# Supplemental Information for: Low-temperature synthesis of cation-ordered bulk Zn<sub>3</sub>WN<sub>4</sub> semiconductor via heterovalent solidstate metathesis

Christopher L. Rom<sup>1,\*</sup>, Shaun O'Donnell<sup>1,2</sup>, Kayla Huang<sup>1,3</sup>, Ryan A. Klein<sup>1,4</sup>, Morgan J. Kramer<sup>4,5</sup>, Rebecca W. Smaha<sup>1</sup>, Andriy Zakutayev<sup>1,\*</sup>

<sup>1</sup> Materials, Chemical, and Computational Science, National Renewable Energy Laboratory, Golden, CO, 80401, USA

<sup>2</sup> Department of Chemistry, Colorado State University, Fort Collins, CO, 80523, USA

<sup>3</sup> University of Illinois Urbana-Champaign, Champaign, IL, 61801, USA

<sup>4</sup>Center for Neutron Research, National Institute of Standards and Technology, Gaithersburg, MD, 20899, USA

<sup>5</sup> Department of Chemistry, Southern Methodist University, Dallas, TX, 75275 USA

\* Corresponding authors: <u>christopher.rom@nrel.gov</u> and <u>andriy.zakutayev@nrel.gov</u>

### **Table of Contents**

Experimental	2
Synthesis of Zn <sub>3</sub> WN <sub>4</sub>	2
Synthesis of Zn <sub>3</sub> MoN <sub>4</sub>	2
In situ synchrotron powder X-ray diffraction analysis	3
<i>Ex situ</i> powder X-ray diffraction analysis of Zn <sub>3</sub> WN <sub>4</sub>	3
Compositional, thermodynamic, and property measurements	4
Additional PXRD measurements	5
In situ variable temperature SPXRD measurements	6
Synthesis of Zn <sub>3</sub> MoN <sub>4</sub>	9
DSC measurements	1
Evidence of decomposition in the $3 ZnCl_2 + Li_6WN_4$ reaction	2
Compositional characterization1	4
Structural models for the high resolution SPXRD measurements	5
Magnetic susceptibility measurements	20
Full structure visualizations of Li <sub>6</sub> WN <sub>4</sub> and Zn <sub>3</sub> WN <sub>4</sub>	!1
References	21

# Experimental

#### Synthesis of Zn<sub>3</sub>WN<sub>4</sub>

*Caution:* Metathesis reactions have the potential to rapidly release heat and gas.  $N_2$  gas formation within sealed ampules can cause explosions at high temperatures, and appropriate precautions should be taken to prevent equipment damage and unsafe conditions.

As some precursors are highly moisture sensitive, all precursors were prepared and stored in an argon-filled glovebox ( $O_2 < 0.1$  ppm,  $H_2O < 0.1$  ppm) unless explicitly mentioned. ZnF<sub>2</sub> ( $\geq$ 99 %, anhydrous, Alfa Aesar), ZnCl<sub>2</sub> ( $\geq$ 99.995%, anhydrous, Sigma Aldrich), ZnBr<sub>2</sub> (99.999%, anhydrous, Aldrich), Li<sub>3</sub>N ( $\geq$ 99.5%, 80 mesh, Sigma Aldrich), W (99.95%, <1 micron powder, Thermofisher Scientific) were used as received.

Li<sub>6</sub>WN<sub>4</sub> was synthesized in a method modified from that of Yuan et al.<sup>[1]</sup> Solid precursors (2.1 Li<sub>3</sub>N + W, ca. 5 mol% excess Li<sub>3</sub>N to account for loss by evaporation) were ground with a mortar and pestle and loaded into Zr crucibles with Zr lids (ca. 1 g loose powder). The Zr crucibles were then loaded into sacrificial quartz tubes (open on one end), which were loaded into quartz process tubes and heated in a tube furnace. Custom endcaps with quick disconnects enabled air-free transfer from the glovebox to the tube furnace (under Ar or N<sub>2</sub>). The samples were reacted under flowing N<sub>2</sub> (50 sccm, 99.999% purity) with a +5 °C/min ramp followed by a 12 h dwell at 850 °C and then natural cooling after turning off the furnace. Samples were confirmed to be phase pure by powder X-ray diffraction (Figure S1).

Syntheses for Zn<sub>3</sub>WN<sub>4</sub> were conducted by grinding together Li<sub>6</sub>WN<sub>4</sub> with Zn $X_2$  (X = Cl, Br) in a ratio of approximately 1:3. The powders were pelletized with 6 mm diameter dies in an arbor press (ca. 100 mg per pellet), loaded into quartz ampules (10 mm OD, 10 mm ID, ca 10 cm length), sealed under vacuum (<0.03 Torr), and heated in a muffle furnace. The optimized synthesis for Zn<sub>3</sub>WN<sub>4</sub> used a ratio of Li<sub>6</sub>WN<sub>4</sub> + 3.1 ZnBr<sub>2</sub> and was scaled up to a 3 g batch, sealed in a quartz ampule under vacuum (< 0.03 Torr) and heated at +5 °C/min to 300 °C for a 1 h dwell, then allowed to cool naturally. Samples were recovered into the glovebox. Reaction byproduct LiX was washed away using anhydrous and degassed methanol that was dried over molecular sieves. For washing, centrifuge tubes were loaded with approximately 500 mg of product powder and 1.5 mL methanol. The tube was agitated with a vortex, centrifuged, and the supernatant was decanted. This wash was repeated for a total of 3 cycles. Recovered powders were dried overnight under vacuum. However, Zn<sub>3</sub>WN<sub>4</sub> ultimately proved to be stable against air and water, and we note that the anhydrous washing may not be necessary.

### Synthesis of Zn<sub>3</sub>MoN<sub>4</sub>

Just like Li<sub>6</sub>WN<sub>4</sub>, Li<sub>6</sub>MoN<sub>4</sub> was synthesized using Li<sub>3</sub>N and Mo ( $\geq$ 99.9%, 1–5 micron powder, Sigma Aldrich), following a method modified from that of Yuan et al.<sup>[1]</sup> Heating the powders at 850 °C for 12 h (with slight Li<sub>3</sub>N excess) resulted in a phase pure Li<sub>6</sub>MoN<sub>4</sub> (Figure S1b). This Li<sub>6</sub>MoN<sub>4</sub> was then mixed with ZnBr<sub>2</sub>, pelletized, sealed in an ampule under vacuum, and heated

at +5 °C/min to 300 °C for a 1 h dwell, followed by natural cooling. The product was then washed with anhydrous methanol.

#### In situ synchrotron powder X-ray diffraction analysis

In situ synchrotron powder X-ray diffraction (SPXRD) measurements were conducted at beamline 17-BM-B of the Advanced Photon Source at Argonne National Laboratory. For these experiments ( $\lambda = 0.24101$  Å), the PerkinElmer plate detector was positioned 700 mm away from the sample. Homogenized precursors were packed into quartz capillaries in an Ar glovebox and flame-sealed under vacuum (<30 mTorr). Capillaries were loaded into a flow-cell apparatus<sup>[3]</sup> and heated at 5 °C/min to the specified temperature. A thermocouple was placed against the tip of the sample capillary, approximately 2 mm horizontally from the position of the X-ray beam. Diffraction pattern images were collected every 30 s by summing 20 exposures of 0.5 s each (10 s of summed exposure), followed by 20 s of deadtime. Images collected from the plate detector were radially integrated using GSAS-II and calibrated using a silicon standard.

Sequential Rietveld refinements were conducted on *in situ* SXPRD datasets using TOPAS Professional v6.<sup>[4]</sup> Lattice parameters, background terms, and scale factors were refined for each phase as a function of temperature, while atomic coordinates and occupancies were held constant at the initial values of the reference structure. A weighted scale factor (W.S.F.) Q was calculated for each phase p as a product of scale factor S, cell volume V, and cell mass M:  $Q_p = S_p \cdot V_p \cdot W_p$ .<sup>[5]</sup> We note that amorphous and liquid phases are inherently not observed in powder diffraction measurements and therefore cannot be accurately included in this analysis. A Lorentzian size broadening term was refined for each phase to model the peak shape using the pattern showing the greatest intensity of the relevant phase; this term was then fixed for the sequential refinements to better account for changes in intensity. To help stabilize the sequential refinement, isotropic displacement parameters ( $B_{iso}$ ) were fixed at 1 Å<sup>2</sup> for all atoms, but we note that this is likely not physical for a variable temperature investigation.

#### Ex situ powder X-ray diffraction analysis of Zn<sub>3</sub>WN<sub>4</sub>

The products of all reactions were characterized by powder X-ray diffraction (PXRD). Laboratory X-ray diffraction patterns were collected on a Rigaku Ultima IV diffractometer with Cu K $\alpha$  X-ray radiation at room temperature. All samples were initially prepared for PXRD measurements inside the glovebox; powder was placed on off-axis cut silicon single crystal wafers to reduce background scattering and then covered with polyimide tape to impede exposure to atmosphere. After Zn<sub>3</sub>WN<sub>4</sub> was determined to be moderately air stable, PXRD patterns were collected without polyimide tape to decrease the background signal.

High resolution synchrotron powder X-ray diffraction (SPXRD) measurements were conducted at beamline 28-ID-2 of the National Synchrotron Light Source II ( $\lambda = 0.1821$  Å) at Brookhaven National Laboratory. Samples were sealed under vacuum in quartz capillaries, which were then nested in Kapton capillaries. Data were collected for 60 seconds at T = 25 °C while spinning. Scattered photon intensity was measured using a Perkin-Elmer XRD 1621 Digital Imaging Detector. The data were reduced using Dioptas.<sup>[6]</sup> Pawley fits and subsequent Rietveld refinements were conducted using TOPAS Academic.<sup>[4]</sup>

Rietveld refinements were conducted for the laboratory PXRD and SPXRD patterns using TOPAS and TOPAS Academic, v6 (Bruker AXS).<sup>[4]</sup> Reference structures were sourced from the Inorganic Crystal Structure Database (ICSD). The *Pmn2*<sub>1</sub> Zn<sub>3</sub>MoN<sub>4</sub> structure (ICSD Col. Code 255744) was used as a starting model for *Pmn2*<sub>1</sub> Zn<sub>3</sub>WN<sub>4</sub>, with the Mo replaced by W.<sup>[7]</sup> For the cation disordered *P6*<sub>3</sub>*mc* Zn<sub>3</sub>WN<sub>4</sub> structure, *P6*<sub>3</sub>*mc* ZnO was used as a starting model, with atomic occupancies adjusted to match the stoichiometry of Zn<sub>3</sub>WN<sub>4</sub>, and lattice parameters adjusted to match the SPXRD pattern. For each structure, lattice parameters, isotropic displacement parameters, and general atomic coordinates were refined. For some models, Zn and W occupancy were refined as detailed in the Discussion section and the Supplemental (Figure S7). Structural visualizations and reference PXRD patterns were generated using VESTA.<sup>[8]</sup>

#### Compositional, thermodynamic, and property measurements

The composition of nominal Zn<sub>3</sub>WN<sub>4</sub> was measured by X-ray Fluorescence spectroscopy (XRF) and combustion analysis. Cation composition was quantified by XRF using a Bruker M4 Tornado with a Rh X-ray source. Samples were pelletized and XRF spectra were collected at 4 points across the pellet. Zn and W ratios were quantified from each spectra using the Bruker M4 software.

Differential scanning calorimetry (DSC) experiments were conducted using a Q20 system from TA Instruments. Samples were prepared in an Ar-filled glovebox. Samples (ca. 10 mg) were loaded into aluminum pans, which were crimped closed with an aluminum lid. The reference pan was also crimped closed under argon. Pans were then transferred out of the glovebox for measurement, and data were collected upon ramping up to 400 °C at a rate of 10 °C/min. Thermodynamic calculations for reaction enthalpies ( $\Delta H_{rxn}$ ) were conducted using formation enthalpy ( $\Delta H_f$ ) values reported in the Materials Project database.<sup>[9,10]</sup> Chemical potential diagrams were generated in the method of Yokokawa<sup>[11]</sup> as detailed by Todd et al.<sup>[12]</sup> and implemented in pymatgen (v2023.9.10).<sup>[13]</sup>

UV-vis measurements were conducted on a Cary 6000 UV-Vis-NIR spectrometer. PTFE was used as a white reflectance standard. Absorbance was calculated with the Kubelka-Munk transformation,  $k/s = (1 - R)^2 / 2R$  (where R is the reflectance, k is the apparent absorption coefficient, and s is the apparent scattering coefficient).

DC susceptibility data were measured on a Quantum Design Physical Properties Measurement System (PPMS) from T = 2 to 305 K in applied fields up to  $\mu_0 H = 14$  T.

# **Additional PXRD measurements**

Laboratory PXRD confirms the phase purity of the precursors that were synthesized for our metathesis reactions (Figure S1). All peaks index to the reference structures from the ICSD. The broad background at  $2\theta < 35^{\circ}$  comes from the polyimide tape used to protect the powder from air and moisture. This background suppresses the measured peak intensities for peaks at  $2\theta < 35^{\circ}$ , leading to the poor fit for that region of the pattern.



 $Li_6MoN_4$  and  $Li_6WN_4$ , respectively).

### In situ variable temperature SPXRD measurements

Figures S2 and S3 show the *in situ* variable temperature SPXRD heatmaps for the  $ZnCl_2$  and  $ZnF_2$  reactions, respectively. The analogous  $ZnBr_2$  reaction is shown in the main text (Figure 1). These data were used for sequential Rietveld analysis, which is presented in Figure 4 in the main text.

*In situ* SPXRD of the 3 ZnCl<sub>2</sub> + Li<sub>6</sub>WN<sub>4</sub> reaction (Figure S2) proceeds similarly to the ZnBr<sub>2</sub>based reaction shown in the main text (Figure 1). Bragg peaks arising from the precursors (3 ZnCl<sub>2</sub> + Li<sub>6</sub>WN<sub>4</sub>) stay steady up to 200 °C, where they begin to decrease in intensity. Simultaneously, Bragg peaks for Zn<sub>3</sub>WN<sub>4</sub>, LiCl, and Li<sub>2</sub>ZnCl<sub>4</sub> begin to increase in intensity. The set of Bragg peaks corresponding to the Li<sub>2</sub>ZnCl<sub>4</sub> phase fade out by 310 °C. We suspect that Li<sub>2</sub>ZnCl<sub>4</sub> is not an essential intermediate but rather the product of a transient side reaction between the precursor ZnCl<sub>2</sub> and product LiCl. The main factor influencing reactivity is the liquidus line in the LiCl-ZnCl<sub>2</sub> system, with the lowest eutectic being at 275 °C.<sup>[14]</sup> The intensity of the Bragg peaks corresponding to the LiCl phase reaches a maximum near ca. 340 °C, and then slowly decreases in intensity up to 550 °C (LiCl melting point is 605 °C).<sup>[15]</sup> Zn<sub>3</sub>WN<sub>4</sub>'s Bragg peaks remain approximately constant in intensity above 300 °C, persisting through the duration of the heating process. These processes are consistent with the following reactions:

 $Li_6WN_4 + 3 ZnCl_2 \rightarrow Zn_3WN_4 + 6 LiCl$ 2 LiCl + ZnCl<sub>2</sub>  $\rightarrow Li_2ZnCl_4$ 

Figure S3 shows that the 3  $ZnF_2 + Li_6WN_4$  reaction does not yield  $Zn_3WN_4$  until 565 °C, a much higher temperature than the  $ZnCl_2$  and  $ZnBr_2$  reactions. Initial reactivity begins near 310 °C, indicated by a decrease in intensities for the set of Bragg peaks corresponding to  $Li_6WN_4$  and  $ZnF_2$ . This initial reactivity is well below the melting point of  $ZnF_2$  (872 °C).<sup>[15]</sup> Concurrent with this initial reaction, an unknown phase briefly grows in (between 310 °C and 404 °C). Extrapolating from the ZnBr<sub>2</sub> and ZnCl<sub>2</sub> reactions, the phase is likely a Li-Zn-F intermediate, but ternary fluorides in this space are poorly characterized.  $Li_2ZnF_4$  has been reported, but the structure is not well described and the unit cell does not match the unknown phase. This intensity of the Bragg peaks arising from the intermediate phase decreases to zero by 404 °C, above which  $Li_6WN_4$  is the only crystalline phase up to 565 °C. At this point,  $Li_6WN_4$  fades out and  $Zn_3WN_4$ , LiF, and a rocksalt phase fit with WN grow in. The presence of this rocksalt phase indicates that the higher reaction temperature and greater exothermicity of the ZnF<sub>2</sub> reaction (compared with the Cl and Br versions) leads to a substantial degree of decomposition of the Zn<sub>3</sub>WN<sub>4</sub> phase. Therefore, this reaction was not explored further.



Figure S2. Heatmap of *in situ* SPXRD data upon heating 3  $ZnCl_2 + Li_6WN_4$  at +10 °C/min. Reference patterns for the reactants and products/intermediates are simulated at the bottom and top, respectively (ICSD Col. Codes 2459 for ZnCl<sub>2</sub>, 66096 for Li<sub>6</sub>WN<sub>4</sub>, 402399 for Li<sub>2</sub>ZnCl<sub>4</sub>, and 27981 for LiCl).<sup>[1]</sup>



Figure S3. Heatmap of *in situ* SPXRD data upon heating  $3 \text{ ZnF}_2 + \text{Li}_6\text{WN}_4$  at  $\pm 10 \text{ °C/min}$ . Reference patterns for the reactants and products/intermediates are simulated at the bottom and top, respectively. (ICSD Col. Codes 9169 for ZnF<sub>2</sub>, 66096 for Li<sub>6</sub>WN<sub>4</sub>, 18012 for LiF).<sup>[1]</sup>

# Synthesis of Zn<sub>3</sub>MoN<sub>4</sub>

The synthesis strategy used for Zn<sub>3</sub>WN<sub>4</sub> was also applied to synthesize Zn<sub>3</sub>MoN<sub>4</sub>, but the product exhibited partial decomposition. Analysis of the SPXRD pattern collected for the Mo analog of Li<sub>6</sub>WN<sub>4</sub>, synthesized via the reaction Li<sub>6</sub>MoN<sub>4</sub> + 3 ZnBr<sub>2</sub>  $\rightarrow$  Zn<sub>3</sub>MoN<sub>4</sub> + 6 LiBr, suggests phase decomposition (Figure S4). In addition to the desired Zn<sub>3</sub>MoN<sub>4</sub> (56 mol%), Rietveld analysis of High resolution SPXRD data show that a rocksalt (RS) structure fit as  $ZnMoN_{0.5}$  also forms (23) mol%), along with a Zn impurity (20 mol%). This RS phase exhibits a substantially larger lattice parameter (a = 4.7106(3) Å) than the defect-RS phase Mo<sub>2</sub>N (a = 4.16 Å to 4.19 Å).<sup>[17,18]</sup> Therefore, we hypothesize it may be a  $(Zn,Mo)N_x$  material, as octahedra  $Zn^{2+}$  has a substantially larger ionic radius (0.74 Å) than octahedral  $Mo^{x+}$  (0.65 Å for  $Mo^{4+}$ , 0.69 Å for  $Mo^{3+}$ ).<sup>[19]</sup> The rocksalt ZnMoN<sub>0.5</sub> phase was created from a  $Fm\overline{3}m$  Mo<sub>2</sub>N starting model. Rietveld analysis with the composition of  $ZnMoN_{0.5}$  provides a reasonable fit. Further analysis of this material is beyond the scope of this manuscript. Additional minor peaks that we have not indexed are present (possibly higher order oxides). These impurity phases suggest that  $Zn_3MoN_4$  is less stable at elevated temperatures than Zn<sub>3</sub>WN<sub>4</sub>. This decomposition occurs despite the excess ZnBr<sub>2</sub> which was intended to serve as a heat sink during the exothermic reaction. Despite the partial decomposition of the phase, Zn<sub>3</sub>MoN<sub>4</sub> is still the major phase in the pattern. As with Zn<sub>3</sub>WN<sub>4</sub>, the SPXRD pattern for Zn<sub>3</sub>MoN<sub>4</sub> shows evidence of cation-ordering in the form of the  $Pmn2_1$  reflections at low angle: e.g., (010), (110), (101), (011). However, these reflections are weaker than in the W case, owing to the lower scattering factor of Mo compared to W. We focused our work on Zn<sub>3</sub>WN<sub>4</sub> because W scatters Xrays more strongly than Mo (facilitating characterization) and because our Zn<sub>3</sub>WN<sub>4</sub> products exhibited higher phase purity.



### **DSC** measurements

The rapid, exothermic nature of the ZnCl<sub>2</sub>-based reaction produced an odd curvature (peak b-ii) in the temperature vs. heat flow trace shown in Figure 3b. The reaction releases so much heat in such a short time that the aluminum DSC pan and stage increase in temperature by approximately 16 °C before cooling back down to the programed temperature (Figure S5b). This peak in the time vs. temperature trace appears as the curved peak (peak b-ii) in the temperature vs. heat flow trace (Figure 3b). In contrast, the ZnBr<sub>2</sub> reaction does not noticeably affect the temperature of the DSC system (Figure S5a).



# Evidence of decomposition in the 3 ZnCl<sub>2</sub> + Li<sub>6</sub>WN<sub>4</sub> reaction

Figure S6 shows that a trace Zn impurity can be detected in reactions between  $Li_6WN_4 + 3.1 ZnCl_2$ , even when heated at low temperatures (ca. 250 °C). This impurity makes the powder appear grey in color. The Zn is likely produced via decomposition of Zn<sub>3</sub>WN<sub>4</sub> during the highly exothermic reaction (See DSC measurements, Figure 3). Surprisingly, we do not observe W or WN. This may indicate that tungsten remains in the Zn<sub>3</sub>WN<sub>4</sub> phase (which would then be Zn-poor), or that the W (or WN) is amorphous.



# **Chemical potential diagrams**

Chemical potential diagrams show the relative stability of  $Zn_3WN_4$  in the Li-Zn-W-N phase space. The Gibbs free energy is calculated as a function of temperature,<sup>[20]</sup> revealing that  $Zn_3WN_4$  is calculated to be thermodynamically stable (i.e., on the convex hull) at the synthesis temperature of 300 °C (Figure S7a,b), but becomes destabilized relative to competing phases (W,  $W_2N_3$ ,  $N_2$ , Zn) by 700 °C (Figure S7c,d). Therefore,  $Zn_3WN_4$  would likely decompose near 700 °C at 1 atm of  $N_2$ . To represent a 4-dimensional space in 3-dimensions, we project down the Li-W-N space onto its intersection with the Zn-W-N axes, so the stability region of Li<sub>6</sub>WN<sub>4</sub> is shown as a grey volume. The Li<sub>6</sub>WN<sub>4</sub> chemical potential space intersects that of Zn<sub>3</sub>WN<sub>4</sub> (Figure S6a,b), indicating that these two phases are thermodynamically stable against one another.<sup>[11]</sup> Prior work on oxides<sup>[12]</sup> and nitrides<sup>[21]</sup> has shown that this intersection in chemical potential space means interdiffusion can effectively occur between the two phases, allowing for the formation of the targeted phase (Zn<sub>3</sub>WN<sub>4</sub> in this case) without nucleating an intermediate at the interface.



Figure S7. Calculated chemical potential diagrams for the (Li)-Zn-W-N system at a,b) 300 °C and c,d) 700 °C. Panels a and c show a 3D perspective while panels b and d show the space on the W-N chemical potential axes.

### **Compositional characterization**

X-ray Fluorescence (XRF) spectroscopy identified a Zn:W ratio of 3.8:1, in excess of the expected 3:1 ratio for Zn<sub>3</sub>WN<sub>4</sub>. XRF was conducted on pelletized powder after washing away the bromide byproduct. The excess zinc may be incorporated in the wurtzite-derived lattice, as suggested by Rietveld analysis (discussed below). A representative raw spectrum is shown in Figure S8. Spectra were collected at 4 different points across the pellet, and each spectrum was fit using the Bruker software to quantify Zn and W atomic ratios. The Zn:W ratio of 3.8:1 was calculated by averaging the Zn:W values from the 4 spectra.





### Structural models for the high resolution SPXRD measurements

Table S1. Refined atomic coordinates for Zn <sub>3</sub> WN <sub>4</sub> in space group <i>Pmn</i> 2 <sub>1</sub> from the SPXRD										
shown in Figure 5a. The unit cell parameters refined to $a = 6.5602(8)$ Å, $b = 5.6813(7)$ Å, and										
$c = 5.3235(2)$ Å. $R_{exp} = 1.182$ %, $\bar{R}_{wp} = 3.989$ %, $R_p = 2.802$ %, GOF = 3.374. This structure is										
visualized in Figure 5d.										
Site	Wyckoff	Atom	x	y	Z	Occupancy	$B_{iso}$ (Å <sup>2</sup> )			
Zn1	4b	Zn	0.251(1)	0.834(1)	0.697(1)	1.00(1)	0.10(4)			
Zn1	4b	W	0.251(1)	0.834(1)	0.697(1)	0.00(1)	0.10(4)			
Zn2	2a	Zn	0	0.337(6)	0.698(1)	1.00(1)	0.20(7)			
Zn2	2a	W	0	0.337(6)	0.698(1)	0.00(1)	0.20(7)			
W1	2a	W	0	0.670(2)	0.221(1)	0.83(2)	0.44(4)			
W1	2a	Zn	0	0.670(2)	0.221(1)	0.17(2)	0.44(4)			
N1	4b	Ν	0.765(2)	0.846(3)	0.095(1)	1	0.30(18)			
N2	2a	Ν	0	0.330(4)	0.116(2)	1	0.4(3)			
N3	2a	Ν	0	0.678(4)	0.555(2)	1	0.4(3)			

The refined lattice parameters for Zn<sub>3</sub>WN<sub>4</sub> are shown in Tables S1.

We considered several structural models of Zn<sub>3</sub>WN<sub>4</sub> when conducting Rietveld analysis against our high resolution SPXRD patterns (Table S2). Several terms were allowed to vary for each approach: sample displacement, lattice parameters, size broadening (Lorentzian), strain broadening (Lorentzian), isotropic displacement parameters, and a 15-term background Chebyshev polynomial. Our most robust model was a single-phase model that allowed for a small degree of cation disorder for *Pmn*2<sub>1</sub> Zn<sub>3</sub>WN<sub>4</sub>, but with each cation site fixed to full occupancy (e.g., the Zn1 site was refined with Zn occupancy set to 1-*x* and W occupancy *x*). This model resulted in an  $R_{wp}$  of 3.989 % and is shown in Figure 5, Table S1, and Figure S9a. For comparison, a simpler model of *Pmn*2<sub>1</sub> Zn<sub>3</sub>WN<sub>4</sub> with fixed cation occupancies (e.g., the Zn1 site fixed with 1.0 Zn occupancy) gave a significantly worse fit to the data ( $R_{wp} = 4.638$  %, Figure S9b). Atomic positions were allowed to refine for both these single phase models. However, our diffuse reflectance spectroscopy measurements suggest that the material is not a perfectly cation-ordered phase. While the small of cation disorder in the *Pmn*2<sub>1</sub> Zn<sub>3</sub>WN<sub>4</sub> structure can explain the optical data,<sup>[22]</sup> we also considered a two-phase model.

Given the two distinct absorption onsets shown in the diffuse reflectance spectrum (Figure 5), we also considered a two-phase model in our Rietveld refinements. For the first two-phase model, we started with the  $Pmn2_1$  Zn<sub>3</sub>WN<sub>4</sub> from the fixed cation occupancy models. We then fixed atomic positions. Next, we created a model for cation disordered Zn<sub>3</sub>WN<sub>4</sub> in a  $P6_3mc$  structure (i.e., the wurtzite structure type), and fixed the lattice parameters of the disordered phase (dis) as functions of the refined lattice parameters for the ordered phase (ord):  $a_{dis} = 0.5a_{ord}$  and  $c_{dis} = c_{ord}$ . This relationship holds true for the order-disorder transition of the related wurtzite-derived material, ZnGeN<sub>2</sub>.<sup>[23]</sup> This disordered structure is consistent with the cation-disordered Zn<sub>3</sub>WN<sub>4</sub> synthesized via thin film sputtering.<sup>[24]</sup> We then refined the size and strain broadening for both phases. The disordered phase exhibited broader peaks (meaning smaller size, larger strain) than the ordered phase, but this analysis is qualitative rather than quantitative as the broadening contribution from the beam profile could not be reliably refined. This refinement resulted in 78 mol% phase fraction

of  $Pmn2_1$ , 22 mol% for  $P6_3mc$ , and an  $R_{wp}$  value of 3.953 % (Figure S9a), comparable to the singlephase model. However, we chose not to present this model in the main text, given that the extra complexity does not substantially improve the goodness of fit.

The best fit was obtained via a two-phase model, but with non-physical lattice parameters. Allowing the  $P6_3mc$  lattice parameters to freely refine results in the best fit we obtained by the Rietveld method ( $R_{wp} = 3.742$  %). However, the model is likely non-physical. In this model, the refined *c* lattice parameter for this  $P6_3mc$  phase increases substantially (5.451(2) Å) compared to the  $Pmn2_1$  phase (c = 5.3228(2) Å), which is not consistent with prior studies of order-disorder transitions in wurtzite-derived structures (e.g., ZnGeN<sub>2</sub>).<sup>[23]</sup> The *c* lattice parameters of ZnGeN<sub>2</sub> are identical in the  $Pna2_1$  (the ordered structure) and  $P6_3mc$  (the disordered structure), because cation disorder does not affect the layer spacing of the hcp anions along the (00*l*) direction. This analysis reveals ambiguities in these two-phase models.

Given the limitations of the two-phase models, we posit that the single phase model provides the most reliable fit to the SPXRD data without over-fitting the pattern.<sup>[25]</sup> Yet, the SPXRD measurements probe the long range average ordered structure. Local ordering—which we do not probe here—may influence the optical absorption properties shown in Figure 7. The impact of local ordering on optical properties has been characterized in the halide perovskite CsSnBr<sub>3</sub>,<sup>[26]</sup> in Fe doped SrTiO<sub>3</sub>,<sup>[27]</sup> and in carbon coated FeF<sub>3</sub>.<sup>[28]</sup>

<b>Rietveld approach</b>	Space group	Composition	$R_{wp}$ (%)	Figure	
Single phase.	Diana?	7no - WasaN	2 080	1a and S6a	
Refined cation occ.	rmn21	Z113.17 VV 0.831N4	5.989		
Single phase.	$D_{mn}$	ZnoWNL	1 628	Séh	
Fixed cation occ.	rmn21	Z113 W 1N4	4.038	300	
Two phases.	$Pmn2_1$ Zn <sub>3</sub> WN <sub>4</sub>				
Fixed cation occ.	Dí		3.953	S7a	
Fixed $a$ and $c$ for $P6_3mc$	$P6_{3}mc$ $Zn_{0.75}W_{0.25}N$				
Two phases.	$Pmn2_1$	Zn <sub>3</sub> WN <sub>4</sub>			
Fixed cation occ.	DC		3.742	S7b	
Refined <i>a</i> and <i>c</i> for $P6_3mc$	P03MC	$Ln_{0.75}$ w $_{0.25}$ N			

Table S2. Summary of models considered for the high resolution SPXRD data for Zn<sub>3</sub>WN<sub>4</sub>.







Figure S11. High resolution SXPRD 2D diffraction images from the 2D detector for the integrated patterns shown in Figure 1 (a)  $Zn_3WN_4$  and (b)  $Zn_3MoN_4$  + decomposition products rocksalt (RS)  $ZnMoN_{0.5}$  and metallic Zn.

Qualitative inspection of the 2D diffraction images show that the diffraction rings for the Zn<sub>3</sub>WN<sub>4</sub> powder are homogeneous (Figure S8). If rings with two different morphologies were present, this would suggest the presence of two distinct phases with different crystallinity, size, and strain. That we see only one morphology of ring in the 2D detector image supports either a single phase or multiple phases with nearly identical crystallinity, size, and strain. In our Rietveld refinement that modeled the data using two phases, the fit to the data is significantly worse when we constrain the size and strain broadening terms to be the same for both phases. These findings support our use of the single-phase model.

### Magnetic susceptibility measurements

Magnetic measurements were performed on  $Zn_3WN_4$  using a Quantum Design Physical Property Measurement System (PPMS). A powder sample of  $Zn_3WN_4$  was loaded into a small packet (0.001125 mg) and secured inside a plastic straw for the measurement. Magnetic susceptibility ( $\chi$ ) of  $Zn_3WN_4$  was measured as a function of temperature shows largely diamagnetic behavior with a trace paramagnetic impurity (Figure S12a). Similarly, magnetization (M) as a function of applied field (H) at 2 K shows that diamagnetism dominates the field-dependent magnetization (Figure S12b). These findings are inconsistent with pure  $Zn_3WN_4$ , which should be purely diamagnetic. Zn impurities, if present, would also give a diamagnetic response. The paramagnetic component suggests the possibility of a reduced tungsten species (e.g.,  $W^{5+}$ ), possibly as a sub-nitride (e.g., Zn\_3WN<sub>4- $\delta$ </sub>), an oxynitride impurity (e.g., Zn\_3WN<sub>4-x</sub>O<sub>x</sub>), or a W-rich phase (e.g., Zn<sub>3- $\delta$ </sub>W<sub>1+ $\delta$ </sub>N<sub>4</sub>).



Figure S12. A) Magnetic susceptibility ( $\chi$ ) of Zn<sub>3</sub>WN<sub>4</sub> as a function of temperature with an applied field ( $\mu_0 H$ ) of 1 T. b) Magnetization (M) as a function of  $\mu_0 H$  at 2 K. The grey outline in (b) is shown in more detail via the inset.

### Full structure visualizations of Li<sub>6</sub>WN<sub>4</sub> and Zn<sub>3</sub>WN<sub>4</sub>



Figure S13. Reaction scheme showing the structure of  $Li_6WN_4$  and  $Zn_3WN_4$ , with Li and Zn atoms included to complement the main text visualization (Figure 7).

#### References

- [1] W. X. Yuan, J. W. Hu, Y. T. Song, W. J. Wang, Y. P. Xu, Powder Diffr. 2005, 20, 18-21.
- [2] D. B. G. Williams, M. Lawton, J. Org. Chem. 2010, 75, 8351-8354.
- [3] P. J. Chupas, K. W. Chapman, C. Kurtz, J. C. Hanson, P. L. Lee, C. P. Grey, *J. Appl. Crystallogr.* **2008**, *41*, 822–824.
- [4] A. A. Coelho, J. Appl. Crystallogr. 2018, 51, 210–218.
- [5] P. K. Todd, A. M. M. Smith, J. R. Neilson, Inorg. Chem. 2019, 58, 15166–15174.
- [6] C. Prescher, V. B. Prakapenka, High Press. Res. 2015, 35, 223–230.
- [7] E. Arca, S. Lany, J. D. Perkins, C. Bartel, J. Mangum, W. Sun, A. Holder, G. Ceder, B. Gorman, G. Teeter, W. Tumas, A. Zakutayev, J. Am. Chem. Soc. 2018, 140, 4293–4301.
- [8] K. Momma, F. Izumi, J. Appl. Crystallogr. 2011, 44, 1272–1276.
- [9] A. Jain, G. Hautier, S. P. Ong, C. J. Moore, C. C. Fischer, K. A. Persson, G. Ceder, *Phys. Rev. B* 2011, 84, 045115.
- [10] A. Jain, S. P. Ong, G. Hautier, W. Chen, W. D. Richards, S. Dacek, S. Cholia, D. Gunter, D. Skinner, G. Ceder, K. A. Persson, *APL Mater.* 2013, 1, 011002.
- [11] H. Yokokawa, J. Phase Equilibria 1999, 20, 258–287.
- [12] P. K. Todd, M. J. McDermott, C. L. Rom, A. A. Corrao, J. J. Denney, S. S. Dwaraknath, P. G. Khalifah, K. A. Persson, J. R. Neilson, *J. Am. Chem. Soc.* 2021, 143, 15185–15194.
- [13] S. P. Ong, W. D. Richards, A. Jain, G. Hautier, M. Kocher, S. Cholia, D. Gunter, V. L. Chevrier, K. A. Persson, G. Ceder, *Comput. Mater. Sci.* 2013, 68, 314–319.
- [14] S. J. Shaw, G. S. Perry, Thermochim. Acta 1989, 155, 87–96.
- [15] J. R. Rumble (ed.), "Physical Constants of Inorganic Compounds" in CRC Handbook of Chemistry and Physics, CRC Press/Taylor & Francis, Boca Raton, FL., 2023
- [16] C. Martin, A. Durif-Varambon, J.-C. Joubert, Bull. Minéralogie 1965, 88, 141-141.
- [17] C. L. Bull, T. Kawashima, P. F. McMillan, D. Machon, O. Shebanova, D. Daisenberger, E. Soignard, E. Takayama-Muromachi, L. C. Chapon, J. Solid State Chem. 2006, 179, 1762–1767.
- [18] B. Cao, J. C. Neuefeind, R. R. Adzic, P. G. Khalifah, Inorg. Chem. 2015, 54, 2128–2136.
- [19] R. D. Shannon, Acta Crystallogr. Sect. A 1976, 32, 751–767.

- [20] C. J. Bartel, S. L. Millican, A. M. Deml, J. R. Rumptz, W. Tumas, A. W. Weimer, S. Lany, V. Stevanović, C. B. Musgrave, A. M. Holder, *Nat. Commun.* 2018, 9, 4168.
- [21] C. L. Rom, A. Novick, M. J. McDermott, A. A. Yakovenko, J. R. Gallawa, G. T. Tran, D. C. Asebiah, E. N. Storck, B. C. McBride, R. C. Miller, A. L. Prieto, K. A. Persson, E. Toberer, V. Stevanović, A. Zakutayev, J. R. Neilson, *J. Am. Chem. Soc.* 2024, 146, 4001–4012.
- [22] J. J. Cordell, G. J. Tucker, A. Tamboli, S. Lany, APL Mater. 2022, 10, 011112.
- [23] E. W. Blanton, K. He, J. Shan, K. Kash, J. Cryst. Growth 2017, 461, 38-45.
- [24] W. Sun, C. J. Bartel, E. Arca, S. R. Bauers, B. Matthews, B. Orvañanos, B.-R. Chen, M. F. Toney, L. T. Schelhas, W. Tumas, J. Tate, A. Zakutayev, S. Lany, A. M. Holder, G. Ceder, *Nat. Mater.* 2019, 18, 732–739.
- [25] B. H. Toby, Powder Diffr. 2006, 21, 67–70.
- [26] D. H. Fabini, G. Laurita, J. S. Bechtel, C. C. Stoumpos, H. A. Evans, A. G. Kontos, Y. S. Raptis, P. Falaras, A. Van der Ven, M. G. Kanatzidis, R. Seshadri, *J. Am. Chem. Soc.* 2016, 138, 11820–11832.
- [27] H. A. J. L. Mourão, O. F. Lopes, W. Avansi, M. J. M. Pires, S. Souza, C. Ribeiro, V. R. Mastelaro, *Mater. Sci. Semicond. Process.* 2017, 68, 140–146.
- [28] N. Yamakawa, M. Jiang, B. Key, C. P. Grey, J. Am. Chem. Soc. 2009, 131, 10525–10536.