Supporting Information

Length evolution of helical micro/nano-scale structures

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The dependence of L on h in both CLHL and COHL cases.

Both *L*_{CLHL} and *L*_{COHL} increase linearly with increasing *h*, as shown in Figures S1 and S2. The plots show a non-linear dependence of *L* on *d*. Just as mentioned in the manuscript, there exists a critical point for *d*, below which the increasing slope of *L* with increasing h is much slower; above the critical point of *d*, the increase becomes much steeper with the increase of *h*.

 Figure S1│ Plots of *L*CLHL vs. *h* with fixed *R* and *d*.

Figure S2 | Plots of L_{COHL} vs. *h* with fixed *R* and *d*.

Derivation of the helical length formula of a COHL structure

The equation of the COHL can be parameterized using the rotation angle θ in a cylindrical coordinate. Therefore, the parametric equation of the COHL can be described in Eq. 1,

$$
\begin{cases}\n x(\theta) = R(\theta)\cos(\theta) \\
y(\theta) = R(\theta)\sin(\theta), \quad 0 \le \theta \le 2k\pi \\
z(\theta) = v\theta\n\end{cases}
$$
\n(1)

where $v = \frac{d}{2\pi}$ and $R(\theta) = R - v\theta \cot \alpha$.

The helical length of the COHL (L_{COHL}) can be obtained through the curve integral as exhibited in Eq. 2.

$$
L_{COHL} = \int_0^{2k\pi} \left\{ [x'(\theta)]^2 + [y(\theta)]^2 + [z(\theta)]^2 \right\}^{1/2} d\theta
$$
 (2)

Substituting Eq.1 into Eq. 2, we then obtain Eq. 3,

$$
L_{\text{COHL}} = v \int_0^{2k\pi} \left[\theta^2 \cot^2 \alpha - \theta \cot \alpha \frac{2R}{v} + c\right]^{1/2} d\theta
$$

where $k = \frac{h}{d}$ and $c = 1 + \cot^2 \alpha + (\mathbb{R}/v)^2$.

To simplify the derivation, we denote the coefficient $\cot^2 \alpha$ as $a,(2R/v)\cot \alpha$ as *b* and the integrand $(a\theta^2 + b\theta + c)^{1/2}$ as *F*. Therefore, Eq. 3 can be solved through the following integrals in Eqs. 4 and 5 by assuming $a > 0$ and $4ac-b^2 > 0$.

$$
\int F d\theta = \frac{2a\theta - 3b}{4a^2} F + \frac{3b^2 - 4ac}{8a^2} \int \frac{d\theta}{F}
$$
\n
$$
\int \frac{d\theta}{F} = \frac{1}{\sqrt{a}} \operatorname{arsinh} \frac{2a\theta + b}{\sqrt{4ac - b^2}}
$$
\n(4)

Substituting Eq. 4 into Eq. 5, we get Eq. 6.

$$
\int F d\theta = \frac{2a\theta - 3b}{4a^2} F + \frac{3b^2 - 4ac}{8a^{2.5}} \operatorname{arsinh} \frac{2a\theta + b}{\sqrt{4ac - b^2}}
$$

(6)

By replacing *a*, *b*, *c* and *F* in Eq. 6 with the original form, we finally obtain the helical length L_{COHL} as shown in Eq. 7.

 L_{COHL}

$$
= \left[\frac{v\theta cot\alpha - R}{2cot\alpha}\sqrt{(\theta cot\alpha - \frac{R}{v})^2 + 1 + cot^2\alpha} + \frac{v(1 + cot^2\alpha)}{2cot\alpha}arsinh(\frac{v\theta cot\alpha - R}{v\sqrt{1 + cot^2\alpha}})\right]\Big|_{0}^{2k}
$$

(7)

Simplifying the result by replacing $R - v\theta \cot \alpha$ with R(θ), we get Eq. 8.

$$
L_{COHL} = \left[\frac{-R(\theta)}{2\cot\alpha} \sqrt{\frac{R^2(\theta)}{v^2} + 1 + \cot^2\alpha} + \frac{v(1 + \cot^2\alpha)}{2\cot\alpha} \arcsin\left(\frac{-R(\theta)}{v\sqrt{1 + \cot^2\alpha}}\right) \right]_{0}^{2k\pi}
$$
\n(8)

By replacing *v* with $v = d/2\pi$, and R(θ) with $R - v\theta \cot \alpha$, we get the final form of the ℎ

helical length L_{COHL} , written as Eq. 9, where $k=\frac{1}{1}$ d_{\perp} L_{COHL}

$$
= \left[\frac{\theta d \cot \alpha - R}{4\pi \cot \alpha}\sqrt{\frac{(4\pi^2 R - d \nu \theta \cot \alpha)^2}{d^2} + 1 + \cot^2 \alpha} + \frac{d(1 + \cot^2 \alpha)}{4\pi \cot \alpha} \right. \frac{d(\theta d \cos \alpha)}{4\pi^2 \sqrt{1}}.
$$