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o-B₂N₂: A Promising Metal-free Photocatalyst for Highly Efficient Conversion of CO₂ to Hydrocarbons

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Figure S1. Band structure (left) and phonon dispersion curves (right) of the unit cell of orthorhombic boron nitride along the high symmetry points Γ -X-S-Y- Γ . The band gap at the Y point is 0.78 eV. Positive phonon frequencies suggests dynamic stability of the structure.



Figure S2. Six possible adsorption sites on the surface of the $o-B_2N_2$ monolayer. Boron and nitrogen atoms are shown in blue and white, respectively.



Figure S4 Optimized geometries of CO_2 molecule adsorbed over pristine orthorhombic boron nitride (o- B_2N_2) when CO_2 molecule interacts vertically over the six available active sites. (E_{ads}) represent the corresponding adsorption energies. The distance between the monolayer is indicated with arrow.

No of site	Active site	Adsorption Energy	Height	Height	Angle	C-O=1.179	
		(eV)	(Å) I	(Å) F	(degree)	(Å)	
Site 1	At B	-0.412865	2.115	3.113	179.8	1.186	
Table S1. Initial (I) and final (F) he	ight, change in the angle,	and bond le	ngth of CO ₂	for the desorpti	ion of CO ₂ over	
pristisite B2N2 monolayArt N		-0.425283	2.011	3.126	180	1.185	
Site 3	At B-B	-0.402461	2.191	3.245	179.6	1.185	
Site 4	At B-N	-0.430514	2.152	3.173	1.185	1.185	
Site 5	At N-N	-0.404501	2.177	3.291	179.8	1.186	
Site 6	At center	-0.413672	2.491	3.424	179.7	1.185	
		Vertically act	tive sites				
Site 7	At B	-0.347096	2.036	2.841	179.9	1.185	
Site 8	At N	0.339909	2.032	2.916	179.9	1.185	
Site 9	At B-B	-0.349769	2.191	2.975	180	1.184	
Site 10	At B-N	-0.349809	2.193	2.913	179.9	1.182	
Site 11	At N-N	-0.34554	2.139	2.95	180	1.186	
Site 12	at center	-0.370259	2.516	3.104	180	1.182	

Table S2. Calculated structural, adsorption, and electronic parameters of B-, N-, and BN-vacancy defected $o-B_2N_2$ monolayer before and after CO_2 adsorption. Here H and V denotes horizontal and vertical adsorptions, respectively.

A defective monolayer of O- B ₂ N ₂		Bond length(^Å)			CO ₂ adsorbed at a Distance	$E_g(eV)$	Adsorptio n Energy of CO ₂ (eV)
		B-B	N-N	B-N	(A)		
Pristine O-BN		1.72	1.45	1.44	_	0.78	-
B vacancy		1.69	1.41	1.38	-	Metallic	-
N vacancy		1.74	1.45	1.42	-	Metallic	-
BN vacancy		1.86	1.60	1.50	-	1.33	-
CO ₂ adsorbed over Pristine		1.72	1.45	1.44	3.425	0.78	-0.41
CO ₂ adsorbed over B	Η	1.745	1.414	1.45	1.354	0.0008	-3.7516
vacancy	V	1.698	1.414	1.404	2.656	0.0236	-0.3540
CO ₂ adsorbed	Н	1.728	1.447	1.441	3.012	0.0064	-0.3815
over it vacuncy	V	1.728	1.450	1.442	2.562	0.0073	-0.3835
CO ₂ adsorbed over BN vacancy	Н	1.868	1.603	1.445	2.989	1.325	-0.4358
	V	1.868	1.609	1.445	2.585	1.337	-0.3748



Figure S5. Band structures of SV-B vacancy defected $o-B_2N_2$ monolayers for two different vacancy concentrations. The Fermi level is set at 0 eV.

Computation Details for defects for Figure S5:

For our *ab initio* calculations we used the density functional theory (DFT) at the level of the generalized gradient approximation (GGA) and the Perdew-Burke- Ernzerhof (PBE)\ as as implemented in the Vienna *Ab-initio* Simulation Package (VASP) [1-4]. The projected augmented wave (PAW) potential is used to describe the core electrons. A kinetic energy cutoff of 520 eV was found to be sufficient to achieve a total energy convergence of the energies of the systems to within 1 meV. Gaussian smearing of 0.05 eV was chosen to accelerate the electronic convergence. The optimization of atomic positions (including full cell optimization) was allowed to proceed without any symmetry constraints until the force on each atom is less than 5 meV/A. In order to assess the effect of the vacancy concentration on the electronic properties we considered two vacancy defected o-BN structures: (i) one shown in Fig 1(a) and, (ii) a much larger 6x6x1 supercell, each containing a single B vacancy giving us defect concentrations of 2.8% and 0.7%, respectively. The k-point sampling was set to be $14 \times 14 \times 1$ for the former, while Gamma point was used for the latter. The supercells are taken to be periodic in the XY-plane with planes separated by 10 Å along the z-direction to avoid interactions between the planes. The band structures obtained using the VASP code is shown

in Figure S5 for the two structures. In the smaller unit cell, the defect-defect interaction causes extended states to develop in the band gap producing metallic behavior, while in the larger supercell the defects are sufficiently far apart to interact with each other and as a result, more bulk like (semiconducting) behavior is seen.

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Figure S6. Optimized geometric structures of the adsorption of CO_2 molecule and formation of various intermediates on the surface of 2D monolayer SV-B vacancy-defected o- B_2N_2 material with the equilibrium structures of COOH path.



Figure S7. Optimized geometric structures of the adsorption of CO_2 molecules and formation of various intermediates on the surface of 2D monolayer SV-B vacancy-defected o- B_2N_2 material with the equilibrium structures of OCOH path



Figure S8. Electron charge density plots of SV-B vacancy defected o- B_2N_2 material and all intermediates formed along the COOH path during the CO₂RR process.



Figure S9. Electron charge density plots of SV-B vacancy defected $o-B_2N_2$ monolayer and all intermediates formed along the OCOH path during the CO_2RR process.



Density of States (states/eV)

Figure S10. Shows total and projected DOS of horizontal CO_2 adsorption over B-vacancy defected o- B_2N_2 monolayer.