00003	0.00005
125	-1.10531	-0.93467	0.00008	-0.00007	-0.00005	-0.00008	-0.00009	-0.00004
150	-1.06790	-0.96483	0.00036	0.00015	0.00010	0.00009	0.00008	0.00002
175	-1.04010	-0.98161	-0.00001	0.00004	0.00002	0.00004	0.00005	0.00005
200	-1.02193	-0.99066	-0.00034	-0.00014	-0.00010	-0.00009	-0.00008	-0.00004
225	-1.01130	-0.99539	-0.00027	-0.00010	-0.00006	-0.00006	-0.00006	-0.00005
250	-1.00558	-0.99779	0.00002	0.00003	0.00005	0.00004	0.00003	0.00001
275	-1.00268	-0.99896	0.00025	0.00011	0.00009	0.00008	0.00007	0.00004
300	-1.00125	-0.99951	0.00029	0.00008	0.00003	0.00003	0.00004	0.00003
350	-1.00023	-0.99986	0.00001	-0.00008	-0.00011	-0.00009	-0.00008	-0.00004
400	-0.99999	-0.99991	-0.00025	-0.00008	-0.00005	-0.00005	-0.00005	-0.00004
450	-0.99998	-0.99991	-0.00014	0.00008	0.00013	0.00011	0.00010	0.00006
500	-0.99991	-0.99991	0.00017	-0.00002	-0.00005	-0.00005	-0.00004	-0.00002
RMSE		_	2.0E-04	7.6E-05	6.3E-05	5.9E-05	5.8E-05	3.4E-05

$\lambda$ . The last row shows the sum of absolute values of the force constants. Please see Section 2.6.1 of the main text for additional details and discussion.										
	$\lambda = 10^{-8}$	λ=10-9	λ=10 <sup>-10</sup>	λ=10 <sup>-11</sup>	λ=10 <sup>-12</sup>	λ=10 <sup>-20</sup>				
k <sub>1</sub>	0.722095562	0.732963059	0.731056623	0.732239879	0.732721173	0.716352037				
k <sub>2</sub>	-0.493924304	-0.453665595	-0.437217721	-0.425469207	-0.422810820	-0.401964886				
k <sub>3</sub>	0.423826705	0.172199834	0.203826522	0.164324310	0.152145226	0.620093926				
<b>k</b> <sub>4</sub>	0.000000000	-0.220855956	-0.436476147	-0.552373758	-0.576139531	-1.588700146				
k5	0.000000000	1.547859193	1.439235648	1.755529569	1.830273684	-0.723790820				
k <sub>6</sub>	-1.370826079	-2.221089984	-1.326968080	-1.099867066	-1.070479469	6.738112940				
k <sub>7</sub>	0.000000000	-0.659979520	-0.700716437	-1.128691139	-1.116208404	-3.372565479				
k <sub>8</sub>	0.000000000	0.000000000	-1.068785502	-1.502396056	-1.627598580	-4.246902958				
k9	0.000000000	0.000000000	0.000000000	-0.249227598	-0.529246504	-5.135999922				
k <sub>10</sub>	0.000000000	0.000000000	0.000000000	0.383964516	0.554896362	-0.243272935				
k <sub>11</sub>	1.404363463	0.000000000	0.992454504	1.802334706	1.728244142	4.554823776				
k <sub>12</sub>	0.158317729	2.789790589	1.532480608	1.860919841	2.130693142	7.914267812				
k <sub>13</sub>	0.000000000	0.000000000	1.120819449	0.963001977	1.684902249	6.919472380				
k <sub>14</sub>	0.000000000	0.000000000	0.000000000	0.000000000	0.392860928	1.508486184				
k <sub>15</sub>	0.000000000	0.000000000	0.000000000	-0.013939077	-1.083956753	-6.637173720				
k <sub>16</sub>	0.000000000	0.000000000	0.000000000	-1.585807685	-2.432520787	-12.808342789				
k <sub>17</sub>	0.000000000	0.000000000	0.000000000	-1.349541481	-2.005604045	-9.225503143				
k <sub>18</sub>	-0.754863080	-1.915303544	-2.374382200	0.007312916	1.475379887	15.592033023				
abs sum	5.32821692	10.71370727	12.36441944	15.57694078	21.54668169	88.94785887				

Table S2: Computed force constants representing the H2 electronic singlet state's bonded interaction energy. Results are shown for different values of the LASSO regularization parameter

### S2. Analytic derivatives of the damped nonbonded potential

When evaluating derivatives for Cases #1, #2, and #3 described below, the following derivatives of the separation functions are useful:

$$\tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] = tanh \left[ \frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}} \right]$$
(S1)

$$\tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] = \tanh \left[ \frac{d_{AB}^{eq,j}}{d_{AB}} - \frac{d_{AB}}{d_{AB}^{eq,j}} \right]$$
(S2)

$$\vec{\nabla}_{C \notin \{A,B\}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] = 0$$
(S3)

$$\vec{\nabla}_{C \notin \{A,B\}} \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] = 0$$
(S4)

$$\vec{\nabla}_{\text{atom}_{-1}}\vec{\nabla}_{\text{C}\notin\{A,B\}}\tau_{AB}\left[d_{AB}, d_{\text{cutoff}}^{\text{nonbonded}}\right] = 0 \quad \text{for any atom}_{-1}$$
(S5)

$$\vec{\nabla}_{C\notin\{A,B\}}\vec{\nabla}_{atom_2}\tau_{AB}\left[d_{AB}, d_{cutoff}^{nonbonded}\right] = 0 \quad \text{for any atom}_2$$
(S6)

$$\vec{\nabla}_{atom_{l}}\vec{\nabla}_{C\notin\{A,B\}}\tau_{AB}\left[d_{AB},d_{AB}^{eq_{l}}\right] = 0 \quad \text{for any atom} \ 1$$
(S7)

$$\vec{\nabla}_{C \notin \{A,B\}} \vec{\nabla}_{atom_2} \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] = 0 \quad \text{for any atom}_2$$
(S8)

$$\vec{\nabla}_{A}\tau_{AB}\left[d_{AB}, d_{cutoff}^{nonbonded}\right] = \left(\frac{d_{cutoff}^{nonbonded}}{d_{AB}^{2}} + \frac{1}{d_{cutoff}^{nonbonded}}\right) \left(1 - \tanh^{2}\left[\frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}}\right]\right) \hat{R}_{AB}$$
(S9)

$$\hat{R}_{AB} = \left(\vec{R}_{B} - \vec{R}_{A}\right) / d_{AB}$$
(S10)

$$\vec{\nabla}_{B}\tau_{AB}\left[d_{AB}, d_{cutoff}^{nonbonded}\right] = -\vec{\nabla}_{A}\tau_{AB}\left[d_{AB}, d_{cutoff}^{nonbonded}\right]$$
(S11)

$$\vec{\nabla}_{A}\tau_{AB}\left[d_{AB}, d_{AB}^{eq,j}\right] = \left(\frac{d_{AB}^{eq,j}}{d_{AB}^{2}} + \frac{1}{d_{AB}^{eq,j}}\right) \left(1 - \tanh^{2}\left[\frac{d_{AB}^{eq,j}}{d_{AB}} - \frac{d_{AB}}{d_{AB}^{eq,j}}\right]\right) \hat{R}_{AB}$$
(S12)

$$\vec{\nabla}_{B}\tau_{AB}\left[d_{AB}, d_{AB}^{eq,j}\right] = -\vec{\nabla}_{A}\tau_{AB}\left[d_{AB}, d_{AB}^{eq,j}\right]$$
(S13)

$$\begin{split} \vec{\nabla}_{A}\vec{\nabla}_{A}\tau_{AB}\left[d_{AB},d_{cutoff}^{nonbonded}\right] &= \vec{\nabla}_{B}\vec{\nabla}_{B}\tau_{AB}\left[d_{AB},d_{cutoff}^{nonbonded}\right] = \\ & \left(\left(3\frac{d_{cutoff}^{nonbonded}}{d_{AB}^{a}} + \frac{1}{d_{cutoff}^{nonbonded}}d_{AB}\right)\left(1 - \tanh^{2}\left[\frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}}\right]\right)\hat{R}_{AB}\hat{R}_{AB} \\ & -2\left(\frac{d_{cutoff}^{nonbonded}}{d_{AB}^{2}} + \frac{1}{d_{cutoff}^{nonbonded}}\right)^{2} \tanh\left[\frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}}\right]\left(1 - \tanh^{2}\left[\frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}}\right]\right)\left(1 - \tanh^{2}\left[\frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}}\right]\right)\hat{R}_{AB}\hat{R}_{AB} \\ & -\frac{1}{d_{AB}}\left(\frac{d_{cutoff}^{nonbonded}}{d_{AB}^{2}} + \frac{1}{d_{cutoff}^{nonbonded}}\right)\left(1 - \tanh^{2}\left[\frac{d_{cutoff}^{nonbonded}}{d_{AB}} - \frac{d_{AB}}{d_{cutoff}^{nonbonded}}\right]\right)\vec{\delta} \end{split}$$

$$(S14)$$

where  $\vec{\delta}$  is the 3×3 identity tensor.

$$\vec{\nabla}_{A}\vec{\nabla}_{B}\tau_{AB}\left[d_{AB},d_{cutoff}^{nonbonded}\right] = \vec{\nabla}_{B}\vec{\nabla}_{A}\tau_{AB}\left[d_{AB},d_{cutoff}^{nonbonded}\right] = -\vec{\nabla}_{A}\vec{\nabla}_{A}\tau_{AB}\left[d_{AB},d_{cutoff}^{nonbonded}\right]$$
(S15)

$$\begin{split} \vec{\nabla}_{A}\vec{\nabla}_{A}\tau_{AB}\left[d_{AB},d_{AB}^{eq,j}\right] &= \vec{\nabla}_{B}\vec{\nabla}_{B}\tau_{AB}\left[d_{AB},d_{AB}^{eq,j}\right] = \\ & \left(\left(3\frac{d_{AB}^{eq,j}}{d_{AB}^{-3}} + \frac{1}{d_{AB}^{eq,j}d_{AB}}\right) \left(1 - \tanh^{2}\left[\frac{d_{AB}^{eq,j}}{d_{AB}} - \frac{d_{AB}}{d_{AB}^{eq,j}}\right]\right) \hat{R}_{AB}\hat{R}_{AB} \\ & -2\left(\frac{d_{AB}^{eq,j}}{d_{AB}^{-2}} + \frac{1}{d_{AB}^{eq,j}}\right)^{2} \tanh\left[\frac{d_{AB}^{eq,j}}{d_{AB}} - \frac{d_{AB}}{d_{AB}^{eq,j}}\right] \left(1 - \tanh^{2}\left[\frac{d_{AB}^{eq,j}}{d_{AB}} - \frac{d_{AB}}{d_{AB}^{eq,j}}\right]\right) \hat{R}_{AB}\hat{R}_{AB} \\ & -\frac{1}{d_{AB}}\left(\frac{d_{AB}^{eq,j}}{d_{AB}^{-2}} + \frac{1}{d_{AB}^{eq,j}}\right) \left(1 - \tanh^{2}\left[\frac{d_{AB}^{eq,j}}{d_{AB}} - \frac{d_{AB}}{d_{AB}^{eq,j}}\right]\right) \vec{\delta} \\ & \vec{\nabla}_{A}\vec{\nabla}_{B}\tau_{AB}\left[d_{AB}, d_{AB}^{eq,j}\right] = \vec{\nabla}_{B}\vec{\nabla}_{A}\tau_{AB}\left[d_{AB}, d_{AB}^{eq,j}\right] = -\vec{\nabla}_{A}\vec{\nabla}_{A}\tau_{AB}\left[d_{AB}, d_{AB}^{eq,j}\right] \tag{S17}$$

In the below formulas, atom\_1 is any chosen atom in the material, and atom\_2 is any chosen atom in the material. *Note*: atom\_2 may either be the same atom or a different atom than atom\_1.

*Case* # 1: The two atoms A and B are inside the same bonded cluster j and a cutoff distance is used for their nonbonded interaction. In this case, we express the effective multibody pairwise potentials as follows:

$$\Phi_{ABx}^{intercluster} = \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq, j} \right] \left( U_{ABx, intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx, intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \right)$$
(S18)

$$\Phi_{ABx}^{intracluster} = \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right)$$
(S19)

 $\Theta_{H}$  is the Heaviside step function, and  $d_{AB}^{eq_{L}j}$  is the equilibrium distance between atoms A and B in the isolated bonded cluster j. The first-order and second-order derivatives expand as follows:

$$\begin{split} \vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intercluster} &= \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \\ \begin{pmatrix} 3\tau_{AB}^{2} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \\ + 2\tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \\ + \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \vec{\nabla}_{atom_{-1}} U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \\ \end{pmatrix} \end{split}$$

$$(S20)$$

$$\begin{split} \vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intracluster} &= \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \\ \begin{pmatrix} 3\tau_{AB}^{2} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \\ &+ 2\tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \\ &+ \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \vec{\nabla}_{atom_{-1}} U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \\ & \end{pmatrix} \end{split}$$

$$(S21)$$

$$\begin{split} \vec{\nabla}_{atom_{\perp}} \vec{\nabla}_{atom$$

$$\begin{split} \bar{\nabla}_{atom_{-2}} \Phi_{ABx}^{instructure} &= \Theta_{H} \Big[ d_{cutoff}^{conbinded} - d_{AB} \Big] \\ 3\tau_{AB}^{-2} \Big[ d_{AB}, d_{cutoff}^{outomined} \Big] \tau_{AB} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( U_{ABx,intrachuster}^{ABx,intrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] - U_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \Big) \bar{\nabla}_{atom_{-2}} \tau_{AB} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \\ + 2\tau_{AB}^{-3} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( U_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] - U_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \Big) \bar{\nabla}_{atom_{-2}} \tau_{AB} \Big[ d_{AB}, d_{Cutoff}^{cut} \Big] \\ + \tau_{AB}^{-3} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \bar{\nabla}_{atom_{-1}} \bar{\nabla}_{atom_{-2}} U_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \\ + 3\tau_{AB}^{-2} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( \bar{\nabla}_{atom_{-1}} \bar{U}_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \\ + 3\tau_{AB}^{-2} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( \bar{\nabla}_{atom_{-1}} \bar{U}_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \Big) \bar{\nabla}_{atom_{-2}} \tau_{AB}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \\ + 2\tau_{AB}^{-3} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( \bar{\nabla}_{atom_{-1}} \bar{U}_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \\ + 2\tau_{AB}^{-3} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( \bar{\nabla}_{atom_{-1}} \bar{U}_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \\ + 2\tau_{AB}^{-3} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( \bar{\nabla}_{atom_{-1}} \bar{U}_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \Big) \Big( \bar{\nabla}_{atom_{-2}} \bar{U}_{ABx,intrachuster}^{antrachuster} \Big[ \Big\{ \vec{R}_{C} \Big\} \Big] \\ + 2\tau_{AB}^{-3} \Big[ d_{AB}, d_{cutoff}^{cut} \Big] \tau_{AB}^{-2} \Big[ d_{AB}, d_{AB}^{cut} \Big] \Big( U_{ABx,intrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachuster}^{antrachust$$

*Case* # 2: The two atoms A and B are inside the same bonded cluster j and a cutoff distance is not used for their nonbonded interaction. In this case, we express the effective multibody pairwise potentials as follows:

$$\Phi_{ABx}^{\text{intercluster}} = \tau_{AB}^{2} \left[ d_{AB}^{}, d_{AB}^{\text{eq},j} \right] \left( U_{ABx,\text{intercluster}}^{\text{nonbonded}} \left[ \left\{ \vec{R}_{C}^{} \right\} \right] - U_{ABx,\text{intercluster}}^{\text{nonbonded}} \left[ \left\{ \vec{R}_{C}^{eq}^{}, j \right\} \right] \right)$$
(S24)

$$\Phi_{ABx}^{intracluster} = \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx, intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx, intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right)$$
(S25)

The first-order and second-order derivatives expand as follows:

$$\vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intercluster} = \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right]$$

$$\begin{pmatrix} 2\tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \\ + \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \vec{\nabla}_{atom_{-1}} U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right]$$

$$(S26)$$

$$\vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intracluster} = \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \\ \left( 2\tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \left( U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{AB}^{eq,j} \right] \right) \\ \left( + \tau_{AB}^{2} \left[ d_{AB}, d_{AB}^{eq,j} \right] \vec{\nabla}_{atom_{-1}} U_{ABx,intracluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \right) \right)$$
(S27)

$$\begin{split} \vec{\nabla}_{atom_{\perp}} \vec{\nabla}_{atom_{\perp}2} \Phi_{ABx}^{intercluster} &= \Theta_{H} \Big[ d_{cutoff}^{innobinded} - d_{AB} \Big] \\ & \left\{ 2\tau_{AB} \Big[ d_{AB}, d_{AB}^{e,j} \Big] \Big( U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] - U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C}^{e,j} \right\} \Big] \right) \vec{\nabla}_{atom_{\perp}1} \vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] \vec{\nabla}_{atom_{\perp}1} U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] \\ & + \tau_{AB}^{2} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] \vec{\nabla}_{atom_{\perp}1} U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] \\ & + 2\tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] (\vec{\nabla}_{atom_{\perp}1} U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] ) \vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] (\vec{\nabla}_{atom_{\perp}1} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] ) \vec{\nabla}_{atom_{\perp}2} U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] \\ & + 2 \Big( U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] - U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C}^{eq,j} \right\} \Big] \right) (\vec{\nabla}_{atom_{\perp}1} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] ) \vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] \\ & \vec{\nabla}_{atom_{\perp}1} \vec{\nabla}_{atom_{\perp}2} \Phi_{ABx}^{nintercluster} \Big[ \left\{ \vec{R}_{C} \right\} \Big] - U_{ABx,intercluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C}^{eq,j} \right\} \Big] \right) (\vec{\nabla}_{atom_{\perp}1} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] (\vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] \\ & \vec{\nabla}_{atom_{\perp}1} \vec{\nabla}_{atom_{\perp}2} \Phi_{ABx}^{nintercluster} \Big[ \left\{ \vec{R}_{C} \right\} \Big] - U_{ABx,intracluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C}^{eq,j} \right\} \Big] ) (\vec{\nabla}_{atom_{\perp}1} \vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] \\ & + \tau_{AB}^{2} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] (\vec{\nabla}_{atom_{\perp}1} U_{ABx,intracluster}^{nonbonded} \Big[ \left\{ \vec{R}_{C} \right\} \Big] \right) (\vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] \\ & + \tau_{AB}^{2} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] (\vec{\nabla}_{atom_{\perp}1} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] ) (\vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] ) (\vec{\nabla}_{atom_{\perp}2} \tau_{AB} \Big[ d_{AB}, d_{AB}^{eq,j} \Big] ) \\ & + 2 \Big( u_{ABx,intracluster}^{2} \Big[ \left\{ \vec{R}_{C} \right\} \Big] - u_{ABx,intracluster}^{2} \Big[ \left\{ \vec{R}_{C} \right\} \Big] ) (\vec{\nabla}_{atom_{\perp}2} \tau_{AB$$

*Case # 3*: The two atoms A and D are not inside the same bonded cluster and a cutoff distance is used for their nonbonded interaction. In this case, we express the effective multibody pairwise potentials as follows:

$$\Phi_{ADx}^{intercluster} = \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right]$$
(S30)  
$$\Phi_{ADx}^{intracluster} = 0$$
(S31)

The first-order and second-order derivatives expand as follows:

$$\begin{split} \vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intercluster} &= \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \\ \left( 3\tau_{AB}^{2} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \left( U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq,j} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \right] \\ + \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \vec{\nabla}_{atom_{-1}} U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \\ \vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intracluster} = 0 \end{split}$$
(S33)

$$\begin{split} \vec{\nabla}_{atom_{-1}} \vec{\nabla}_{atom_{-2}} \Phi_{ABx}^{intercluster} &= \Theta_{H} \left[ d_{cutoff}^{nonbonded} - d_{AB} \right] \\ \left( 3\tau_{AB}^{2} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \left( U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \right) \vec{\nabla}_{atom_{-1}} \vec{\nabla}_{atom_{-2}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \\ &+ \tau_{AB}^{3} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \vec{\nabla}_{atom_{-1}} \vec{\nabla}_{atom_{-2}} U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \\ &+ 3\tau_{AB}^{2} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \left( \vec{\nabla}_{atom_{-1}} U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] \right) \vec{\nabla}_{atom_{-2}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \\ &+ 3\tau_{AB}^{2} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \left( \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \right) \vec{\nabla}_{atom_{-2}} U_{ABx,intercluster}^{nonbonded} \\ &+ 6\tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \left( U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right] - U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C}^{eq-j} \right\} \right] \right) \left( \vec{\nabla}_{atom_{-1}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \right) \vec{\nabla}_{atom_{-2}} \tau_{AB} \left[ d_{AB}, d_{cutoff}^{nonbonded} \right] \\ &\quad \vec{\nabla}_{atom_{-1}} \vec{\nabla}_{atom_{-2}} \Phi_{ADx}^{intracluster} = 0 \end{aligned}$$

$$(S35)$$

*Case* # 4: The two atoms A and D are not inside the same bonded cluster and a cutoff distance is not used for their nonbonded interaction. In this case, we express the effective multibody pairwise potentials as follows:

$$\Phi_{ADx}^{intercluster} = U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right]$$
(S36)

$$\Phi_{ADx}^{intracluster} = 0 \tag{S37}$$

The first-order and second-order derivatives expand as follows:

$$\vec{\nabla}_{atom_{-1}} \Phi_{ABx}^{intercluster} = \vec{\nabla}_{atom_{-1}} U_{ABx,intercluster}^{nonbonded} \left[ \left\{ \vec{R}_{C} \right\} \right]$$
(S38)

$$\vec{\nabla}_{\text{atom}_{-1}} \Phi_{\text{ADx}}^{\text{intracluster}} = 0 \tag{S39}$$

$$\vec{\nabla}_{atom_{-1}}\vec{\nabla}_{atom_{-2}}\Phi^{intercluster}_{ABx} = \vec{\nabla}_{atom_{-1}}\vec{\nabla}_{atom_{-2}}U^{nonbonded}_{ABx,intercluster}\left[\left\{\vec{R}_{C}\right\}\right]$$
(S40)

$$\vec{\nabla}_{\text{atom}_{-1}}\vec{\nabla}_{\text{atom}_{-2}}\Phi_{\text{ADx}}^{\text{intracluster}} = 0 \tag{S41}$$

### S3. Analytic derivatives of the Manz stretch potential

$$U_{AB}^{Manz\_stretch} \left[ d_{AB} \right] = \frac{3k_{AB}}{5\gamma_{AB}^{\circ 2}} \left( 1 - \left( \frac{5}{2} \right) exp \left[ -\gamma_{AB}^{\circ} \left( d_{AB} - d_{AB}^{ref} \right) \right] + \left( \frac{3}{2} \right) exp \left[ -\frac{5}{3}\gamma_{AB}^{\circ} \left( d_{AB} - d_{AB}^{ref} \right) \right] \right)$$
(S42)

$$\frac{dU_{AB}^{Manz\_stretch}\left[d_{AB}\right]}{dd_{AB}} = \frac{3k_{AB}}{2\gamma_{AB}^{\circ}} \left(exp\left[-\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right] - exp\left[-\frac{5}{3}\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right]\right)$$
(S43)

$$\frac{d^{2}U_{AB}^{Manz\_stretch}\left[d_{AB}\right]}{dd_{AB}^{2}} = k_{AB}\left(-\frac{3}{2}exp\left[-\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right] + \frac{5}{2}exp\left[-\frac{5}{3}\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right]\right)$$
(S44)

$$\frac{d^{3}U_{AB}^{Manz\_stretch}\left[d_{AB}\right]}{dd_{AB}^{3}} = k_{AB}\gamma_{AB}^{\circ}\left(\frac{3}{2}exp\left[-\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right] - \frac{25}{6}exp\left[-\frac{5}{3}\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right]\right)$$
(S45)

$$\frac{d^{4}U_{AB}^{Manz\_stretch}\left[d_{AB}\right]}{dd_{AB}^{4}} = k_{AB}\gamma_{AB}^{\circ}{}^{2}\left(-\frac{3}{2}\exp\left[-\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right] + \frac{125}{18}\exp\left[-\frac{5}{3}\gamma_{AB}^{\circ}\left(d_{AB}-d_{AB}^{ref}\right)\right]\right)$$
(S46)

### S4. Analytic derivatives of my new angle-bending potential

Consider a bond angle defined by the atoms A, B, and C, where B is the middle atom. Swope and Ferguson<sup>S1</sup> and Dubbeldam et al.<sup>S2</sup> gave universal formulas for first and second derivatives of any angle-bending potential  $U_{\alpha}[\theta]$  with respect to changes in the Cartesian coordinates of atoms A, B, C. Dubbeldam et al.<sup>S2</sup> also gave universal formulas for first and second derivatives of any angle-bending potential with respect to changes in the unit cell's size and shape for unrelaxed homogeneous strain. Their formulas require the following inputs:

$$f_1 = \frac{dU_{\measuredangle}[\theta]}{d\cos[\theta]}$$
(S47)

$$f_{2} = \frac{d^{2}U_{\measuredangle}[\theta]}{\left(d\cos[\theta]\right)^{2}}$$
(S48)

For conciseness, we first define

$$U_{new} \left[\theta\right] = k \frac{2\left(\cos\theta - \cos\theta_{eq}\right)^{2}}{\sin^{2}\theta + 3\sin^{2}\theta_{eq}} \left(\frac{\tanh\left[2\sin\left[\theta/2\right]\right]}{\tanh\left[2\sin\left[\theta_{eq}/2\right]\right]}\right) = k \frac{2\left(\cos\theta - \cos\theta_{eq}\right)^{2}}{w\left[\theta\right]}$$
(S49)

$$p[\theta] = \sqrt{2(1 - \cos\theta)} = 2\sin[\theta/2]$$
(S50)

$$tp[\theta] = tanh[p[\theta]]$$
(S51)

$$w[\theta] = 1 - \cos^2 \theta + 3(1 - \cos^2 \theta_{eq}) \left( \frac{tp[\theta]}{tp[\theta_{eq}]} \right)$$
(S52)

Its derivatives are

$$\mu[\theta] = -3 \left( \frac{1 - \cos^2 \theta_{eq}}{p[\theta]} \right) \left( \frac{1 - (tp[\theta])^2}{tp[\theta_{eq}]} \right)$$
(S53)

$$\frac{\mathrm{dw}[\theta]}{\mathrm{d\cos}[\theta]} = -2\cos\theta + \mu[\theta]$$
(S54)

$$\frac{d^2 w[\theta]}{\left(d\cos[\theta]\right)^2} = -2 + \frac{\mu[\theta]}{p[\theta]} \left(2tp[\theta] + \frac{1}{\left(p[\theta]\right)^2}\right)$$
(S55)

The derivatives of my new angle-bending function are:

$$\frac{dU_{new}\left[\theta\right]}{d\cos\left[\theta\right]} = \frac{1}{w\left[\theta, \theta_{eq}\right]} \left(4k\left(\cos\theta - \cos\theta_{eq}\right) - U_{new}\left[\theta\right]\frac{dw\left[\theta\right]}{d\cos\left[\theta\right]}\right)$$
(S56)

$$\frac{d^{2}U_{new}\left[\theta\right]}{\left(d\cos\left[\theta\right]\right)^{2}} = \frac{1}{w\left[\theta, \theta_{eq}\right]} \left(4k - 2\frac{dU_{new}\left[\theta\right]}{d\cos\left[\theta\right]}\frac{dw\left[\theta\right]}{d\cos\left[\theta\right]} - U_{new}\left[\theta\right]\frac{d^{2}w\left[\theta\right]}{\left(d\cos\left[\theta\right]\right)^{2}}\right)$$
(S57)

Eqn (S54)–(S57) were derived by substituting  $t = \cos[\theta]$  into eqn (S49) or (S52) and then differentiating with respect to t. Eqn (S56) and (S57) can be used to compute the potential's derivatives with respect to bond angle changes:

$$\frac{dU_{new}[\theta]}{d\theta} = -\sin[\theta]\frac{dU_{new}[\theta]}{d\cos[\theta]}$$
(S58)

$$\frac{d^{2}U_{new}\left[\theta\right]}{d\theta^{2}} = -\cos\left[\theta\right]\frac{dU_{new}\left[\theta\right]}{d\cos\left[\theta\right]} + \sin^{2}\left[\theta\right]\frac{d^{2}U_{new}\left[\theta\right]}{\left(d\cos\left[\theta\right]\right)^{2}}$$
(S59)

If  $\theta = \theta_{eq} = \pi$ , then the following formulas should be used to avoid division by zero:

$$\mathbf{U}_{\text{new}} \left[ \boldsymbol{\theta} = \boldsymbol{\theta}_{\text{eq}} = \boldsymbol{\pi} \right] = 0 \tag{S60}$$

$$\frac{dU_{new}}{d\cos[\theta]}\Big|_{\theta=\theta_{eq}=\pi} = k$$
 (S61)

$$\frac{\mathrm{d}^{2}\mathrm{U}_{\mathrm{new}}}{\left(\mathrm{d}\cos\left[\theta\right]\right)^{2}}\bigg|_{\theta=\theta_{\mathrm{ex}}=\pi}=0$$
(S62)

$$\left. \frac{\mathrm{d}U_{\mathrm{new}}}{\mathrm{d}\theta} \right|_{\theta = \theta_{\mathrm{eq}} = \pi} = -k\sin\left[\theta\right]$$
(S63)

$$\frac{d^2 U_{new}}{d\theta^2}\Big|_{\theta=\theta_{eq}=\pi} = -k\cos[\theta]$$
(S64)

#### References

- S1. W. C. Swope and D. M. Ferguson, Alternative expressions for energies and forces due to angle bending and torsional energy, *J. Comput. Chem.*, 1992, **13**, 585-594, DOI: 10.1002/jcc.540130508.
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