Electronic supplementary information (ESI)

August 9, 2024

Inverse temperature. The virtual inverse temperature T_{inv} (further only inverse temperature) is commonly introduced in Monte carlo simulations, with background in statistical physics. For a simple illustration, the probability of a state a with an energy E_a would be

$$P(E_a) = e^{\frac{-E_a}{k_B T}}$$

and when using β , called the thermodynamic beta $\beta = \frac{1}{k_B T}$:

 $P(E_a) = e^{-\beta E_a}$

This can be compared with a Gaussian distribution for a variable x with zero mean:

$$P(x) \sim e^{-\frac{x^2}{\sigma^2}}$$

where x is a distance from a mean value and σ^2 the standard deviation. The connection between the energy E and x for a harmonic model (harmonic oscillator), corresponding to the Gaussian distribution for the prior, is $E = \frac{1}{2}kx^2$. The inverse temperature $T_{inv} = \beta$ can be then borrowed from the statistical physics, for the Monte carlo simulation, resulting in

$$P(x) \sim e^{-\frac{T_{\text{inv}}x^2}{2\frac{1}{k}}} = e^{-\frac{T_{\text{inv}}x^2}{\sigma^2}},$$

where the harmonic constant $1/k = \sigma^2$ and the inverse temperature T_{inv} (here ≤ 1) broadens the sampled distribution as if at higher temperature.

Forces using different DFT functionals. The atomic force components were computed using the ω B97X DFT functional with D4 dispersion correction. One could argue though, that a double hybrid such as the DSD-PBEP86 functional with D3-BJ dispersion correction would give more accurate results. Figure 1 shows the correlation between the cartesian forces between these two DFT functionals. This validates the ω B97X DFT functional, without the need to using more expensive functionals.



Figure 1: DSD-PBEP86-D3-BJ vs $\omega B97X\text{-}D4$ force difference components

Tables 1 and 2 give the non-bonded and the bonded parameters respectively of the atomic resolution discussed in the main manuscript (Subsection 3.1).

Table 1: atomic resolution, non-bonding parameters						
vdW form	pcO^a	r^b_{\min}	ε^c	D		
$LJ_{vdw,12-6}$						
0	-0.7468	2.0598	0.0176			
Н		1.409	0.000001			
$LJ_{\rm vdw,10-6}$						
$LJ_{\rm vdw,8-6}$						
$LJ_{\rm vdwpr,12-6}$						
0-0	-0.7214	4.2582	0.0121			
O-H		3.1224	0.000289			
H-H		3.9662	0.000094			
$LJ_{\rm vdwpr,10-6}$						
0-0	-0.7456	3.9419	0.1020			
O-H		3.2433	0.001164			
H-H		4.1820	0.000275			
$LJ_{\rm vdwpr,8-6}$						
0-0	-0.7788	4.1869	0.3159			
O-H		5.2451	0.000171			
H-H		4.5799	0.000795			
$Buckingham^d_{ydwpr}$		А	В	С		
0-0	-0.6314	3.9057	0.038010	34487.5		
O-H		0.021760	84085.9	55650.5		
H-H		0.001190	489571.7	6645.4		
$Morse^{e}_{vdwpr}$		Re	А	D		
0-0	-0.6742	3.1061	0.5846	2.08488		
O-H		1.8662	0.2040	3.6677		
H-H		0.1356	15.5413	4.3203		

^a Partial charge at the oxygen site. The partial change of hydrogen is always $-0.5 \times pcO$

^b Distance (Å) at the energy minimum, an atom-specific parameter con-tribution in the case of vdw, or the full distance in the case of vdwpr. c Minimum energy (kcal/mol), an atom-specific parameter contribution in

the case of vdw, or the full energy in the case of vdwpr.

 d See Equation 4 in the main manuscript. ^{*e*} See Equation 5 in the main manuscript.

Table 2: Atomic resolution, bonding parameters completing two models of Table 1 $\overline{vdW \text{ form } r + \lfloor \tilde{\lambda} \rfloor} = \frac{|\lambda|}{|\lambda|} \frac{|\lambda|}{|\lambda|}$

vdW form	r_{\min} [A]	k[kcal/mol/A]
$LJ_{\rm vdw,12-6}$		
Bond	0.9627	485.04
Angle	105.84	30.05
$LJ_{\rm vdwpr,12-6}$		
Bond	0.9633	468.12
Angle	105.72	28.94

The correlation plot for the Cartesian components for the full forces in the atomic resolution is found in Figure 2.



Figure 2: Correlation plot between the Cartesian components of the full forces for QC and for $LJ_{vdw,12-6}$. The forces are in [kcal/mol/Å]. Atomic resolution.

The non-bonded and bonded optimized parameters for the molecular resolution (main manuscript Subsection 3.2) are found in Tables 3 and 4

vdW form	pcO	$r_{\rm min}$	ε	D
$LJ_{\rm vdw,12-6}$				
0	-0.8731	2.1167	0.018994	
Н		0.2978	0.000014	
$LJ_{\rm vdw,10-6}$				
0	-0.9092	2.0998	0.068610	
Н		1.3540	0.000003	
$LJ_{\rm vdw,8-6}$				
O	-0.9343	2.1717	0.2734	
Н		1.6801	0.000021	
$LJ_{\rm vdwpr,12-6}$				
0-0	-0.8377	3.6813	0.1017	
O-H	ATTRAC	0.056357	11.6304	
H-H		4.004571	0.000063	
$LJ_{\rm vdwpr,10-6}$				
0-0	-0.8482	3.5888	0.4075	
O-H	ATTRAC	0.2580	19.3712	
H-H		4.8174	0.000060	
$LJ_{\rm vdwpr,8-6}$				
0-0	-0.8298	3.8445	1.0910	
O-H	ATTRAC	0.0647	4.2633	
H-H		5.2328	0.000341	
Morse _{vdwpr}		Re	А	D
0-0	-0.7621	3.1187	0.7975	2.058559
O-H		1.8928	0.0848	4.5712
H-H		2.1499	0.0201	0.2903

 Table 3: Molecular resolution, non-bonding parameters

Table 4: Molecul	ar resolution, h	oonding parameters
vdW form	r_{\min} [Å]	k[kcal/mol/Å]
$LJ_{\rm vdw,12-6}$		
Bond	0.9609	697.46
Angle	106.83	44.34
$\overline{LJ_{\text{vdwpr},12-6}}$		
Bond	0.9613	635.57
Angle	106.69	39.41

The non-bonding optimized parameters for the intermediate resolution (main manuscript Subsection 3.3) are found in Table 5

Table 5: Intermediate resolution, non-bonding parameters for vdw and vdwpr 12-6.

vdW form	рсO	r_{\min}	ε
$LJ_{\rm vdw,12-6}$			
Ο	-0.7433	2.21736	0.007791
Η		1.1992	0.000001
$LJ_{\rm vdwpr,12-6}$			
O-O	-0.7161	4.7146	0.003801
O-H		2.8109	0.000707
H-H		5.7877	0.000001

Tables 6 and 7 show the non-bonded optimized parameters for different weights in the low-forcedifference region (main manuscript Subsection 3.4.1). Tables 8 and 9 show the bonded parameters for a weight of $\omega = 10$ and $\omega = 100$ respectively.

vdW form	pcO	r_{\min}	ε
$LJ_{\rm vdw,12-6,\omega=1}$			
0	-0.8048	2.1641	0.0098
Н		1.7688	0.000000
$LJ_{\rm vdw,12-6,\omega=10}$			
0	-0.9458	2.1714	0.009660
Н		1.6568	0.000000
$LJ_{\rm vdw,12-6,\omega=100}$			
0	-1.0101	2.0325	0.02221
Н		1.6157	0.000001
$LJ_{vdwpr,12-6,\omega=1}$			
0-0	-0.7723	4.2706	0.0118
O-H		4.4845	0.000004
H-H		4.2539	0.000046
$LJ_{\rm vdwpr,12-6,\omega=10}$			
0-0	-0.8900	4.4151	0.008243
O-H		4.7717	0.000003
H-H		3.9891	0.000136
$LJ_{\text{vdwpr},12-6,\omega=100}$			
0-0	-0.9421	4.2624	0.0129
O-H		4.5448	0.000006
H-H		3.8961	0.000203

Table 6: Parameters obtained at atomic resolution, with weight ω

vdW form	рсO	r_{\min}	ε
$LJ_{\mathrm{vdw},12-6,\omega=1}$			
0	-0.7951	2.3433	0.004121
Н		1.6223	0.000000
$LJ_{\mathrm{vdw},12-6,\omega=10}$			
0	-0.9204	2.3541	0.004316
Н		1.3655	0.000001
$LJ_{\rm vdw,12-6,\omega=100}$			
0	-0.9790	2.4007	0.003571
Н		1.7323	0.000000
$LJ_{vdwpr,12-6,\omega=1}$			
0-0	-0.7635	4.9181	0.002377
O-H		3.1531	0.000208
H-H		5.6973	0.000001
$LJ_{vdwpr,12-6,\omega=10}$			
0-0	-0.8695	4.7714	0.003748
O-H		3.0150	0.000473
H-H		5.8030	0.000001
$LJ_{\text{vdwpr},12-6,\omega=100}$			
0-0	-0.9179	4.7793	0.003860
O-H		3.2950	0.000175
H-H		5.9420	0.000001

Table 7: Parameters obtained at intermediate resolution, with weight \boldsymbol{w}

Table 8: Bonded parameters for a weight of $\omega = 10$. AR = atomic resolution, IR = intermediate resolution.

vdW form	r_{\min} [Å]	$\rm k[kcal/mol/Å]$
$LJ_{\rm vdw,12-6,AR}$		
Bond	0.9599	664.09
Angle	107.37	44.09
$LJ_{\rm vdwpr,12-6,AR}$		
Bond	0.9607	615.89
Angle	107.75	27.21
$LJ_{\rm vdw, 12-6, IR}$		
Bond	0.9504	679.58
Angle	107.15	45.09
$LJ_{\rm vdwpr,12-6,IR}$		
Bond	0.9611	634.58
Angle	106.95	41.68

vdW form	r_{\min} [Å]	k[kcal/mol/Å]
$LJ_{\rm vdw,12-6,AR}$		
Bond	0.9592	728.32
Angle	107.73	49.04
$LJ_{\rm vdwpr,12-6ARA}$		
Bond	0.9601	657.00
Angle	107.08	59.26
$LJ_{\rm vdw,12-6,IR}$		
Bond	0.9582	560.81
Angle	107.37	53.464
$LJ_{\rm vdwpr,12-6,IR}$		
Bond	0.9605	689.16
Angle	107.24	45.87

Table 9: Bonded parameters for a weight of $\omega = 100$. AR = atomic resolution, IR = intermediate resolution.

Table 10 shows the atomic resolution parameters after modifying the LJ repulsion power and the Coulomb interaction (main manuscript Subsection 3.4.2).

vdW form	pcO	r_{\min}	ε
$LJ_{vdwpr,12-6,cf}$			
0-0	-0.6300	4.3260	0.009666
O-H		3.7597	0.000016
H-H		4.0631	0.000076
cf	0.5774		
$LJ_{vdwpr,12-6,N-6,cf}$			
0-0	-0.6770	4.4301	0.0371
O-H		5.1601	0.000012
H-H		4.8425	0.000096
cf	0.6957		
LJ-N	9.65		

Table 10: Parameters with cf, LJ-N for the atomic resolution

In Figure 3 we show the correlation plot for the non-bonding cartesian components of several common simple water models with respect to our QC.



Figure 3: Non-bonding force components for several common water models compared with forces from the QC, obtained with ORCA.

In Figure 4 is shown the correlation distribution between the original AMOEBA force field for water and our QC forces.



Figure 4: Non-bonding force components for the original AMOEBA water model compared with forces from QC calculation.

Next is shown the original AMOEBA and AMOEBA-ASTA-0 parameters in Tinker format: Amoeba2018

vdw vdw	90 91			3.4050 2.6550) 0) 0	.1100 .0135	0.910	
bond angle	90 91	91 90	91	556.85 48.70	5 0 0 1	.9572 08.50		
ureybrad	91	90	91	-7.60) 1	.5537		
multipole	349	-350	-350		-0. 0. 0.	51966 00000 37928 00000	0.00000	0.14279
					0.	00000	0.00000	0.03881
multipole	350	349	350		0.: -0. -0.	25983 03859 03673	0.00000	-0.05818
					0.	00000	-0.10739	0 14412
polarize polarize	349 350		0 0	.8370 .4960	0.3900 0.3900	350 349		
AMOE	BA-	ASTA	- 0 ¹					
vdw 90 vdw 91	3.507 2.734	715 C 165 C	.20262 .00654	2826 4387 0.91	.0			
bond angle ureybrad	90 91 91	91 90 90	91 91	531.44 0 47.71 10 -11.9277).9632)5.33 1.5114			
multipole	349	-350	-350	-0.48971 0.0 0.0 0.57542 0.0 -0.3 0.0 0.0	0.046 86762 -0.207	72 80		

 $^{^{1}}$ The parameters are given for convenience of easy comparison with the original AMOEBA water parameters. They were not well tested with Tinker however. The development and testing has been done only with OpenMM.

```
multipole
            350
                  349
                       350 0.244855
                            0.02775
                                      0.0 -0.18587
                            0.03933
                            0.0 0.04819
                            0.01055 0.0 -0.08752
polarize
          349
                 0.1653
                            0.4799 350
polarize
          350
                 0.3763
                            0.8161 349
```



Figure 5: Distribution of χ^2 for different amounts of non-bonding parameters. For the case of the simple water models, our ASTA approach has 5 parameters, *i.e.*, 2 pairs of LJ for H and O and the oxygen charge. For OPC3 there are only 3, since the LJ for the H is not taken into account. For the case of AMOEBA water model we have 17 parameters. These distributions are obtained as skewed-Gaussian approximations of simulations mimicking the actual MCMC simulation of this study. We can see that for the case of the OPC3, the χ^2 is well sampled with a non-zero probability density to up to $\chi^2 \approx 7.0$. On the other hand, the obtained χ^2 for ASTA for the case of LJ_{12-6} was 2600, while for the OPC3 was 6500. This means that the parameters computed with ASTA are 6500 - 2600 = 3900 away from the correct subspace to sample the set of parameters providing good bulk properties. This implies, that according to Equation 6 from the main manuscript, we would need to sample the parameters for ASTA at high temperature: $T_{inv} = \frac{7}{3900} = 0.0018$. This also means, that the parameters for good bulk properties would not be compatible with the most accurate inter-molecular forces achievable by this force field form. For the case of AMOEBA, since we have more parameters, we can observe that it can be well sampled to up to $\chi^2 \approx 25.0$. Also, for the case of AMOEBA-ASTA-0, we got a χ^2 of 1234, while for the AMOEBA-ASTA-TEST it was of 1045, meaning that the AMOEBA-ASTA-TEST is a $\Delta \chi^2 = 189$ away from parameters providing good bulk properties. This means that this time we would only need a slightly increased temperature: $T_{\text{inv}} = \frac{25}{189} = 0.13$ in our MCMC, in order to reach the right subspace. Therefore, due to both, larger number of parameters shifting the χ^2 distribution to the right, and smaller $\Delta \chi^2$ at $T_{\rm inv} = 1$, the accurate inter-molecular forces and good bulk properties can be obtained by the same parameter set.