

1            **ESIPT-Suppressed 2-(2'-hydroxyphenyl)benzoxazole derivative as a new**  
2            **photoinitiator for multiphoton polymerization**

3 João V. P. Valverde<sup>a</sup>, André L. S. Romero<sup>a</sup>, Renan Cunha<sup>a</sup>, Rafael Q. Garcia<sup>a</sup>, Timothée  
4 Stoerkler<sup>b</sup>, Julien Massue<sup>b</sup>, Leonardo De Boni<sup>a</sup> and Cleber R. Mendonça<sup>a,\*</sup>

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6 <sup>a</sup> *São Carlos Institute of Physics, University of São Paulo, CP 369, 13560-970 São Carlos, SP, Brazil*

7 <sup>b</sup> *Institut de Chimie et Procédés pour l'Energie, l'Environnement et la Santé (ICPEES), Université de*  
8 *Strasbourg, UMR CNRS 7515, 67087 Strasbourg Cedex 02, France*

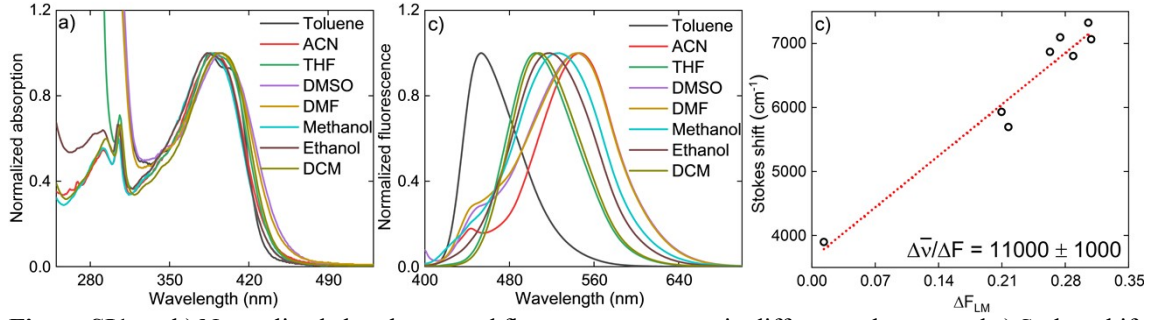
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10 \*Author to whom correspondence should be addressed: [crmendon@ifsc.usp.br](mailto:crmendon@ifsc.usp.br) and  
11 [joaovalverde@ifsc.usp.br](mailto:joaovalverde@ifsc.usp.br)

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13            **Supplementary Information**

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15 **SI1. Linear photophysical properties**

16            To determine the linear photophysical properties, **HBO-NBu<sub>2</sub>** was dissolved in  
17 dichloromethane (DCM) at a concentration of approximately  $10.0^{-5} \text{ mol/L}$  and placed  
18 it in a quartz cell with a  $10.0 \text{ mm}$  path length. We collected the one-photon absorption  
19 (1PA) and fluorescence spectra using a UV-Vis spectrophotometer (UV-1800 model,  
20 Shimadzu) and a spectrofluorometer (F-7000 model, Hitachi), respectively. For the  
21 fluorescence measurements, we adjusted the solution concentration to an absorbance of  
22 approximately 0.5 to minimize fluorescence reabsorption. We determined the  
23 fluorescence quantum yield ( $\phi_{fl}$ ) using the well-known Brouwer method<sup>1,2</sup>, with  
24 coumarin 500 dissolved in methanol ( $\phi_{fl} = 68.1\%$ )<sup>3</sup> as the standard sample. We  
25 performed solvatochromism measurements by dissolving the molecule in different  
26 solvents, namely toluene, tetrahydrofuran (THF), dichloromethane (DCM),  
27 dimethylsulfoxide (DMSO), dimethylformamide (DMF), ethanol, acetonitrile (ACN),  
28 and methanol (see **Figs. SI1a, SI1b, and SI1c**). The solvatochromism measurement and  
29 the Lippert-Mataga equation<sup>4,5</sup> allowed us to determine the difference between the  
30 permanent dipole moment of the excited state and the ground one ( $|\Delta\vec{\mu}|$ )<sup>6</sup>. Finally, we  
31 measured the excitation anisotropy spectrum using the L-configuration<sup>7</sup>, in which we  
32 polarize the excitation and fluorescence emission beams in different directions  
33 (specifically in vertical and horizontal polarization combinations).

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**Figure S11** – a,b) Normalized absorbance and fluorescence spectra in different solvents, and c) Stokes shift versus the Lippert–Mataga polarity function.

### 39 SI2. Measurement of two- and three-photon absorption

40 We measured the two- and three-photon absorption (2PA and 3PA) spectrum  
41 using the multiphoton-excited fluorescence (MPEF) technique<sup>8–10</sup>. This method directly  
42 compares the fluorescence intensity excited by multiphoton absorption (MPA) with that  
43 of 1PA, using the latter to calibrate detection efficiency and, consequently,  $\phi_{fl}$ . Thus, we  
44 determined the 2PA and 3PA cross-sections ( $\sigma_{2PA}$  and  $\sigma_{3PA}$ ) by comparing the  
45 multiphoton excited fluorescence signal with the one-photon excited one, ensuring  
46 detection under identical experimental conditions. We obtained  $\sigma_{2PA}$  and  $\sigma_{3PA}$  using the  
47 following equations<sup>8</sup>:

48

$$\sigma_{2PA} = \frac{1}{\sqrt{2}} \left( \frac{\pi}{\ln(2)} \right)^2 \frac{P_{1PA}^3 (h\nu_{2PA})^2 F_{2PA}}{P_{2PA} h\nu_{1PA} F_{1PA}} (T\tau\Delta x\Delta y) \left( \frac{1 - 10^{-Abs(\lambda_{1PA})}}{Abs(\lambda_{1PA}) \cdot \ln(10)} \right) \sigma_{1PA}(\lambda_{1PA}), \quad (1)$$

49 and

$$\sigma_{3PA} = 3 \left( \frac{\sqrt{3}\pi}{4\ln(2)} \right)^3 \frac{P_{1PA} (h\nu_{3PA})^3 F_{3PA}}{(P_{3PA})^3 h\nu_{1PA} F_{1PA}} (T\tau\Delta x\Delta y)^2 \left( \frac{1 - 10^{-Abs(\lambda_{1PA})}}{Abs(\lambda_{1PA}) \cdot \ln(10)} \right) \sigma_{1PA}(\lambda_{1PA}). \quad (2)$$

50

51 In this equation,  $P_{nPA}$  represents the average power of the beam laser,  $h\nu_{nPA}$  is the photon  
52 energy, and  $F_{nPA}$  is the fluorescence signal. The subscript  $nPA$  refers to  $n$ -photon  
53 excitation, where  $n$  can be 1, 2, or 3.  $T$  denotes the laser repetition rate,  $\tau$  is the pulse  
54 duration (FWHM), and  $\Delta x$  and  $\Delta y$  represent the spatial width (FWHM) of the beam laser.  
55 Additionally,  $\sigma_{1PA}$  and  $Abs$  indicate the 1PA cross-section and the absorbance at the  
56 specific 1PA wavelength, respectively.

57

The experimental setup of the MPEF technique consists of a regenerative  
58 amplified Yb:KGW femtosecond laser system (Pharos PH1 model, Light Conversion,

59 *Inc.*) with the following specifications:  $800 \mu\text{J}$  pulse energy, with  $220 \text{ fs}$  pulse width,  
60 centered at  $1030 \text{ nm}$ , and repetition rate of  $7.5 \text{ kHz}$ . We used this laser system to pump  
61 an optical parametric amplifier (OPA) (Orpheus model, Light Conversion, *Inc.*). The  
62 OPA, combined with two second- and fourth-harmonic generation systems (Lyra-FH  
63 model, Light Conversion, *Inc.*), allows us to generate tunable pulses from  
64  $220 - 3000 \text{ nm}$ , with pulse energies between  $0.07$  and  $67 \mu\text{J}$  and  $100 - 180 \text{ fs}$  pulse  
65 width, characterized by an autocorrelator (GECO model, Light Conversion, *Inc.*). Below,  
66 we give a brief description of the experimental setup.

67 In our setup, we initially direct the beam emerging from the OPA to a spatial filter,  
68 using dichroic mirrors to eliminate residual wavelengths. The spatial filter generates a  
69 Gaussian intensity profile beam for all excitation wavelengths. We then align the  
70 excitation beam by 1PA, 2PA, and 3PA through two pinholes placed before and after the  
71 sample. This alignment ensures we collect the fluorescence signal under the same  
72 experimental conditions. We focused the laser beam on the sample using a  $f = 25 \text{ cm}$   
73 lens, positioned  $8 \text{ cm}$  before the focal point. A beam splitter between the lens and the  
74 sample cell directs  $< 4\%$  of the laser beam to a CCD camera to determine  $\Delta x$  and  $\Delta y$ .  
75 We collect the fluorescence signal perpendicular to the excitation beam, using the  
76 combination of two  $f = 5 \text{ cm}$  lenses and a  $R = 10 \text{ cm}$  spherical mirror, and focus the  
77 horizontally-elongated fluorescence image on the entrance plane of a monochromator.  
78 We coupled a photomultiplier tube to the monochromator, connected to a lock-in  
79 amplifier, and triggered it with the laser repetition rate to detect the fluorescence emission  
80 signal. Finally, we obtain the fluorescence signal from the average of  $10^4$  acquisitions,  
81 each with an integration time of  $1 \text{ s}$ .

82 We performed the 3PA measurement as described above, with the difference that,  
83 in this case, we focused the laser beam more strongly on the sample, positioned  $5 \text{ cm}$   
84 before the focal point.

85 We have taken all the precautions as discussed by Makarov *et al.*<sup>8</sup>, including  
86 keeping the monochromator slit opening the same 1PA, 2PA, and 3PA measurements.  
87 The photomultiplier voltage, integration time, and lock-in amplification are the same in  
88 all three situations. Furthermore, we have eliminated any spurious signals, such as laser  
89 scattering. Regarding sample preparation, we placed the solution in a quartz cell with a  
90  $10 \text{ mm}$  path length at the same concentration mentioned for 1PA. Finally, one performed

91 2PA measurements in the 550 – 900 *nm* range with a spectral resolution of 10 *nm* and  
92 by 3PA at 1030 *nm*.

93

### 94 **SI3. Femtosecond transient absorption spectroscopy**

95 We performed femtosecond transient absorption (TA) measurements using a  
96 homemade pump-probe setup based on the same laser system and OPA described in the  
97 previous section (at a repetition rate of  $\sim 1.9$  *kHz*). In this setup, we split the laser system  
98 beam using a 50% beam splitter. We use one of the beams to pump the OPA, generating  
99 a tunable pump pulse at 400 *nm* ( $\sim 100$  *fs*). We attenuate the second beam power (4.0  
100 mW) to generate a white-light continuous probe pulse (450 – 1030 *nm*) by focusing the  
101 beam onto a 2.0 *mm* thick sapphire window. We then focus both beams on the sample  
102 (spatially overlapping) and collect the spectrum of the transmitted probe pulse with a  
103 portable spectrometer (USB2000+ model, Ocean Optics). We vary the probe-probe time  
104 delay via a computer-controlled translation stage. We performed the measurements under  
105 magic angle polarization conditions ( $54.7^\circ$ ). Finally, we conducted the global analysis of  
106 the TA spectra using a custom-developed program, applying the methodology described  
107 by Van Stokkum *et al.*<sup>11</sup>.

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### 109 **SI4. Time-resolved fluorescence**

110 We determined the fluorescence lifetime ( $\tau_{fl}$ ) of the **HBO-NBu<sub>2</sub>** molecule by the  
111 time-resolved fluorescence technique using the same laser system and OPA described  
112 above (at a repetition rate of 300 *Hz*). We excited the sample at 400 *nm* and monitored  
113 the fluorescence signal as a function of the time. We obtained  $\tau_{fl}$  by the signal  
114 convolution method<sup>12</sup>, which convolutes the IRF with a monoexponential decay  
115 function<sup>6</sup>.

116

### 117 **SI5. Quantum chemical calculations**

118 To better understand the properties of **HBO-NBu<sub>2</sub>**, we performed quantum  
119 chemical calculations (QCC) based on the framework of density functional theory (DFT)  
120 and its time-dependent extension (TD-DFT), employing the Gaussian 09 software<sup>13</sup>. First,  
121 we conducted geometry optimization and vibrational frequency calculations using the  
122 hybrid functional M06-2X<sup>14</sup> and Pople's standard 6-311G(d,p) basis set<sup>15</sup>. We did not  
123 identify negative vibrational modes, indicating that the fully optimized structure

124 corresponds to a stable minimum. We then performed TD-DFT calculations to determine  
 125 the 20 lowest-energy singlet electronic transitions, employing the hybrid functional  
 126 B3LYP<sup>16</sup> and the Pople's extended 6-311G++(d,p) basis set. With the same methodology,  
 127 we also calculated the molecular orbitals and Le Bahers' parameters<sup>17</sup>, such as effective  
 128 charge displacement length ( $D_{CT}$ ), transferred charge ( $q_{CT}$ ) and  $|\Delta\vec{\mu}|$ . We performed all  
 129 QCC considering DCM as the solvent medium by using the Polarizable Continuum  
 130 Model (PCM) with integral equation formalism (IEF-PCM)<sup>18,19</sup>. We summarize in **Table**  
 131 **SI1** and **SI2** the optimized molecular geometry (equilibrium) and electronic transition  
 132 data of **HBO-NBu<sub>2</sub>**, respectively.

133

134 **Table SI1** – Optimized structures of **HBO-NBu<sub>2</sub>** obtained through IEFPCM-M062x/6-311G(d,p)  
 135 calculations in DCM medium.

<b>HBO-NBu<sub>2</sub></b>			
	IEFPCM-TD-M062x/6-311G(d,p)		
C	10.5684	-0.7757	0.7784
C	10.6737	0.1834	-0.2407
C	9.5509	0.7341	-0.8444
C	8.3095	0.2922	-0.3932
C	8.2313	-0.6615	0.6201
C	9.3335	-1.2240	1.2360
N	7.0108	0.6262	-0.7682
C	6.2461	-0.0982	-0.0093
O	6.9062	-0.9046	0.8583
C	4.7961	-0.1153	-0.0119
C	4.0950	0.7217	-0.9050
C	2.7010	0.7019	-0.9033
C	2.0042	-0.1333	-0.0319
C	2.7060	-0.9678	0.8584
C	4.0841	-0.9505	0.8596
C	0.5755	-0.1428	-0.0428
C	-0.6321	-0.1567	-0.0476
C	-2.0610	-0.1706	-0.0532
C	-2.7777	0.6475	-0.9382
C	-4.1622	0.6212	-0.9492

C	-4.8728	-0.2064	-0.0730
C	-4.1547	-1.0191	0.8118
C	-2.7676	-1.0079	0.8185
N	-6.2974	-0.1913	-0.1130
C	-7.0286	4.7478	0.7046
C	-7.3368	3.3131	1.1260
C	-6.5756	2.2877	0.2865
C	-6.8736	0.8643	0.7397
C	-10.2321	-3.1733	-1.0236
C	-8.9002	-3.0061	-0.2960
C	-8.2653	-1.6410	-0.5560
C	-6.9237	-1.4918	0.1556
H	11.4712	-1.1779	1.2205
H	11.6579	0.5002	-0.5625
H	9.6278	1.4739	-1.6305
H	9.2368	-1.9624	2.0203
H	2.1741	1.3485	-1.5930
H	2.1595	-1.6137	1.5321
H	4.6371	-1.5869	1.5390
H	-2.2384	1.2884	-1.6250
H	-4.7170	1.2308	-1.6518
H	-4.6836	-1.6664	1.5014
H	-2.2211	-1.6431	1.5050
H	-7.5821	5.4695	1.3080
H	-7.2952	4.9094	-0.3430
H	-5.9628	4.9637	0.8146
H	-8.4123	3.1263	1.0391
H	-7.0811	3.1769	2.1821
H	-5.5035	2.4940	0.3581
H	-6.8538	2.3905	-0.7681
H	-6.5375	0.7300	1.7823
H	-7.9568	0.7235	0.7337
H	-10.6726	-4.1535	-0.8318

H	-10.0990	-3.0691	-2.1034
H	-10.9484	-2.4132	-0.7012
H	-8.2069	-3.7929	-0.6117
H	-9.0500	-3.1371	0.7809
H	-8.9604	-0.8582	-0.2364
H	-8.1042	-1.5024	-1.6297
H	-7.0558	-1.6612	1.2372
H	-6.2531	-2.2723	-0.2124
O	4.7052	1.5471	-1.7678
H	5.6760	1.4616	-1.6656

136

137 **Table S12** – Theoretical photophysical properties of **HBO-NBu<sub>2</sub>** molecule obtained by IEFPCM-TD-  
 138 B3LYP/6-311++G(d,p) calculation in DCM medium, such as electronic state, transition energy, transition  
 139 wavelength, and oscillator strength.

<b>HBO-NBu<sub>2</sub></b>			
IEFPCM-TD-B3LYP/6-311++G(d,p)			
Electronic state	Energy (eV)	Wavelength h (nm)	Oscillator Strength
S <sub>0</sub> -S <sub>1</sub>	3.1647	391.77	1.5185
S <sub>0</sub> -S <sub>2</sub>	3.6818	336.75	0.3616
S <sub>0</sub> -S <sub>3</sub>	4.0260	307.96	0.1658
S <sub>0</sub> -S <sub>4</sub>	4.2857	289.30	0.0847
S <sub>0</sub> -S <sub>5</sub>	4.3659	283.98	0.0034
S <sub>0</sub> -S <sub>6</sub>	4.4335	279.65	0.0119
S <sub>0</sub> -S <sub>7</sub>	4.4599	278.00	0.0157
S <sub>0</sub> -S <sub>8</sub>	4.6161	268.59	0.0000
S <sub>0</sub> -S <sub>9</sub>	4.6960	264.02	0.0331
S <sub>0</sub> -S <sub>10</sub>	4.7474	261.16	0.0217
S <sub>0</sub> -S <sub>11</sub>	4.8511	255.58	0.0243
S <sub>0</sub> -S <sub>12</sub>	5.0113	247.41	0.0239
S <sub>0</sub> -S <sub>13</sub>	5.0472	245.65	0.0019
S <sub>0</sub> -S <sub>14</sub>	5.1138	242.45	0.0089
S <sub>0</sub> -S <sub>15</sub>	5.1996	238.45	0.0065
S <sub>0</sub> -S <sub>16</sub>	5.2193	237.55	0.0242

S <sub>0</sub> -S <sub>17</sub>	5.2453	236.37	0.0158
S <sub>0</sub> -S <sub>18</sub>	5.2764	234.98	0.0652
S <sub>0</sub> -S <sub>19</sub>	5.3469	231.88	0.0150
S <sub>0</sub> -S <sub>20</sub>	5.3782	230.53	0.0044

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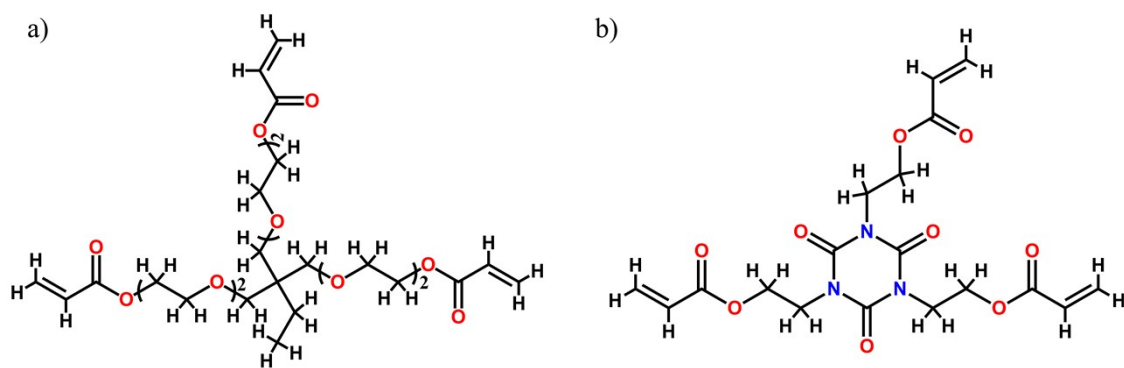
## 141 **SI6. Multiphoton polymerization**

142 We employed a homemade setup to investigate the viability of **HBO-NBu<sub>2</sub>** as a  
 143 photoinitiator for MPP. Depending on the absorption process, we used two different laser  
 144 systems. For 2PP, we used an amplified femtosecond erbium laser system (FemtoFiber  
 145 Scientific model, Toptica Photonics, *Inc.*) with a 200 fs pulse width at 775 nm (second  
 146 harmonic of 1550 nm), operating at a repetition rate of 80 MHz (0.75 nJ/pulse). For  
 147 3PP, we used an amplified Yb:KGW femtosecond laser system (Carbide, Light  
 148 Conversion, *int.*) with a 244 fs pulse width at 1030 nm and operating at a repetition rate  
 149 of 1 MHz (5.0 μJ/pulse). We directed the beam to a rotating polarizer and a half-wave  
 150 plate to control the excitation power. We focused the beam onto the sample with an  
 151 objective lens (62 ×, NA = 0.80) and positioned the sample on a computer-controlled  
 152 x-y-z translation stage, set at a writing speed of 20 μm/s. We adjusted the incident power  
 153 from 6 to 14 mW (0.075 to 0.18 nJ) for 2PA and used a fixed power of 6.0 mW (6 nJ)  
 154 for 3PA.

155 The sample used here consists of a mixture of two commercial acrylic resins,  
 156 *ethoxylated(6)trimethylolpropane triacrylate* (SR499 – Sartomer) and *tris(2-*  
 157 *hydroxyethyl)isocyanurate triacrylate* (SR368 – Sartomer) (see **Fig. SI2**), employed in  
 158 equal mass proportions (wt%). This combination provides mechanical stability to the  
 159 final structure and minimizes shrinkage after polymerization<sup>20</sup>. We incorporated different  
 160 concentrations of **HBO-NBu<sub>2</sub>**, previously dissolved in DCM (0.01 to 0.08 wt%). After  
 161 solvent evaporation, we placed a small sample drop on a glass substrate with a  
 162 micrometric spacer and covered it with a coverslip. After completing the microfabrication  
 163 process, we washed the samples by immersing them in ethanol heated to 75 °C for 15  
 164 min. We repeated this procedure three times to ensure the complete removal of residues.

165





SR499 – ethoxylated(6)trimethylolpropane triacrylate

SR368 – tris(2-hydroxyethyl)isocyanurate triacrylate

166  
 167 **Figure SI2** – Representation of the chemical structure of the monomers: a) ethoxylated(6)  
 168 trimethylolpropane triacrylate (SR499 – Sartomer) and b) tris(2-hydroxyethyl) isocyanurate triacrylate  
 169 (SR368 – Sartomer).  
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## 171 SI7. REFERENCES

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