

Size of quorum sensing communities

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Cylindrical colony, reflective boundary

In this supplementary we will derive the steady state concentration of a circular colony with radius $r_0 = \mathcal{R}$ and height $z_0 = \mathcal{H}$ growing on a signal-reflecting surface. The signal concentration is denoted s , D its diffusion constant, and κ_s the intracellular signal production rate per volume. When ρ_v is the (v/v) cell density the volume normalized production is $k(r, z) = \rho_v(r, z)\kappa_s(r, z)$.

The steady state diffusion equation in cylindrical coordinates (assuming rotational symmetry) reads

$$\frac{\partial^2 s}{\partial r^2} + \frac{1}{r} \frac{\partial s}{\partial r} + \frac{\partial^2 s}{\partial z^2} = -\frac{k(r, z)}{D} \quad (27)$$

where $k(r, z)$ is the source term, which is assumed to be homogeneous in the colony, $k(r, z) = k\theta(z_0 - z)\theta(r_0 - r)$ and $\theta(x)$ is the Heaviside unit step function. Eq. (27) is most easily solved by using the Hankel transformation (of order zero)³³.

$$\tilde{f}(q) = \mathcal{F}(f)(q) = \int_0^\infty f(r)J_0(qr)r dr$$

where $J_0(x)$ is the Bessel function of order 0. A useful property of the Hankel transformation is

$$\mathcal{F}(f'' + r^{-1}f')(q) = -q^2\mathcal{F}(f)(q) = -q^2\tilde{f}(q)$$

Applying the Hankel transformation to (27) therefore gives

$$-q^2\tilde{s}(q, z) + \frac{d^2\tilde{s}(q, z)}{dz^2} = -\tilde{k}(q, z) \quad (28)$$

where

$$\begin{aligned} \tilde{k}(q, z) &= \frac{k}{D}\theta(z_0 - z) \int_0^{r_0} J_0(qr)r dr \\ &= \frac{k}{D}\theta(z_0 - z)q^{-2} \int_0^{qr_0} J_0(x)xdx \end{aligned}$$

Set $A(qr_0) = \int_0^{qr_0} J_0(x)xdx$, so

$$-q^2\tilde{s}(q, z) + \frac{d^2\tilde{s}(q, z)}{dz^2} = -\frac{k\theta(z_0 - z)A(qr_0)}{Dq^2} \quad (29)$$

The solution to the homogeneous part of (29) is

$$\tilde{s}_h(q, z) = a(q)e^{-qz} + b(q)e^{qz}$$

where a and b are integration constants. The particular solution can be chosen as

$$\tilde{s}_p(q, z) = \frac{k\theta(z_0 - z)A(qr_0)}{Dq^4}$$

The full solution therefore becomes

$$\tilde{s}(q, z) = \begin{cases} a_-(q)e^{-qz} + b_-(q)e^{qz} + \frac{kA(qr_0)}{Dq^4} & , z < z_0 \\ a_+(q)e^{-qz} + b_+(q)e^{qz} & , z > z_0 \end{cases} \quad (30)$$

The integration constants are determined from the boundary conditions. Since $\tilde{s} \rightarrow 0$ for $z \rightarrow \infty$, we have $b_+(q) = 0$. The reflective boundary condition gives

$$\frac{\partial \tilde{s}(0)}{\partial z} = 0 \Rightarrow a_-(q) = b_-(q)$$

The two pieces of the solution has to be continuous with continuous derivative at $z = z_0$. This second condition implies

$$\begin{aligned} a_-(q)(e^{qz_0} - e^{-qz_0}) &= -a_+(q)e^{-qz_0} \\ \Rightarrow a_+(q) &= a_-(q)(1 - e^{2qz_0}) \end{aligned}$$

The first condition gives

$$\begin{aligned} a_-(q)(e^{-qz_0} - e^{qz_0}) &= a_-(q)(e^{-qz_0} + e^{qz_0}) + \frac{kA(qr_0)}{Dq^4} \\ \Rightarrow a_-(q) &= -\frac{kA(qr_0)e^{-qz_0}}{2Dq^4} \end{aligned}$$

Inserting into (30) we obtain

$$\tilde{s}(q, z) = \frac{kA(qr_0)}{2Dq^4} \begin{cases} 2 - e^{-q(z+z_0)} - e^{q(z-z_0)} & , z < z_0 \\ e^{-q(z-z_0)} - e^{-q(z+z_0)} & , z > z_0 \end{cases}$$

For convenience we define

$$\begin{aligned} I_-(q, z) &= 2 - e^{-q(z+z_0)} - e^{q(z-z_0)} \\ I_+(q, z) &= e^{-q(z-z_0)} - e^{-q(z+z_0)} \end{aligned}$$

To obtain

$$s(r, z) = \begin{cases} s_-(r, z) & \text{for } 0 \leq z \leq z_0 \\ s_+(r, z) & \text{for } z > z_0 \end{cases}$$

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we need to apply the inverse Hankel transformation again.

$$s_{\pm}(r, z) = \frac{k}{2D} \int_0^{\infty} \frac{A(qr_0)}{q^4} I_{\pm}(q, z) J_0(qr) q dq$$

Using that $A(qr_0) = qr_0 J_1(qr_0)$, we finally get

$$s_{\pm}(r, z) = \frac{kr_0}{2D} \int_0^{\infty} \frac{I_{\pm}(q, z)}{q^2} J_0(qr) J_1(qr_0) dq \quad (31)$$

Note that $I_{\pm}(q, z) > 0$ and $I_{\pm}(q, z) = 2qz_0 + \mathcal{O}(q^2)$. This ensures that s is well-defined ($J_1(x) = \mathcal{O}(x)$) and positive. To make the dimensionality of (31) clearer, set $x = qr_0$, $r' = r/r_0$, $z' = z/r_0$ and $z'_0 = z_0/r_0$ (aspect ratio) to obtain

$$s_{\pm}(r', z') = \frac{kr_0^2}{2D} \int_0^{\infty} \frac{I_{\pm}(x, z')}{x^2} J_0(xr') J_1(x) dx \quad (32)$$

where

$$\begin{aligned} I_{-}(x, z') &= 2 - e^{-x(z'+z'_0)} - e^{x(z'-z'_0)} \\ &= 2(1 - e^{-xz'_0} \cosh(xz')) \\ I_{+}(x, z') &= e^{-x(z'-z'_0)} - e^{-x(z'+z'_0)} \\ &= 2e^{-xz'} \sinh(xz'_0) \end{aligned}$$

The integral in (32) can be evaluated in specific points. At the top, $r' = 0$, $z' = z'_0$ we have

$$\begin{aligned} s_{-}(0, z'_0) &= \frac{kr_0^2}{D} \int_0^{\infty} \frac{1 - e^{-xz'_0} \cosh(xz'_0)}{x^2} J_1(x) dx \\ &= \frac{kr_0^2}{2D} \left(2z'_0(-2z'_0 + \sqrt{4z_0'^2 + 1}) + \sinh^{-1}(2z'_0) \right) \\ &\approx \frac{kr_0^2}{2D} \left(2z'_0 - 2z_0'^2 + \frac{4}{3}z_0'^3 + \mathcal{O}(z_0'^4) \right) \end{aligned}$$

$$s_{-}(0, z'_0) = 0.958 \frac{kr_0^2}{2D}, \quad z'_0 = 1$$

$$s_{-}(0, z'_0) = 2 \frac{kr_0 z_0}{2D}, \quad z'_0 \ll 1$$

At the border of the supporting surface, $r' = 1$, $z' = 0$, we get

$$\begin{aligned} s_{-}(1, 0) &= \frac{kr_0^2}{D} \int_0^{\infty} \frac{1 - e^{-xz'_0}}{x^2} J_0(x) J_1(x) dx \\ &\approx \frac{kr_0^2}{2D} (1.273z'_0 - 0.5z_0'^2 + \mathcal{O}(z_0'^3)) \end{aligned}$$

$$s_{-}(1, 0) = 0.926 \frac{kr_0^2}{2D}, \quad z'_0 = 1$$

$$s_{-}(1, 0) = 1.273 \frac{kr_0 z_0}{2D}, \quad z'_0 \ll 1$$

Finally, in the center, we have

$$\begin{aligned} s_{-}(0, 0) &= \frac{kr_0^2}{D} \int_0^{\infty} \frac{1 - e^{-xz'_0}}{x^2} J_1(x) dx \\ &= \frac{kr_0^2}{2D} \left(z'_0(-z'_0 + \sqrt{z_0'^2 + 1}) + \sinh^{-1}(z'_0) \right) \\ &\approx \frac{kr_0^2}{2D} \left(2z'_0 - z_0'^2 + \frac{z_0'^3}{3} + \mathcal{O}(z_0'^5) \right) \end{aligned}$$

$$s_{-}(0, 0) = 1.296 \frac{kr_0^2}{2D}, \quad z'_0 = 1$$

$$s_{-}(0, 0) = 2 \frac{kr_0 z_0}{2D}, \quad z'_0 \ll 1$$

Note that in all points the algebraic prefactor is of order unity. In order to get the expression needed for a thin biofilm in the main text, recall that $z_0 = \mathcal{H}$, $r_0 = \mathcal{R}$, and $k = \rho_v \kappa_s$. Specifically the last equation reads

$$s(0, 0) = 2 \frac{\kappa_s \mathcal{R} \mathcal{H} \rho_v}{2D}, \quad \mathcal{H} \ll \mathcal{R} \quad (33)$$